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Abdelhak MEKAHLIA, Eric SEMAIL, Hussein ZAHR, Franck SCUILLER - Reduced-Order Model of Rotor Cage in Multiphase Induction Machines: Application on the Prediction of Torque Pulsations - Mathematical and computational applications - Vol. 25, n°1, p.11 - 2020

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Reduced-Order Model of Rotor Cage in Multiphase Induction Machines: Application on the Prediction of Torque Pulsations

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Abstract: For three-phase induction machines supplied by sinusoidal current, it is usual to model the n-bar squirrel-cage by an equivalent two-phase circuit. For a multiphase induction machine which can be supplied with different harmonics of current, the reduced-order model of the rotor must be more carefully chosen in order to predict the pulsations of torque. The proposed analysis allows to avoid a wrong design with non-sinusoidal magnetomotive forces. An analytical approach is proposed and confirmed by Finite-Element modelling at first for a three-phase induction machine and secondly for a five-phase induction machine.

Keywords: multiphase induction machine; reduced order; rotor cage; torque pulsations

1. Introduction

In electromechanical energy conversion, the three-phase induction machines are the most used. The rules and methods of the design of this kind of machines have been widely investigated in [1] and [2].

The mathematical modeling of induction machines is characterized by the complexity of the squirrel cage rotor. Modeling and parameters determination of this kind of rotors is investigated in [3].

The three-phase induction machines with classical design and sinusoidal currents are limited regarding torque density and speed range [4,5]. With multiphase induction machines, non-sinusoidal stator current harmonics allow to improve the torque density [6–9]. The speed range can be extended thanks to the sequential injection of current harmonics [10,11]. Several works have also investigated the speed range extension in multiphase induction machines with two polarities by pole-phase modulation [12–14].

Thanks to the aforementioned advantages, multiphase induction machines are becoming attractive in several sectors especially aeronautics [15], naval [16], and automotive industries [17]. With these increasing demands on multiphase induction machines, the design rules and theories, classically developed for the case of three phases, must be reconsidered.

In fact, the induction machine is characterized by complex harmonic interactions due to the squirrel cage rotor structure, if this one is not well-designed dangerous phenomena could occur, especially torque ripple and mechanical vibrations. For the three-phase machines several design rules have been developed [2,18]. It must be mentioned that in the case of a three-phase machine, only one stator current harmonic is used to produce torque, which is not the case for a multiphase machine.

With the possibility of injecting more than one harmonic, more degrees of freedom regarding the machine supply and control are appearing [19,20]. However, on the other hand, most of the classical design approaches, adapted for three-phase machines, cannot be directly used for multiphase machines due to the impact of injecting several stator harmonics on the design parameters [21]. Furthermore, it becomes more difficult to design a squirrel cage rotor without the risk of parasitic phenomena occurring during all the supplying modes (with different sequences). As an example, a previous work has shown that the rotor bars number has to be carefully chosen to avoid important torque ripple under different supplying harmonics (sequences) [22]. Hence, it is necessary to understand deeply the interactions between time and space harmonics and their impact on torque production in multiphase induction machine [23].

Generalized mathematical modeling, based on voltage and flux equations, allows to interpret these harmonic interactions. Thanks to the symmetrical component transformations, the study of space and time harmonics is easier [24]. This mathematical approach requires to identify the machine parameters, inductances and resistances, several works were done in this topic [25,26].

In three-phase induction machines, the stator winding, whose dimension is $N_{ph} = 3$, can be modeled by one two-phase equivalent circuit corresponding to one α - β plane and a zero-sequence (not excited for balanced stator supply) thanks to α - β -0 Concordia Transformation. This mathematical transformation is also applied on the squirrel cage rotor, whose dimension is N_{bar} , which gives:

- $\frac{N_{bar}-1}{2}$ α - β rotor planes, excited by independent harmonic sets, and **one** zero-sequence if N_{bar} is odd
- $\frac{N_{bar}}{2} - 1$ α - β rotor planes and **two** zero-sequences if N_{bar} is even.

However, in the classical simplified modeling approaches, the rotor cage is modeled by only one equivalent two-phase circuit, which corresponds to the α - β plane excited by the fundamental sequence (only the first space harmonic considered). This simplified approach supposes a sinusoidal distribution of the winding, so other space harmonics can be neglected comparing to the fundamental. In fact, this simplification is not always available, especially for the windings with a non-sinusoidal distribution, as the tooth concentrated windings, where all space harmonics exist and with significant amplitudes, which induce in the rotor bars more than one current harmonic. Hence different α - β rotor planes are excited.

Furthermore, in the case of multiphase machines (more than 3 phases), the stator winding is modeled by:

- $\frac{N_{ph}-1}{2}$ α - β stator planes, and **one** zero-sequence if N_{ph} is odd;
- $\frac{N_{ph}}{2} - 1$ α - β stator planes, and **two** zero-sequences if N_{ph} is even.

Every α - β stator plane can be excited by a stator sequence, usually several current harmonics are injected to improve torque production in multiphase machines, these injected harmonics belong to independent sets, so they excite different α - β stator planes. All these excited α - β planes must be considered in modeling, so the stator winding is represented by more than one equivalent two-phase circuit, each one interacts with one α - β rotor plane at least, depending on the induced rotor current harmonics and the importance of their amplitudes.

Considering all the α - β rotor planes, excited by important space harmonics, is important to predict torque pulsations due to harmonic interactions. Recent work has investigated the analytical estimation of torque ripples, due to interactions between stator and rotor space harmonics, for three-phase induction machine [23].

In this paper, a mathematical approach allowing to find the excited α - β rotor planes which must be considered in the modelling of multiphase induction machine is proposed. This approach allows to justify the reduction of the rotor dimension, which is " N_{bar} " in the natural base model (Non-transformed electrical equations). Firstly, the generalized mathematical model of induction machine is presented. Then, the α - β transformation is applied on the stator and rotor equations.

Thanks to the transformed stator and rotor current vectors and the mutual inductance matrix, a new arithmetic approach of prediction of torque pulsating components frequencies is proposed.

This arithmetic tool is used to predict pulsating components frequencies for firstly a three-phase machine, and secondly a five-phase machine supplied by two different sequences. The results are validated by F-E simulations.

2. Induction Machine Model

Based on inductance and resistance matrices in stator and rotor, voltage equations can be written as follows [27]:

$$\underline{V}_s = \underline{R}_s \cdot \underline{I}_s + \frac{d\phi_s}{dt}, \quad (1)$$

$$\underline{V}_r = 0 = \underline{R}_r \cdot \underline{I}_r + \frac{d\phi_r}{dt}, \quad (2)$$

$$\underline{\phi}_s = \underline{L}_{ss} \cdot \underline{I}_s + \underline{L}_{sr}(\theta) \cdot \underline{I}_r, \quad (3)$$

$$\underline{\phi}_r = \underline{L}_{rr} \cdot \underline{I}_r + \underline{L}_{rs}(\theta) \cdot \underline{I}_s, \quad (4)$$

The stator matrix can be written as follows:

$$\underline{R}_s = \begin{bmatrix} R_s & 0 & \dots & 0 \\ 0 & R_s & 0 & \vdots \\ \vdots & \ddots & \ddots & 0 \\ 0 & \dots & 0 & R_s \end{bmatrix}, \quad (5)$$

$$\underline{L}_{ss} = \begin{bmatrix} L_{s11} + L_{ls} & L_{s12} & \dots & \dots & L_{s1N_{ph}} \\ L_{s1N_{ph}} & L_{s11} + L_{ls} & L_{s12} & \dots & L_{s1(N_{ph}-1)} \\ \vdots & \ddots & \ddots & \ddots & \vdots \\ L_{s12} & \dots & \dots & L_{s1N_{ph}} & L_{s11} + L_{ls} \end{bmatrix}, \quad (6)$$

The rotor cage is represented as shown in Figure 1. Each rotor loop (composed by two adjacent half-bars and two ring portions) is considered like a phase [28].

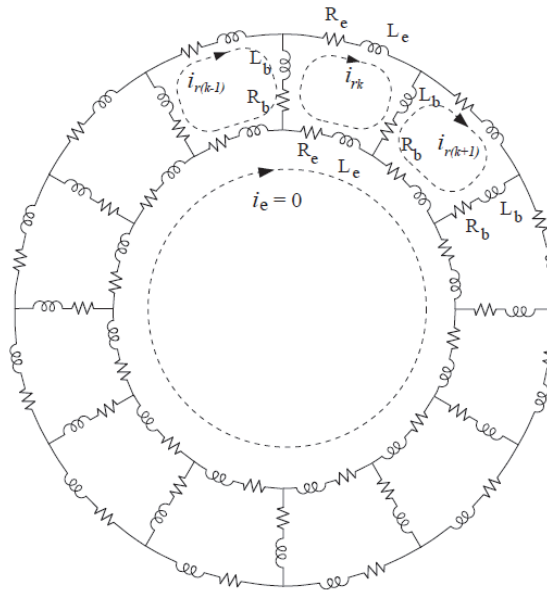


Figure 1. Squirrel cage rotor representation, figure taken from [29] (p. 60).

According to the rotor representation approach shown in Figure 1, rotor resistance and inductance matrices can be written as follows:

$$\underline{R}_r = \begin{bmatrix} 2(R_b + R_e) & -R_b & 0 & \dots & -R_b \\ -R_b & 2(R_b + R_e) & -R_b & \ddots & 0 \\ \vdots & \ddots & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & \ddots & \vdots \\ -R_b & 0 & \dots & -R_b & 2(R_b + R_e) \end{bmatrix}, \quad (7)$$

$$\underline{L}_{rr} = \begin{bmatrix} L_{r11} + 2(L_b + L_e) & L_{r12} - L_b & L_{r13} & \dots & L_{r1(N_{bar}-1)} & L_{r1N_{bar}} - L_b \\ L_{r21} - L_b & L_{r11} + 2(L_b + L_e) & L_{r23} - L_b & L_{r24} & \dots & L_{r2N_{bar}} \\ \vdots & \ddots & \ddots & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & \ddots & \ddots & \vdots \\ L_{rN_{bar}1} - N_b & L_{rN_{bar}2} & \dots & L_{rN_{bar}(N_{bar}-2)} & L_{rN_{bar}(N_{bar}-1)} - N_b & L_{r11} + 2(L_b + L_e) \end{bmatrix}, \quad (8)$$

Mutual inductance between stator and rotor can be written as follows:

$$\underline{L}_{sr}(\theta) = \sum_k \text{real} \left(\hat{M}_{sr}_k \begin{bmatrix} e^{-j.k.p\theta} & \dots & e^{j.k.(N_{bar}-1)\frac{2\pi}{N_{bar}}-p\theta} \\ \vdots & \ddots & \vdots \\ e^{j.k.(N_{ph}-1)\frac{2\pi}{N_{ph}}-p\theta} & \dots & e^{j.k.(N_{ph}-1)\frac{2\pi}{N_{ph}}-p(N_{bar}-1)\frac{2\pi}{N_{bar}}-p\theta} \end{bmatrix} \right) \quad (9)$$

Mutual inductance between a stator phase and a rotor phase (rotating) is a periodic function depending on the electrical angle (mechanical angle multiplied by the number of pole-pairs) and whose period is an electrical turn. This periodic spatial function contains harmonics, called space harmonics. Hence, Fourier transformation of the mutual inductance spatial function allows to find the complex coefficients $(\hat{M}_{sr})_k$, where “k.p” is the space harmonic range (k is positive integer).

As we can see in Equation (9), to simplify the expression of the mutual inductance matrix, the complex coefficients $(\hat{M}_{sr})_k$ are used. For example, the first element (1,1) of the matrix $L_{sr}(\theta)$ from Equation (9) is: $\text{real}((\hat{M}_{sr})_k.e^{-j.k.p\theta})$, which is equal to: $M_{srk}.\cos(-kp\theta + \varphi_{srk})$, where M_{srk} is the amplitude of the complex coefficient $(\hat{M}_{sr})_k$, and φ_{srk} is its angle.

3. α - β -0 Transformed Model

α - β -0 generalized transformation matrix [24], for the **odd** dimension “n”, is defined as follows:

$$\underline{A}_n = \sqrt{\frac{2}{n}} \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & \dots & \frac{1}{\sqrt{2}} \\ 1 & \cos\left(\frac{2\pi}{n}\right) & \dots & \cos\left(\frac{2(n-1)\pi}{n}\right) \\ 0 & \sin\left(\frac{2\pi}{n}\right) & \dots & \sin\left(\frac{2(n-1)\pi}{n}\right) \\ \vdots & \ddots & \dots & \vdots \\ 1 & \cos\left(\frac{(n-1)\pi}{n}\right) & \dots & \cos\left(\frac{(n-1)^2\pi}{n}\right) \\ 0 & \sin\left(\frac{(n-1)\pi}{n}\right) & \dots & \sin\left(\frac{(n-1)^2\pi}{n}\right) \end{bmatrix}, \quad (10)$$

If the dimension “n” is **even**, a new line is added at the end of the matrix \underline{A}_n (last line). This line has this form: $\frac{2}{\sqrt{n}} \left[\frac{1}{\sqrt{2}} \quad -\frac{1}{\sqrt{2}} \quad \dots \quad \frac{1}{\sqrt{2}} \quad -\frac{1}{\sqrt{2}} \right]$.

Equations (1)–(4) are transformed as follows:

$$\underline{V}_{s\alpha\beta} = \underline{A}_s \cdot \underline{R}_s \cdot \underline{A}_s^{-1} \cdot \underline{I}_{s\alpha\beta} + \frac{d\phi_{s\alpha\beta}}{dt}, \quad (11)$$

$$\underline{V}_{r\alpha\beta} = 0 = \underline{A}_r \cdot \underline{R}_r \cdot \underline{A}_r^{-1} \cdot \underline{I}_{r\alpha\beta} + \frac{d\phi_{r\alpha\beta}}{dt}, \quad (12)$$

$$\underline{\phi}_{s\alpha\beta} = \underline{A}_s \cdot \underline{L}_{ss} \cdot \underline{A}_s^{-1} \cdot \underline{I}_{s\alpha\beta} + \underline{A}_s \cdot \underline{L}_{sr}(\theta) \cdot \underline{A}_r^{-1} \cdot \underline{I}_{r\alpha\beta}, \quad (13)$$

$$\underline{\phi}_{r\alpha\beta} = \underline{A}_r \cdot \underline{L}_{rr} \cdot \underline{A}_r^{-1} \cdot \underline{I}_{r\alpha\beta} + \underline{A}_r \cdot \underline{L}_{rs}(\theta) \cdot \underline{A}_s^{-1} \cdot \underline{I}_{s\alpha\beta}, \quad (14)$$

where \underline{A}_s is α - β -0 transformation matrix with the stator dimension " N_{ph} ", and \underline{A}_r is related to the rotor dimension " N_{bar} ".

To investigate the rotor dimension reduction, this transformation is applied on the stator-rotor mutual inductance matrix $\underline{L}_{sr}(\theta)$. The use of the transformed matrix, and the transformed stator and rotor currents vectors, lead to the prediction of torque components frequencies.

3.1. α - β -0 Transformation of the Stator Current Vector

The stator winding can be supplied by different sequences " u ", defined as follows:

$$u \in [0, 1 \dots N_{ph} - 1] \quad (15)$$

For example, in the case of classical three-phase machines, the sequence " $u = 0$ " corresponds to the zero-sequence, the sequence " $u = 1$ " corresponds to the direct order of supplying the phases and the sequence " $u = 2$ " corresponds to the inverse order.

If we consider that the winding is supplied by one sequence " u ", the current of the stator phase " i " in the natural base, can be expressed as follows:

$$I_{s_i} = I_{s_u} \cdot \sin\left(2\pi \cdot f_{s_u} \cdot t - (i - 1) \cdot u \cdot \frac{2\pi}{N_{ph}} + \varphi_{s_u}\right), \quad (16)$$

The stator current vector in the natural base can be expressed as follows:

$$\underline{I}_s = \begin{pmatrix} I_{s_0} \\ I_{s_1} \\ \vdots \\ I_{s_{N_{ph}-1}} \end{pmatrix}, \quad (17)$$

If only one sequence " u " is supplying the stator, Concordia transformation reduces the dimension of the current vector from N_{ph} to 2 (corresponding to one α - β stator plane). The different cases of the transformed stator current vector $\underline{I}_{s\alpha\beta}$, which depend on the imposed stator sequence, are shown in Table 1.

This table shows the projection of four different cases of sequences (u_0, u_1, u_2, u_3), where:

- $u_0 = N_{ph}$ (Zero-sequence) \rightarrow Example of 6 phases: $u_0 = 5$;
- $u_1 < \text{floor}\left(\frac{N_{ph}}{2}\right)$ (\leq if N_{ph} is odd) \rightarrow Example of 6 phases: $u_1 = 1$ or 2;
- $u_2 > \text{floor}\left(\frac{N_{ph}}{2}\right)$ \rightarrow Example of 6 phases: $u_2 = 4$ or 5;
- $u'_0 = \frac{N_{ph}}{2}$ (considered as a "second zero-sequence", exists only when N_{ph} is even) \rightarrow Example of 6 phases: $u'_0 = 3$.

Table 1. Transformed stator current vector $\underline{I}_{s_{\alpha\beta}}$.

Index	Transformed Current vector $\underline{I}_{s_{\alpha\beta}}$
1	$\sqrt{N_{ph}} \cdot I_{s_{u_0}} \cdot \sin(2\pi f_{s_{u_0}} \cdot t + \varphi_{s_{u_0}})$
\vdots	\vdots
$2 \cdot u_1$	$\sqrt{\frac{N_{ph}}{2}} \cdot I_{s_{u_1}} \cdot \sin(2\pi f_{s_{u_1}} \cdot t + \varphi_{s_{u_1}})$
$2 \cdot u_1 + 1$	$-\sqrt{\frac{N_{ph}}{2}} \cdot I_{s_{u_1}} \cdot \cos(2\pi f_{s_{u_1}} \cdot t + \varphi_{s_{u_1}})$
\vdots	\vdots
$2 \cdot (N_{ph} - u_2)$	$\sqrt{\frac{N_{ph}}{2}} \cdot I_{s_{u_2}} \cdot \sin(2\pi f_{s_{u_2}} \cdot t + \varphi_{s_{u_2}})$
$2 \cdot (N_{ph} - u_2) + 1$	$\sqrt{\frac{N_{ph}}{2}} \cdot I_{s_{u_2}} \cdot \cos(2\pi f_{s_{u_2}} \cdot t + \varphi_{s_{u_2}})$
\vdots	\vdots
N_{ph}	$\sqrt{N_{ph}} \cdot I_{s_{u'_0}} \cdot \sin(2\pi f_{s_{u'_0}} \cdot t + \varphi_{s_{u'_0}})$

3.2. α - β -0 Transformation of Mutual Inductance Matrix

Stator to rotor mutual inductance matrix $L_{sr}(\theta)$ is characterized by the presence of several space harmonics, α - β -0 transformation allows to separate these space harmonics into different planes.

The transformed matrix $\underline{L}_{sr_{\alpha\beta}}(\theta)$ has a dimension of $N_{ph} \times N_{bar}$, its general form is presented in Figure 2.

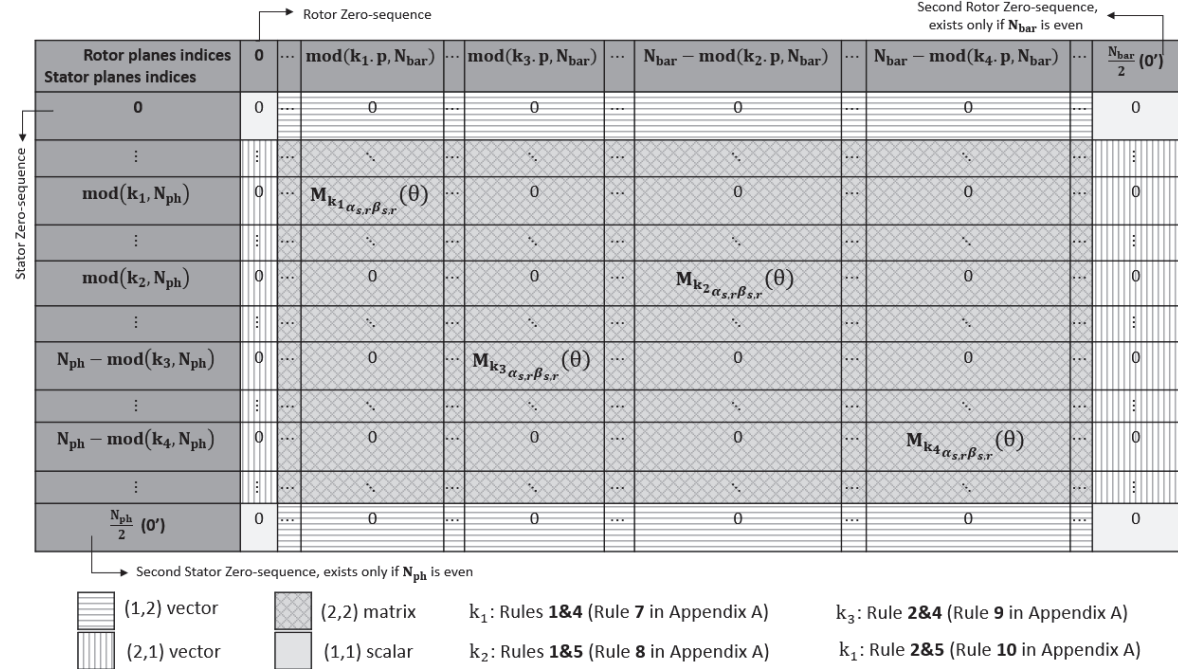


Figure 2. Transformed Stator-Rotor mutual inductance matrix $\underline{L}_{sr_{\alpha\beta}}(\theta)$.

As we can see in Figure 2, each space harmonic appears only in one cell. The different cells in Figure correspond to two-by-two matrices, two-element vectors or scalars (as described in the legend in the bottom of Figure). A cell corresponding to two-by-two matrix represents the intersection between

an α - β stator plane and an α - β rotor plane. The general form of these matrices can be written as follows:

$$\underline{M}_{k\alpha_{s,r}\beta_{s,r}}(\theta) = \frac{\sqrt{N_{ph} \cdot N_{bar}}}{2} \cdot M_{srk} \begin{pmatrix} \delta_{11} \cos(-kp\theta + \varphi_{srk}) & \delta_{12} \sin(-kp\theta + \varphi_{srk}) \\ \delta_{21} \sin(-kp\theta + \varphi_{srk}) & \delta_{22} \cos(-kp\theta + \varphi_{srk}) \end{pmatrix}, \quad (18)$$

where $\delta_{ij} = \pm 1$ depending on the category of the harmonic (see the rules in Appendix A).

The following rules (from 1 to 6) define the distribution of a space harmonic “k” in $\underline{L}_{sr\alpha\beta}(\theta)$:

- Rule 1: If $\text{mod}(k, N_{ph}) \leq \frac{N_{ph}-1}{2}$, the harmonic “k” is projected on the stator plane number $\text{mod}(k, N_{ph})$, so the two lines whose indices are $2 \cdot \text{mod}(k, N_{ph})$ and $2 \cdot \text{mod}(k, N_{ph}) + 1$ in the matrix $\underline{L}_{sr\alpha\beta}(\theta)$.
- Rule 2: If $\text{mod}(k, N_{ph}) > \frac{N_{ph}-1}{2}$, the harmonic “k” is projected on the stator plane number $N_{ph} - \text{mod}(k, N_{ph})$, so the two lines whose indices are $2 \cdot (N_{ph} - \text{mod}(k, N_{ph}))$ and $2 \cdot (N_{ph} - \text{mod}(k, N_{ph})) + 1$ in the matrix $\underline{L}_{sr\alpha\beta}(\theta)$.
- Rule 3: If $\text{mod}(k, N_{ph}) = 0$, the harmonic “k” is projected on the first stator zero-sequence (0), so the first line in the matrix $\underline{L}_{sr\alpha\beta}(\theta)$.
- Rule 3-bis (Only when N_{ph} is even): If $\text{mod}(k, N_{ph}) = \frac{N_{ph}}{2}$, the harmonic “k” is projected on the second stator zero-sequence, so the last line in the matrix $\underline{L}_{sr\alpha\beta}(\theta)$.
- Rule 4: If $\text{mod}(k.p, N_{bar}) \leq \frac{N_{bar}-1}{2}$, the harmonic “k” is projected on the rotor plane number $\text{mod}(k.p, N_{bar})$, so the two lines whose indices are $2 \cdot \text{mod}(k.p, N_{bar})$ and $2 \cdot \text{mod}(k.p, N_{bar}) + 1$ in the matrix $\underline{L}_{sr\alpha\beta}(\theta)$.
- Rule 5: If $\text{mod}(k.p, N_{bar}) > \frac{N_{bar}-1}{2}$, the harmonic “k” is projected on the rotor plane number $N_{bar} - \text{mod}(k.p, N_{bar})$, so the two lines whose indices are $2 \cdot (N_{bar} - \text{mod}(k.p, N_{bar}))$ and $2 \cdot (N_{bar} - \text{mod}(k.p, N_{bar})) + 1$ in the matrix $\underline{L}_{sr\alpha\beta}(\theta)$.
- Rule 6: If $\text{mod}(k.p, N_{bar}) = 0$, the harmonic “k” is projected on the first rotor zero-sequence (0), so the column whose index is 1, in the matrix $\underline{L}_{sr\alpha\beta}(\theta)$.
- Rule 6-bis (Only when N_{bar} is even): If $\text{mod}(k.p, N_{bar}) = \frac{N_{bar}}{2}$, the harmonic “k” is projected on the second rotor zero-sequence (0'), so the last column in the matrix $\underline{L}_{sr\alpha\beta}(\theta)$.

When several harmonics are projected on the same stator and rotor planes, their matrices (or vectors or scalars, see Appendix A) $\underline{M}_{k\alpha_{s,r}\beta_{s,r}}(\theta)$ are superposed. For example, we consider the harmonics “k₁” and “k₂”, if: $\text{mod}(k_1, N_{ph}) = \text{mod}(k_2, N_{ph})$ and $\text{mod}(k_1.p, N_{bar}) = \text{mod}(k_2.p, N_{bar})$, then $\underline{M}_{k_1\alpha_{s,r}\beta_{s,r}}(\theta)$ and $\underline{M}_{k_2\alpha_{s,r}\beta_{s,r}}(\theta)$ are superposed in $\underline{L}_{sr\alpha\beta}(\theta)$.

3.3. α - β -0 Transformation of Rotor Current Vector

For the windings with an integral number of slots per phase and per pole-pair, **2.spp = integral** (including the special fractional like spp = 0.5, 1.5 ...), the sequence “u” generates a set of space harmonics listed in F_u , which is defined as follows [30]:

$$F_u = \{Z \cdot N_{ph} + u\} p \quad Z = 0, 1, -1, 2, -2, 3, -3 \dots \quad (19)$$

It should be mentioned that all the space harmonics are multiples of “p”.

As examples:

- For $N_{ph} = 3$ and “u = 1”, $F_u = p \cdot \{1, -2, 4, -5, 7, -8, 10, -11, 13 \dots\}$;

- For $N_{ph} = 5$: for "u = 1", $F_u = p * \{1, -4, 6, -9, 11, -14, 16, -19, 21 \dots\}$; for u = 2, $F_u = p * \{2, -3, 7, -8, 12, -13, 17 \dots\}$; for u = 3, $F_u = p * \{-2, 3, -7, 8, -12, 13, -17 \dots\}$. As we can see, the sequences "u = 2" and "u = 3" generate the same set of space harmonics, but with opposite signs.

If the stator sequence "u" has a frequency " f_s ", each space harmonic "v.p" ("v" is integer, can be positive or negative) belonging to the set " F_u " induces in the rotor bar currents a frequency, called " f_{rv} ", which can be determined as follows (taking into account the mechanical speed of the rotor):

$$f_{rv} = f_s - v.p. \frac{\Omega_{mec}}{2\pi} \quad (20)$$

So, the current in a rotor bar "i", can be expressed as following:

$$I_{r_i} = \sum_v \left[I_{r_v} \cdot \sin \left(2\pi \cdot f_{r_v} \cdot t - (i-1) \cdot v.p. \frac{2\pi}{N_{bar}} + \varphi'_{r_v} \right) \right] \quad (21)$$

where I_{r_v} is the amplitude of the harmonic range "v" of rotor current, and φ'_{r_v} is its angle.

The transformed rotor current vector contains several harmonics (according to Equation (19)), separated into different α - β rotor planes as shown in the Table 2. Each α - β rotor plane can be excited by more than one harmonic, but to simplify the presentation only one harmonic per plane is considered in the table.

Table 2. Transformed rotor current vector $\underline{I}_{r_{\alpha\beta}}$.

Index	Transformed Current vector $\underline{I}_{r_{\alpha\beta}}$
1	$\sqrt{N_{bar}} \cdot I_{r_{v_0}} \cdot \sin(2\pi f_{r_{v_0}} \cdot t + \varphi_{r_{v_0}})$
\vdots	\vdots
$2 \cdot \text{mod}(v_1 \cdot p, N_{bar})$	$\sqrt{\frac{N_{bar}}{2}} \cdot I_{r_{v_1}} \cdot \sin(2\pi f_{r_{v_1}} \cdot t + \varphi_{r_{v_1}})$
$2 \cdot \text{mod}(v_1 \cdot p, N_{bar}) + 1$	$-\sqrt{\frac{N_{bar}}{2}} \cdot I_{r_{v_1}} \cdot \cos(2\pi f_{r_{v_1}} \cdot t + \varphi_{r_{v_1}})$
\vdots	\vdots
$2 \cdot \text{mod}(v_2 \cdot p, N_{bar})$	$\sqrt{\frac{N_{bar}}{2}} \cdot I_{r_{v_2}} \cdot \sin(2\pi f_{r_{v_2}} \cdot t + \varphi_{r_{v_2}})$
$2 \cdot \text{mod}(v_2 \cdot p, N_{bar}) + 1$	$\sqrt{\frac{N_{bar}}{2}} \cdot I_{r_{v_2}} \cdot \cos(2\pi f_{r_{v_2}} \cdot t + \varphi_{r_{v_2}})$
\vdots	\vdots
$2 \cdot (N_{bar} - \text{mod}(v_3 \cdot p, N_{bar}))$	$\sqrt{\frac{N_{bar}}{2}} \cdot I_{r_{v_3}} \cdot \sin(2\pi f_{r_{v_3}} \cdot t + \varphi_{r_{v_3}})$
$2 \cdot (N_{bar} - \text{mod}(v_3 \cdot p, N_{bar})) + 1$	$\sqrt{\frac{N_{bar}}{2}} \cdot I_{r_{v_3}} \cdot \cos(2\pi f_{r_{v_3}} \cdot t + \varphi_{r_{v_3}})$
\vdots	\vdots
$2 \cdot (N_{bar} - \text{mod}(v_4 \cdot p, N_{bar}))$	$\sqrt{\frac{N_{bar}}{2}} \cdot I_{r_{v_4}} \cdot \sin(2\pi f_{r_{v_4}} \cdot t + \varphi_{r_{v_4}})$
$2 \cdot (N_{bar} - \text{mod}(v_4 \cdot p, N_{bar})) + 1$	$-\sqrt{\frac{N_{bar}}{2}} \cdot I_{r_{v_4}} \cdot \cos(2\pi f_{r_{v_4}} \cdot t + \varphi_{r_{v_4}})$
\vdots	\vdots
N_{bar}	$\sqrt{N_{bar}} \cdot I_{r_{v_5}} \cdot \sin(2\pi f_{r_{v_5}} \cdot t + \varphi_{r_{v_5}})$

This table shows the projection of four different cases of rotor harmonics ($v_0, v_1, v_2, v_3, v_4, v_5$), where:

- $\text{mod}(|v_0| \cdot p, N_{bar}) = N_{bar}$ (Rotor zero-sequence);
- $\text{mod}(|v_1| \cdot p, N_{bar}) < \text{floor}\left(\frac{N_{bar}}{2}\right)$ (\leq if N_{bar} is odd), and $v_1 > 0$;
- $\text{mod}(|v_2| \cdot p, N_{bar}) < \text{floor}\left(\frac{N_{bar}}{2}\right)$ (\leq if N_{bar} is odd), and $v_2 < 0$;
- $\text{mod}(|v_3| \cdot p, N_{bar}) > \text{floor}\left(\frac{N_{bar}}{2}\right)$, and $v_3 > 0$;

- $\text{mod}(|v_4| \cdot p, N_{\text{bar}}) > \text{floor}\left(\frac{N_{\text{bar}}}{2}\right)$, and $v_4 < 0$;
- $\text{mod}(|v_5| \cdot p, N_{\text{bar}}) = \frac{N_{\text{bar}}}{2}$ (only when N_{bar} is even).

When different harmonics are projected on the same α - β rotor plane, their α and β components are superposed. This property will be at the origin of torque ripples.

As we can see in the Table 2, the harmonic ranges “ v ” and their frequencies “ f_{r_v} ”, are written in the absolute value. In fact, according to the Equation (19), “ v ” can be positive or negative. Therefore, the frequencies “ f_{r_v} ”, determined by the Equation (20), can also have positive or negative values.

In the case when “ f_{r_v} ” is positive, the angle “ φ_{r_v} ” in Concordia base is equal to “ φ'_{r_v} ” in the natural base. When “ f_{r_v} ” is negative, $\varphi_{r_v} = \pi - \varphi'_{r_v}$.

4. Torque Calculation

The torque can be calculated in the natural base by the following expression [27]:

$$T = \underline{I}'_s \cdot \frac{dL_{sr}(\theta)}{d\theta} \cdot \underline{I}_r \quad (22)$$

In the case when the stator is supplied by a sequence “ u ”, the torque can be expressed in α - β -0 base as following:

$$T = \begin{bmatrix} I_{s\alpha} & I_{s\beta} \end{bmatrix} \cdot \sum_v \frac{dM_{\alpha\beta v}}{d\theta} \cdot \underline{I}_{r_{\alpha\beta v}} \quad (23)$$

Using the transformed stator and rotor current vectors (Tables 1 and 2), and the transformed mutual inductance matrix (Figure 2), the expressions of the torque developed from the interaction between space and time harmonics can be determined.

As described before, under stator sequence “ u ”, a set of space harmonics “ F_u ” are excited (defined in Equation (19)). These space harmonics induce in the rotor bars currents time harmonics. Each rotor time harmonic is characterized by its amplitude “ I_{r_v} ”, its frequency “ f_{r_v} ”, its polarity “ $v \cdot p$ ” and its phase “ φ_{r_v} ” (defined in Concordia base).

A space harmonic “ $v \cdot p$ ” (represented by its projection on Concordia base $M_{v\alpha s, r\beta s, r}(\theta)$) interacts with the rotor time harmonic of the same range to produce a constant torque (as shown in Equations (24), (25) and (26)).

The torque developed from the interaction the space and the rotor time harmonic “ $v = u$ ” is expressed as follows:

$$T_{u-u} = \frac{N_{\text{ph}} \cdot N_{\text{bar}}}{4} \cdot p \cdot u \cdot I_{s_u} \cdot I_{r_{v=u}} \cdot (M_{sr})_{v=u} \cdot \sin(\varphi_{s_u} - \varphi_{r_{v=u}} + \varphi_{sr_{v=u}}) \quad (24)$$

For the other induced rotor harmonics, the resultant harmonic constant torque depends on the sign of “ v ”:

- if $v < 0$:

$$T_{v^- - v^-} = -\frac{N_{\text{ph}} \cdot N_{\text{bar}}}{4} \cdot p \cdot |v| \cdot I_{s_u} \cdot I_{r_v} \cdot (M_{sr})_v \cdot \sin(\varphi_{s_u} - \varphi_{r_v} - \varphi_{sr_v}) \quad (25)$$

- if $v > 0$:

$$T_{v^+ - v^+} = -\frac{N_{\text{ph}} \cdot N_{\text{bar}}}{4} \cdot p \cdot |v| \cdot I_{s_u} \cdot I_{r_v} \cdot (M_{sr})_v \cdot \sin(\varphi_{s_u} + \varphi_{r_v} + \varphi_{sr_v}) \quad (26)$$

On the other hand, when many space and rotor time harmonics are superposed on the same α - β rotor plane, they all interact with each other. For example, if two harmonic ranges “ v_1 ” and “ v_2 ” are in the same rotor Concordia plane, the space harmonic “ $v_1 \cdot p$ ” interacts with the time harmonics “ $v_2 \cdot p$ ”

and produces pulsating torque component. The expressions of pulsating torque components depend on the harmonic ranges, and are detailed in Appendix B.

5. Application

To validate the mathematical approach of prediction of pulsating torque components proposed in this paper, several multiphase induction machine topologies were inspected. Transient F-E simulations (Maxwell 2D software) are done to determine the steady-state torque temporal variation. The simulations are done by imposing sinusoidal stator currents.

5.1. Three-Phase Induction Machine

The machine parameters are shown in Appendix C.

Table 3 shows the imposed parameters in this simulation.

Table 3. Simulation parameters for the three-phase induction machine (IM).

Parameter	Value
Injected sequence	1
Stator current frequency	50 Hz
Peak current	80 A
Mechanical speed	1470 rpm

The simulated machine has a 36 stator slots, two investigated numbers of rotor bars 48 and 49, and two pole-pairs. The torque variation is extracted from the simulation results (for the machine with 48 bars), and a harmonic analysis is done and shown in Figure 3.

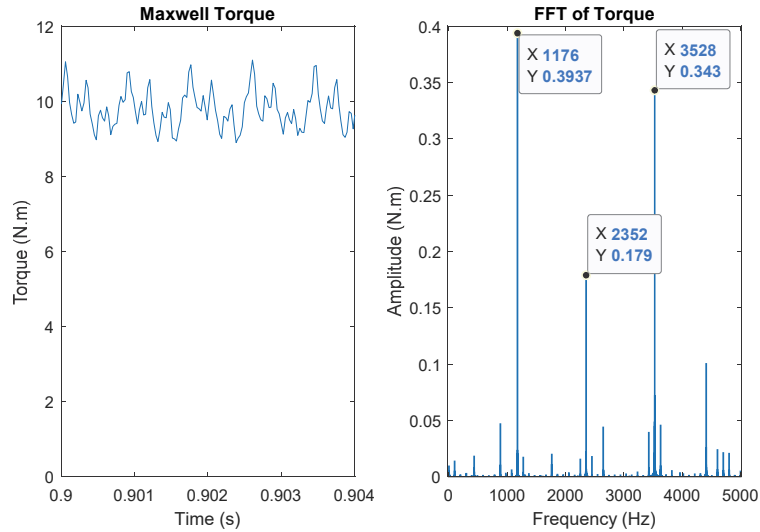


Figure 3. FFT of Finite-Element torque for the three-phase IM with 48 bars.

As we can see in Figure 3, there are three predominant pulsating frequencies (1176 Hz, 2352 Hz and 3528 Hz). To understand the origin of these frequencies, it is important to determine the interactions between space and time harmonics. The distribution of space harmonics (up to $v = 50$) in the matrix $L_{sr_{\alpha\beta}}(\theta)$, according to Figure 2, is shown in Figure 4. To simplify the presentation, each space harmonic “ $v.p$ ” is represented by the number “ v ” in Figure ($v = 1, \dots, 50$).

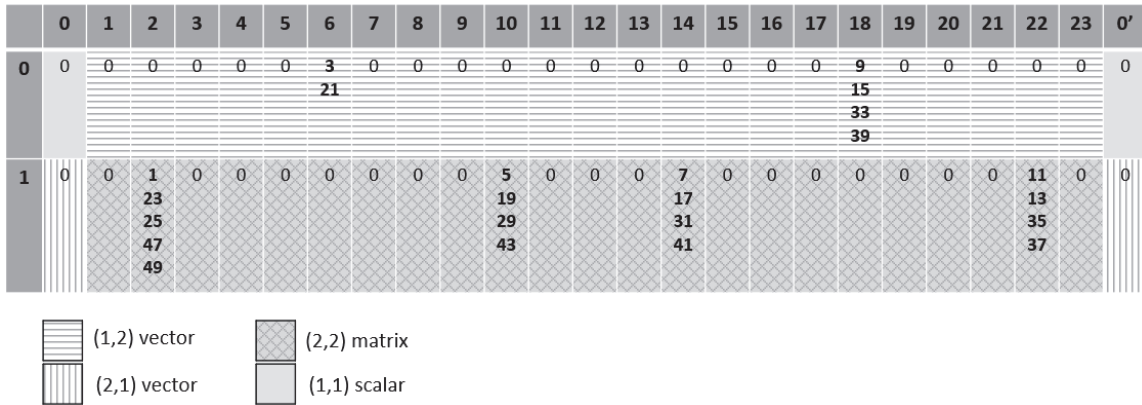


Figure 4. Space harmonics distribution in $L_{sr,\alpha\beta}(\theta)$ for the three-phase IM with 48 bars (in line 1, 4 planes among 23 are excited, each one by several harmonics which lead to important interactions).

All the even order harmonics are null, according to the winding distribution, so only odd order harmonics appear in the matrix.

The first line (0) in the table corresponds to the stator zero-sequence. For this zero-sequence, the harmonics 3 and 21 can excite the plane n°6 of the rotor and the harmonics 9, 15, 33 and 39 can excite the plane n°18 of the rotor. With a wye-connected three-phase machine, the zero-sequence is not excited. As consequence the plane n°6 and n°18 are not excited.

The second line (1) corresponds to the only possible stator plane in the case of this three-phase machine, which can be excited by sequences “u = 1” or “u = 2”. For this stator plane, **four** different rotor planes are excited (n°2, n°10, n°14, n°22), each one of them contains several superposed harmonics. As example the rotor plane n°2 is excited by the space harmonics n°1, 23, 25, 47 and 49.

In Figure 4, the number of bars is even, so a second rotor zero-sequence appears in **column 0'**.

According to the Equations (20), (27) and Table A1 in Appendix B, the frequencies of pulsating torque components can be predicted (shown in Table 4).

Table 4. Predicted torque pulsating components frequencies for the three-phase IM with 48 bars (in Hz).

Rotor Plane	2					10				
	1p	23p	25p	47p	49p	5p	19p	29p	43p	
1p	0	1176	1176	2352	2352	5p	0	1176	1176	2352
23p	1176	0	2352	3528	3528	19p	1176	0	2352	1176
25p	1176	2352	0	3528	1176	29p	1176	2352	0	3528
47p	2352	3528	3528	0	4704	43p	2352	1176	3528	0
49p	2352	3528	1176	4704	0					
Rotor Plane	14				22					
	7p	17p	31p	41p	11p	13p	35p	37p		
7p	0	1176	1176	2352	11p	0	1176	1176	2352	
17p	1176	0	2352	1176	13p	1176	0	2352	1176	
31p	1176	2352	0	3528	35p	1176	2352	0	3528	
41p	2352	1176	3528	0	37p	2352	1176	3528	0	

In this table, we find the same three predominant frequencies in Figure 3, which are 1176, 2352 and 3528 Hz. Each frequency results from many interactions between space and time harmonics. For example, the frequency 1176 Hz (the highest torque pulsation) results from the interactions: 1p-23p, 1p-25p, 25p-49p, 5p-19p, 5p-29p, 19p-43p, 7p-17p, 7p-31p, 17p-41p, 11p-13p, 11p-35p and 13p-37p.

According to the Table A3 in Appendix C, the 10 most important space harmonics (excited by the sequence “u = 1”) regarding their amplitudes in the winding function are: v = 1, 5, 7, 11, 13, 17, 19, 23, 35, 37. An interaction between two of these harmonics is considered important regarding the amplitude of torque pulsation. These significant interactions are colored in gray in the Table 4.

We observe that most of these colored interactions generate a pulsation at 1176 Hz, which is consistent with F-E results (Figure 3), where the highest amplitude of torque ripple is at 1176 Hz.

Basing on Figure 3, this machine with 48 bars, presents many superpositions between space harmonics (in the same rotor planes). All these harmonics interactions can be avoided by taking 49 bars instead of 48. Figure 5 shows how the choice of 49 bars improves the separation between harmonics, all the harmonics below $v = 50$ are perfectly separated.

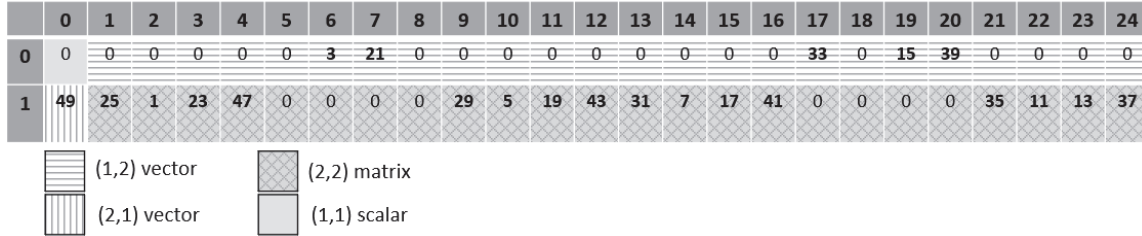


Figure 5. Space harmonics distribution in $L_{sr\alpha\beta}(\theta)$ for the three-phase IM with 49 bars (in line 1, 17 planes among 24 are excited, each one by one harmonic, so without interactions).

The Figure 6 shows the comparison between the three-phase IM with 48 and 49 bars regarding the torque (F-E results). It proves that the torque is smoother with 49 bars thanks to the good separation between space harmonics into several α - β rotor planes (as shown in Figure 5).

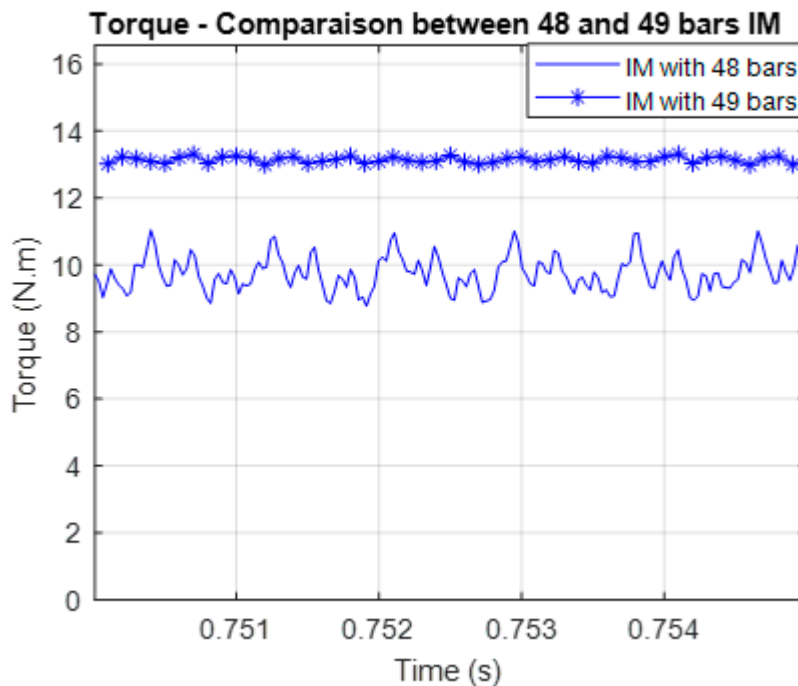


Figure 6. Comparison of the developed torque with 48 and 49 bars, three-phase IM.

This result is consistent with the rules related to the choice of rotor bars number described in the reference [2] (pp. 340–341). According to this reference, for this machine topology (three phases, 36 stator slots and two pole-pairs), the number of bars 48 presents harmful parasitic torque, which is explained thanks to this new approach of prediction of torque pulsations.

After this first example, we conclude that even for three-phase induction machines, the classical approach of rotor cage modeling (simply by one α - β rotor plane) can only be used if the combination of phase number and bar number does not lead to harmful interactions between space and time harmonics (like the combination $N_{ph} = 3$, $N_{bar} = 49$ and $p_1 = 2$). However, when the combination causes harmful

interactions (like the one with 48 bars in this example), it is important to consider several α - β rotor planes in modeling (4 planes in this example, for 48 bars) to predict the parasitic torque pulsations.

5.2. Five-Phase Induction Machine with Double-Layer Tooth Concentrated Winding

The machine parameters are shown in Appendix D. The Table 5 shows simulation parameters.

Table 5. Simulation parameters for the five-phase IM with tooth concentrated winding.

Parameter	Value
Injected sequences	1, 3
Stator current frequencies	50, 150 Hz
Peak current	50 A _{peak}
Mechanical speed	744, 740.6 rpm

Two different stator sequences are imposed in separate simulations ($u = 1$ and $u = 3$). According to the Equation (19), the sequence “ $u = 1$ ” induces in the rotor bars the set of harmonics “ $v = 1, -4, 6, -9, 11, -14, 16, -19 \dots$ ”. The sequence “ $u = 3$ ” induces the set of rotor current harmonics “ $v = -2, 3, -7, 8, -12, 13, -17, 18 \dots$ ”.

This machine was studied in a previous paper [22]. In this paper it has been observed that the number of bars 64 presents important torque pulsations, especially referring to the number 65 characterized by a very low torque harmonics content. Nevertheless, according to the rules in the reference [2] (pp. 340–341), for this machine topology (5 phases, 20 stator slots and 4 pole-pairs) the number of bars 64 is not forbidden. In fact, the reported rules for selecting the rotor bars number are defined for three-phase machines and are not directly applicable to multiphase machines.

To explain the origin of torque pulsations for this machine with 64 bars, the developed torque under both sequences is extracted from the simulation results, and a harmonic analysis is done and shown in Figures 7 and 9.

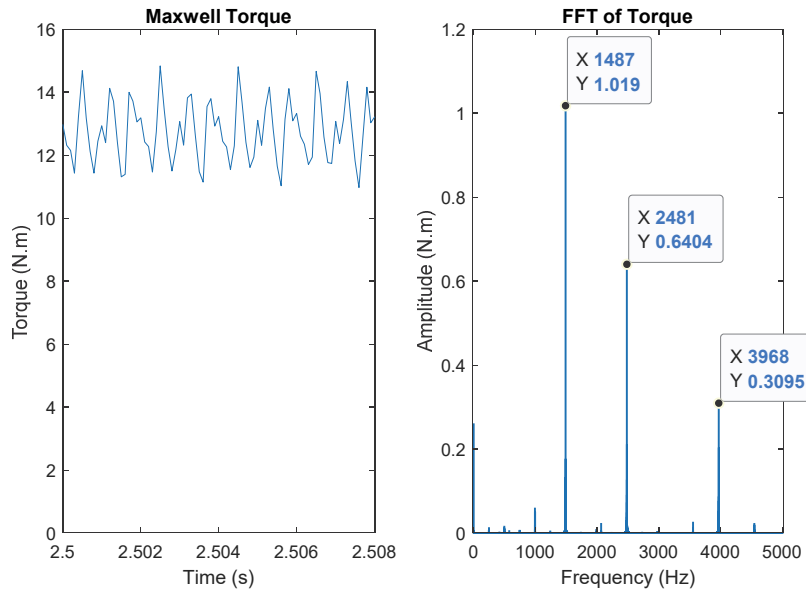


Figure 7. FFT of Finite-Element torque for the five-phase IM with 64 bars. Under sequence “ $u = 1$ ” (the planes of line (1), Figure 8 are excited).

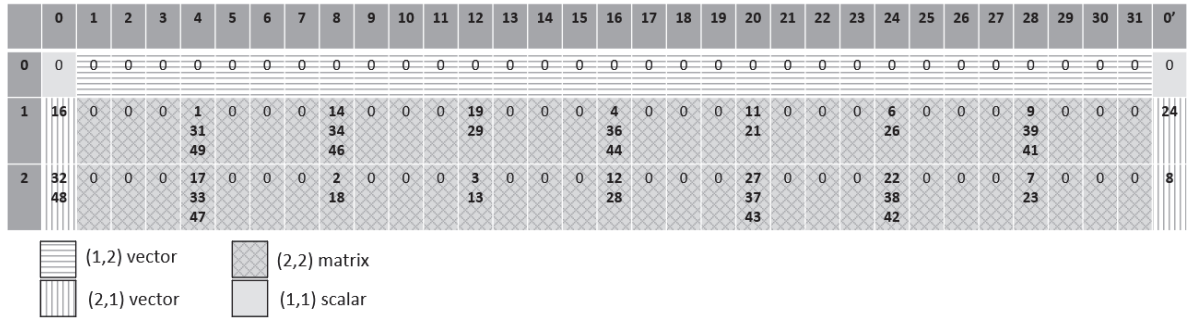


Figure 8. Space harmonics distribution in $L_{sr\alpha\beta}(\theta)$ for the five-phase IM with 64 bars (in each line, 7 rotor planes, among 31, and both zero-sequences are excited, which lead to harmful interactions).

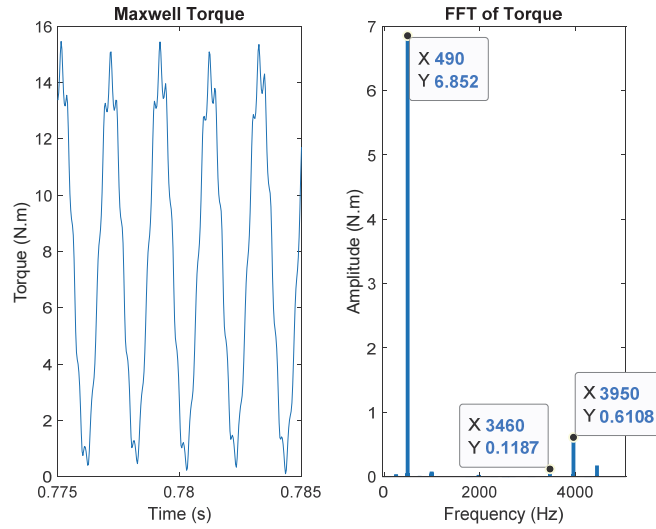


Figure 9. FFT of Finite-Element torque for the five-phase IM with 64 bars. Under sequence “ $u = 3$ ” (the planes of line (2), Figure 8 are excited).

With this number of bars, the developed torque under both sequences contains important pulsating components. To understand the origin of these components, the distribution of space harmonics (up to $v = 50$) in the matrix $L_{sr\alpha\beta}(\theta)$, according to the Figure 2, is shown in the Figure 8.

This figure shows the interactions between space and time harmonics.

The space harmonics generated by the sequence “ $u = 1$ ” are distributed in the second line of Figure, and the sequence “ $u = 3$ ” corresponds to the third line. Under both stator sequences, **seven** rotor planes, the **first rotor zero-sequence** (column 0) and the **second rotor zero-sequence** (column 0') are excited, each one of them contains several superposed harmonics.

According to the Appendix B, the frequencies of pulsating torque components can be predicted for the sequence “ $u = 1$ ” in Table 6 and for the sequence “ $u = 3$ ” in Table 7. According to the Table 6, the frequencies of torque pulsations for the sequence “ $u = 1$ ” (space harmonics considered up to “ $v = 50$ ”) are: 1487, 2481 and 3968 Hz. These frequencies are the same as the predominant frequencies in F-E results (Figure 7). The interactions in the Table 7 produce three main frequencies: 490, 3460 and 3950 Hz, the same as F-E results in Figure 9.

Table 6. Predicted torque pulsating components frequencies for the five-phase IM with 64 bars (in Hz). Sequence “u = 1”.

Rotor Plane	0 (First Rotor Zero-Sequence)			4			8				
		16p			1p	31p	49p		14p	34p	46p
16p	0 & 1487			1p	0	1487	2481	14p	0	2481	1487
				31p	1487	0	3968	34p	2481	0	3968
				49p	2481	3968	0	46p	1487	3968	0
Rotor Plane	12			16			20				
	19p	29p			4p	36p	44p		11p	21p	
	0	2481		4p	0	1487	2481	11p	0	1487	
	2481	0		36p	1487	0	3968	21p	1487	0	
				44p	2481	3968	0				
Rotor Plane	24			28			0' (Second Rotor Zero-Sequence)				
	6p	26p			9p	39p	41p		24p		
6p	0	1487		9p	0	2481	1487				
26p	1487	0		39p	2481	0	3968	24p	0 & 2481		
				41p	1487	3968	0				

Table 7. Predicted torque pulsating components frequencies for the five-phase IM with 64 bars (in Hz). Sequence “u = 3”.

Rotor Plane	0 (First Rotor Zero-Sequence)			4			8				
		32p			17p	33p	47p		2p	18p	
	0 & 3460	3950 & 490		17p	0	490	3460	2p	0	490	
				33p	490	0	3950	18p	490	0	
	3950 & 490	0 & 4440		47p	3460	3950	0				
Rotor Plane	12			16			20				
	3p	13p			12p	28p		27p	37p	43p	
	0	490		12p	0	490		27p	0	3460	490
	490	0		28p	490	0		37p	3460	0	3950
								43p	490	3950	0
Rotor Plane	24			28			0' (Second Rotor Zero-Sequence)				
	22p	38p	42p		7p	23p		8p			
6p	0	490	3460	7p	0	490		0 & 490			
26p	490	0	3950	23p	490	0	8p				
	3460	3950	0								


In Appendix D, the 20 most important space harmonics regarding amplitude are indicated in the Table A4. To have an idea about the importance of a torque pulsation, every interaction between two harmonics among the 20 most important ones (Table A4) is colored in gray.


Under the sequence “u = 1”, the only interaction between a space and a time harmonic whose amplitudes are among the 20 most important, is in the cell corresponding to the rotor zero-sequence containing the harmonic “16.p”, whose interaction generates two components (as explained in Appendix B): a constant torque at 0 Hz, and a pulsating component of a frequency of 1487 Hz (which has the highest pulsation amplitude in Figure 7).


Under the sequence “u = 3”, several interactions are colored in gray (interaction between two harmonics among the 20 most important in Table A4). All these marked interactions generate a frequency of 490 Hz, which has the highest pulsation amplitude in Figure 9.

The Figure 10 shows the space harmonics distribution in $L_{sr\alpha\beta}(\theta)$ for the machine with 65 bars. We observe that this number of bars allows to separate the harmonics better than the case of 64 bars. In fact, the harmonics with the most important amplitudes (Table A4) do not interact between each other. Furthermore, with this number of bars, the rotor zero-sequence is not excited under both sequences.

	0	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	25	26	27	28	29	30	31	32	
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
1	0	16 49	0	0	1	0	31 34	0	0	14	0	19 46	0	0	29 36	0	4	0	0	21 44	0	11	0	0	6	0	26 39	0	0	9	0	24 41	0	0
2	0	0	32 33	17 48	0	0	0	18 47	2	0	0	0	3	13	0	0	0	12	28 37	0	0	0	27 38	22 43	0	0	0	23 42	7	0	0	0	0	8

 (1,2) vector

 (2,2) matrix

 (2,1) vector


 (1,1) scalar

Figure 10. Space harmonics distribution in $L_{sr\alpha\beta}(\theta)$ for the five-phase IM with 65 bars (in line 1, 13 planes among 32 are excited. In line 2, 12 planes among 32 are excited. No zero-sequence is excited. This leads to less interactions than 64 bars).

The Figure 11 shows the comparison of the developed torque between the machines with 64 and 65 bars. With 65 bars the torque is smoother under both sequences thanks to the separation of harmonics (shown in Figure 10). FFT analysis of torque for the new number of bars “65”, is done and shown in Figure 12 for both sequences.

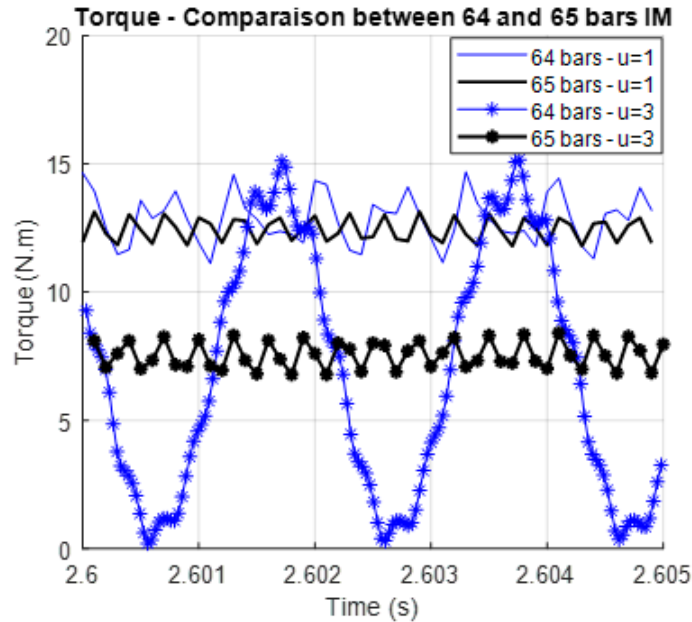


Figure 11. Comparison of the developed torque with 64 and 65 bars, five-phase IM.

Under the sequence “ $u = 1$ ” (corresponds to the black curve in Figure 11), the predominant pulsating component is at the frequency of 3224 Hz. After the application of equations in Appendix B, it appears that this frequency of 3224 Hz is generated by all the interactions shown in Figure 10 (Line 1), so the interactions: 16p-49p, 31p-34p, 19p-49p, 29p-36p, 21p-44p, 26p-39p and 24p-41p (Considering only 50 space harmonics, from 1p to 50p). Comparing to the case of 64 bars (Table 6), the interactions in the machine with 65 bars are mostly between harmonics with relatively low amplitudes (Table A4) which explains the reduction of torque ripple thanks to this combination ($N_{ph} = 5$, $N_{bar} = 65$ and $p_1 = 4$).

Under the sequence “ $u = 3$ ” (corresponds to the blue curve in Figure 11), the predominant pulsating component is at the frequency of 3211 Hz. Using the same approach (Appendix B), we determine that this frequency is generated by all the interactions in the last line of Figure 10, so: 32p-33p, 17p-48p, 18p-47p, 28p-37p, 27p-38p, 22p-43p, 23p-42p. These harmonics have also relatively low amplitudes (Table A4).

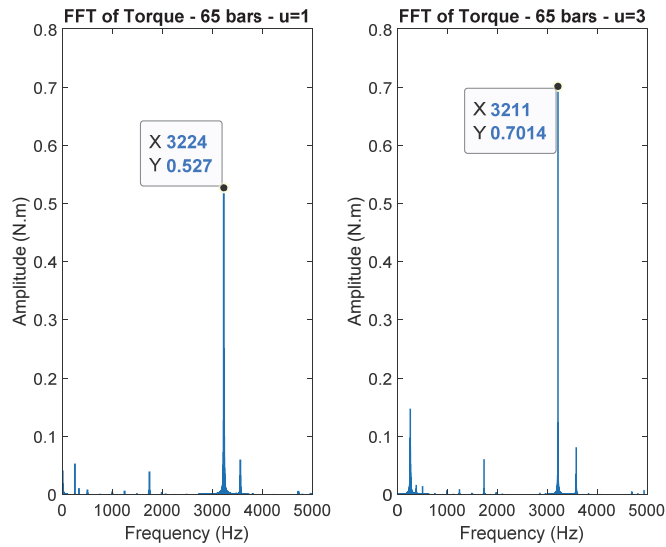


Figure 12. FFT analysis of torque curves for the IM with 65 bars, under sequences “ $u = 1$ ” and “ $u = 3$ ”.

For five-phase induction machine, two α - β stator planes must be considered in modeling. The number of α - β rotor planes to be considered depends on the number of rotor bars. If there are interactions between harmonics with important amplitudes, the rotor planes containing these interactions must be considered to predict torque pulsations.

6. Conclusion

In order to be able to design multiphase induction drives which can be supplied with different harmonics of currents either for speed extension or for torque increase, the paper provides a mathematical approach to point out the excited α - β rotor planes which must be considered at two levels for the modelling. At first, during the electromagnetic design, the aim is to avoid wrong choices of bar numbers leading to torque pulsations. Secondly during the design of the vector control of the machine, it is necessary to use the simplest modeling of the rotor, with the minimum dimension, in order to define the vector control. The selection of the minimum number of rotor planes which contribute to the torque leads to a reduced-order modelling of the rotor cage. It has been shown that the rotor dimension, which is “ N_{bar} ” in the natural base model, can be much lower in the Concordia frame.

In this study, a general arithmetic tool (based on equations in Appendix B) was developed to predict, and consequently to avoid, torque ripples due to wrong combinations of phase number and rotor bar number. The result of the analysis is presented graphically in simple table where it is appearing clearly when interactions between space and time harmonics occur. This tool requires as inputs: the numbers of stator phases, rotor bars, pole-pairs (p_1), the stator frequency, the stator sequence “ u ” and the mechanical speed. It should be mentioned that this tool does not predict the pulsations due to slotting effect.

This arithmetic tool is validated on two different multiphase induction machines topologies (three and five phases) with different numbers of rotor bars. In the provided applications, the induced pulsating torque frequencies are predicted thanks to the arithmetic tool and validated by the FFT analysis of the steady-state torque determined by F-E simulations (Maxwell 2D software, with imposed sinusoidal stator currents).

Through the first application, it has been concluded that even for three-phase induction machines the simplified modeling of the rotor cage by only one equivalent two-phase circuit is not always relevant. Some numbers of bars lead to important interactions between space and time harmonics, and consequently important torque pulsations. To predict these pulsations, it is important to consider all the excited rotor planes (and not only the one containing the fundamental space harmonic). With a difference of one bar quite significant differences are observed and explained.

For five-phase induction machines, it is possible to use two independent sequences to supply the stator. The proposed mathematical approach permitted to model the rotor under both sequences, to predict precisely torque pulsations frequencies for two different numbers of bars. These results allow to explain the important difference in torque ripple, between the five-phase machines with 64 and 65 rotor bars, observed in a previous paper [22].

To design a multiphase induction machine the combination of phase number and rotor bar number must be carefully chosen, considering the different possible polarities related to the used supply sequences. In the literature [1,2], the reported rules regarding this choice are defined for three-phase machines and are not directly applicable to multiphase induction machines. Thanks to the arithmetic tool and the corresponding graphical representation proposed in this paper, the wrong combinations can be predicted and avoided thanks to very fast calculations.

Further work is ongoing regarding the identification of mutual inductance coefficients " M_{sr_v} " to predict precisely the amplitudes of torque pulsations and the mean torque under different supplying sequences for multiphase induction machines and under different levels of magnetic saturation.

Author Contributions: A.M. proposed the used mathematical approach, worked on the numerical applications, analyzed the results, proposed the structure of the paper and wrote it. E.S. improved the structure of the paper and its content and provided advanced theoretical elements. F.S. improved the structure of the paper and its content and helped in the analysis of results. H.Z. improved the content of the paper. All authors validated the final paper. All authors have read and agreed to the published version of the manuscript.

Conflicts of Interest: The authors declare no conflict of interest.

Nomenclature

N_{ph}	Number of stator phases
N_{bar}	Number of rotor bars
spp	Number of slots per pole and per phase
p	Number of pole-pairs
\underline{V}_s	Stator voltage vector, dimension N_{ph}
\underline{I}_s	Stator current vector, dimension N_{ph}
\underline{I}_r	Rotor current vector, dimension N_{bar}
$\underline{\phi}_s$	Stator flux vector, dimension N_{ph}
$\underline{\phi}_r$	Rotor flux vector, dimension N_{bar}
\underline{R}_s	Stator resistance matrix, dimension $N_{ph} \times N_{ph}$ (Diagonal matrix)
\underline{R}_r	Rotor resistance matrix, dimension $N_{bar} \times N_{bar}$
\underline{L}_{ss}	Stator inductance matrix, dimension $N_{ph} \times N_{ph}$
\underline{L}_{rr}	Rotor inductance matrix, dimension $N_{bar} \times N_{bar}$
$\underline{L}_{sr} = \underline{L}_{rs}'$	Mutual inductance matrix (between stator and rotor), dimension $N_{ph} \times N_{bar}$
θ	Mechanical angle
R_s	Stator phase resistance
R_b	Rotor bar resistance
R_e	Resistance of end-connection ring portion (connecting 2 adjacent bars)
L_{sij}	Mutual inductance between two stator phases "i" and "j"
L_{ls}	Leakage stator inductance
L_{rij}	Mutual inductance between two rotor phases "i" and "j"
L_b	Rotor bar leakage inductance
L_e	Leakage inductance of end-connection ring portion
I_{su}	Stator current amplitude (for the stator sequence "u")
f_{su}	Stator current frequency (for the stator sequence "u")
φ_{su}	Stator current phase (for the stator sequence "u")
I_{rv}	Amplitude of induced rotor current harmonic "v"
f_{rv}	Frequency of the induced rotor current harmonic "v"
φ_{rv}	Phase of the induced rotor current harmonic "v"
Ω_{mec}	Mechanical speed

Appendix A

As seen in Figure 2, each space harmonic “k” is projected on a cell represented by $M_{k\alpha_{s,r}\beta_{s,r}}(\theta)$ (Equation (18)), which can be a two-by-two matrix, a two-element vector or a scalar, depending on the rules described below:

- Rule 7: If $\text{mod}(k, N_{\text{ph}}) \leq \text{floor}\left(\frac{N_{\text{ph}}}{2}\right)$ and $\text{mod}(k.p, N_{\text{bar}}) \leq \text{floor}\left(\frac{N_{\text{bar}}}{2}\right)$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \frac{\sqrt{N_{\text{ph}} \cdot N_{\text{bar}}}}{2} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) & \sin(-kp\theta + \varphi_{\text{sr}_k}) \\ -\sin(-kp\theta + \varphi_{\text{sr}_k}) & \cos(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 8: If $\text{mod}(k, N_{\text{ph}}) \leq \text{floor}\left(\frac{N_{\text{ph}}}{2}\right)$ and $\text{mod}(k.p, N_{\text{bar}}) > \text{floor}\left(\frac{N_{\text{bar}}}{2}\right)$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \frac{\sqrt{N_{\text{ph}} \cdot N_{\text{bar}}}}{2} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) & -\sin(-kp\theta + \varphi_{\text{sr}_k}) \\ -\sin(-kp\theta + \varphi_{\text{sr}_k}) & -\cos(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 9: If $\text{mod}(k, N_{\text{ph}}) > \text{floor}\left(\frac{N_{\text{ph}}}{2}\right)$ and $\text{mod}(k.p, N_{\text{bar}}) \leq \text{floor}\left(\frac{N_{\text{bar}}}{2}\right)$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \frac{\sqrt{N_{\text{ph}} \cdot N_{\text{bar}}}}{2} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) & \sin(-kp\theta + \varphi_{\text{sr}_k}) \\ \sin(-kp\theta + \varphi_{\text{sr}_k}) & -\cos(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 10: If $\text{mod}(k, N_{\text{ph}}) > \text{floor}\left(\frac{N_{\text{ph}}}{2}\right)$ and $\text{mod}(k.p, N_{\text{bar}}) > \text{floor}\left(\frac{N_{\text{bar}}}{2}\right)$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \frac{\sqrt{N_{\text{ph}} \cdot N_{\text{bar}}}}{2} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) & -\sin(-kp\theta + \varphi_{\text{sr}_k}) \\ \sin(-kp\theta + \varphi_{\text{sr}_k}) & \cos(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 11: If $\text{mod}(k, N_{\text{ph}}) = 0$ and $\text{mod}(k.p, N_{\text{bar}}) \leq \text{floor}\left(\frac{N_{\text{bar}}}{2}\right)$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{\frac{N_{\text{ph}} \cdot N_{\text{bar}}}{2}} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) & \sin(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 12: If $\text{mod}(k, N_{\text{ph}}) = 0$ and $\text{mod}(k.p, N_{\text{bar}}) > \text{floor}\left(\frac{N_{\text{bar}}}{2}\right)$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{\frac{N_{\text{ph}} \cdot N_{\text{bar}}}{2}} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) & -\sin(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 13 (Only when N_{ph} is even): If $\text{mod}(k, N_{\text{ph}}) = \frac{N_{\text{ph}}}{2}$ and $\text{mod}(k.p, N_{\text{bar}}) \leq \text{floor}\left(\frac{N_{\text{bar}}}{2}\right)$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{\frac{N_{\text{ph}} \cdot N_{\text{bar}}}{2}} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) & \sin(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 14 (Only when N_{ph} is even): If $\text{mod}(k, N_{\text{ph}}) = \frac{N_{\text{ph}}}{2}$ and $\text{mod}(k.p, N_{\text{bar}}) > \text{floor}\left(\frac{N_{\text{bar}}}{2}\right)$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{\frac{N_{\text{ph}} \cdot N_{\text{bar}}}{2}} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) & -\sin(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 15: If $\text{mod}(k, N_{\text{ph}}) \leq \text{floor}\left(\frac{N_{\text{ph}}}{2}\right)$ and $\text{mod}(k.p, N_{\text{bar}}) = 0$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{\frac{N_{\text{ph}} \cdot N_{\text{bar}}}{2}} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) \\ -\sin(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 16: If $\text{mod}(k, N_{\text{ph}}) > \text{floor}\left(\frac{N_{\text{ph}}}{2}\right)$ and $\text{mod}(k.p, N_{\text{bar}}) = 0$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{\frac{N_{\text{ph}} \cdot N_{\text{bar}}}{2}} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) \\ \sin(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 17 (Only when N_{bar} is even): If $\text{mod}(k, N_{\text{ph}}) \leq \text{floor}\left(\frac{N_{\text{ph}}}{2}\right)$ and $\text{mod}(k.p, N_{\text{bar}}) = \frac{N_{\text{bar}}}{2}$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{\frac{N_{\text{ph}} \cdot N_{\text{bar}}}{2}} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) \\ -\sin(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 18 (Only when N_{bar} is even): If $\text{mod}(k, N_{\text{ph}}) > \text{floor}\left(\frac{N_{\text{ph}}}{2}\right)$ and $\text{mod}(k.p, N_{\text{bar}}) = \frac{N_{\text{bar}}}{2}$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{\frac{N_{\text{ph}} \cdot N_{\text{bar}}}{2}} \cdot M_{\text{sr}_k} \begin{pmatrix} \cos(-kp\theta + \varphi_{\text{sr}_k}) \\ \sin(-kp\theta + \varphi_{\text{sr}_k}) \end{pmatrix}$$

- Rule 19: If $\text{mod}(k, N_{\text{ph}}) = 0$ and $\text{mod}(k.p, N_{\text{bar}}) = 0$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{N_{\text{ph}} \cdot N_{\text{bar}}} \cdot M_{\text{sr}_k}(\cos(-kp\theta + \varphi_{\text{sr}_k}))$$

- Rule 20 (Only when N_{bar} is even): If $\text{mod}(k, N_{\text{ph}}) = 0$ and $\text{mod}(k.p, N_{\text{bar}}) = \frac{N_{\text{bar}}}{2}$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{N_{\text{ph}} \cdot N_{\text{bar}}} \cdot M_{\text{sr}_k}(\cos(-kp\theta + \varphi_{\text{sr}_k}))$$

- Rule 21 (Only when N_{ph} is even): If $\text{mod}(k, N_{\text{ph}}) = \frac{N_{\text{ph}}}{2}$ and $\text{mod}(k.p, N_{\text{bar}}) = 0$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{N_{\text{ph}} \cdot N_{\text{bar}}} \cdot M_{\text{sr}_k}(\cos(-kp\theta + \varphi_{\text{sr}_k}))$$

- Rule 22 (Only when N_{ph} and N_{bar} are even): If $\text{mod}(k, N_{\text{ph}}) = \frac{N_{\text{ph}}}{2}$ and $\text{mod}(k.p, N_{\text{bar}}) = \frac{N_{\text{bar}}}{2}$:

$$\underline{M_{k\alpha_{s,r}\beta_{s,r}}}(\theta) = \sqrt{N_{\text{ph}} \cdot N_{\text{bar}}} \cdot M_{\text{sr}_k}(\cos(-kp\theta + \varphi_{\text{sr}_k}))$$

In the case where two, or more, harmonics are projected on the same stator and rotor planes, the matrices $\underline{M_{\alpha\beta_k}}$ related to these harmonics are superposed.

Appendix B

Under a stator sequence “u”, the pulsating torque component due to the interaction between a space harmonic “ $v_1.p$ ” and a rotor time harmonic “ $v_2.p$ ” can be expressed, in the general form, as follows:

$$T_{v_1-v_2} = \frac{N_{\text{ph}} \cdot N_{\text{bar}}}{4} \cdot p \cdot |v_1| \cdot I_{s_u} \cdot I_{r_{v_2}} \cdot (M_{\text{sr}})_{v_1} \cdot \sin(\delta_s(\omega_{s_u} \cdot t + \varphi_{s_u}) + \delta_r(|\omega_{r_{v_2}}| \cdot t + \varphi_{r_{v_2}}) + \delta_{\text{sr}}(-|v_1| \cdot p \cdot \Omega_{\text{mec}} \cdot t + \varphi_{\text{sr}_{v_1}})) \quad (27)$$

where: δ_s , δ_r and $\delta_{\text{sr}} = \pm 1$. These coefficients depend on the rules (Appendix A) related to the harmonics “ $v_1.p$ ” and “ $v_2.p$ ”. The different cases of the coefficients δ_s , δ_r and δ_{sr} are shown in the Table A1 (Only rules 7 to 10 are considered).

Table A1. Pulsating torque components due to interactions of harmonics of rules 7, 8, 9 and 10.

		$\text{mod}(u, N_{\text{ph}}) < \text{floor}(\frac{N_{\text{ph}}}{2})$				$\text{mod}(u, N_{\text{ph}}) > \text{floor}(\frac{N_{\text{ph}}}{2})$			
Rule of "v ₂ .p"	Rule of "v ₁ .p"	Rule 7	Rule 8	Rule 9	Rule 10	Rule 7	Rule 8	Rule 9	Rule 10
Rule 7	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$
	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = 1$	$\delta_r = 1$	$\delta_r = 1$	$\delta_r = 1$
	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = 1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$
Rule 8	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$
	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = 1$	$\delta_r = 1$	$\delta_r = 1$	$\delta_r = 1$	$\delta_r = 1$
	$\delta_{sr} = 1$	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$
Rule 9	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$
	$\delta_r = 1$	$\delta_r = 1$	$\delta_r = 1$	$\delta_r = 1$	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = -1$
	$\delta_{sr} = 1$	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$
Rule 10	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$	$\delta_s = -1$	$\delta_s = 1$
	$\delta_r = 1$	$\delta_r = 1$	$\delta_r = 1$	$\delta_r = 1$	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = -1$	$\delta_r = -1$
	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$	$\delta_{sr} = 1$	$\delta_{sr} = -1$	$\delta_{sr} = 1$

When a rotor zero-sequence is excited (rules 15, 16, 17 and 18 in Appendix A), the general form of the developed torque component due to the interaction of "v₃.p" and "v₄.p" (both exciting rotor zero-sequences) can be written as follows:

$$T_{v_3-v_4} = \frac{N_{\text{ph}} \cdot N_{\text{bar}}}{4} \cdot p \cdot |v_3| \cdot I_{s_u} \cdot I_{r_{v_4}} \cdot (M_{\text{sr}})_{v_3} \cdot \sin(\delta_{s1}(\omega_{s_u} \cdot t + \varphi_{s_u}) + \delta_{r1}(|\omega_{r_{v_4}}| \cdot t + \varphi_{r_{v_4}}) + \delta_{sr1}(-|v_3| \cdot p \cdot \Omega_{\text{mec}} \cdot t + \varphi_{sr_{v_3}})) + \sin(\delta_{s2}(\omega_{s_u} \cdot t + \varphi_{s_u}) + \delta_{r2}(|\omega_{r_{v_4}}| \cdot t + \varphi_{r_{v_4}}) + \delta_{sr2}(-|v_3| \cdot p \cdot \Omega_{\text{mec}} \cdot t + \varphi_{sr_{v_3}})) \quad (28)$$

The different cases of the coefficients δ_{s1} , δ_{s2} , δ_{r1} , δ_{r2} and δ_{sr1} , δ_{sr2} are shown in the Table A2.

Table A2. Pulsating torque components due to interactions of harmonics of rules 15, 16, 17 and 18.

		$\text{mod}(u, N_{\text{ph}}) < \text{floor}(\frac{N_{\text{ph}}}{2})$				$\text{mod}(u, N_{\text{ph}}) > \text{floor}(\frac{N_{\text{ph}}}{2})$			
Rule of "v ₄ .p"	Rule of "v ₃ .p"	Rule 15	Rule 16	Rule 17	Rule 18	Rule 15	Rule 16	Rule 17	Rule 18
Rule 15	$\delta_{s1} = 1$	$\delta_{s1} = 1$				$\delta_{s1} = 1$	$\delta_{s1} = 1$		
	$\delta_{r1} = -1$	$\delta_{r1} = -1$				$\delta_{r1} = 1$	$\delta_{r1} = 1$		
	$\delta_{sr1} = 1$	$\delta_{sr1} = 1$				$\delta_{sr1} = -1$	$\delta_{sr1} = -1$		
	$\delta_{s2} = -1$	$\delta_{s2} = -1$				$\delta_{s2} = -1$	$\delta_{s2} = -1$		
	$\delta_{r2} = -1$	$\delta_{r2} = -1$				$\delta_{r2} = 1$	$\delta_{r2} = 1$		
Rule 16	$\delta_{sr2} = -1$	$\delta_{sr2} = -1$				$\delta_{sr2} = 1$	$\delta_{sr2} = 1$		
	$\delta_{s1} = 1$	$\delta_{s1} = 1$				$\delta_{s1} = 1$	$\delta_{s1} = 1$		
	$\delta_{r1} = 1$	$\delta_{r1} = 1$				$\delta_{r1} = -1$	$\delta_{r1} = -1$		
	$\delta_{sr1} = -1$	$\delta_{sr1} = -1$				$\delta_{sr1} = 1$	$\delta_{sr1} = 1$		
	$\delta_{s2} = -1$	$\delta_{s2} = -1$				$\delta_{s2} = -1$	$\delta_{s2} = -1$		
Rule 17	$\delta_{r2} = 1$	$\delta_{r2} = 1$				$\delta_{r2} = -1$	$\delta_{r2} = -1$		
	$\delta_{sr2} = 1$	$\delta_{sr2} = 1$				$\delta_{sr2} = -1$	$\delta_{sr2} = -1$		
			$\delta_{s1} = 1$	$\delta_{s1} = 1$				$\delta_{s1} = 1$	$\delta_{s1} = 1$
			$\delta_{r1} = -1$	$\delta_{r1} = -1$				$\delta_{r1} = 1$	$\delta_{r1} = 1$
			$\delta_{sr1} = 1$	$\delta_{sr1} = 1$				$\delta_{sr1} = -1$	$\delta_{sr1} = -1$
Rule 18			$\delta_{s2} = -1$	$\delta_{s2} = -1$				$\delta_{s2} = -1$	$\delta_{s2} = -1$
			$\delta_{r2} = -1$	$\delta_{r2} = -1$				$\delta_{r2} = 1$	$\delta_{r2} = 1$
			$\delta_{sr2} = -1$	$\delta_{sr2} = -1$				$\delta_{sr2} = 1$	$\delta_{sr2} = 1$
			$\delta_{s1} = 1$	$\delta_{s1} = 1$				$\delta_{s1} = 1$	$\delta_{s1} = 1$
			$\delta_{r1} = 1$	$\delta_{r1} = 1$				$\delta_{r1} = -1$	$\delta_{r1} = -1$
		$\delta_{sr1} = -1$	$\delta_{sr1} = -1$				$\delta_{sr1} = 1$	$\delta_{sr1} = 1$	
		$\delta_{s2} = -1$	$\delta_{s2} = -1$				$\delta_{s2} = -1$	$\delta_{s2} = -1$	
		$\delta_{r2} = 1$	$\delta_{r2} = 1$				$\delta_{r2} = -1$	$\delta_{r2} = -1$	
		$\delta_{sr2} = 1$	$\delta_{sr2} = 1$				$\delta_{sr2} = -1$	$\delta_{sr2} = -1$	

For the rotor zero-sequence harmonics, the interaction between a space harmonic "v.p" and the time harmonic (of rotor current) of the same range "v.p" generates a constant torque and a pulsating component in the same time.

Appendix C

The three-phase induction machine has an overlapping distributed winding, its geometry is shown in Figure A1.

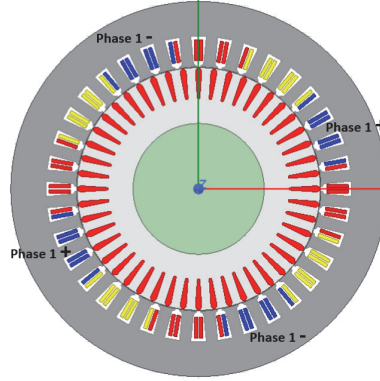


Figure A1. Geometry of the three-phase induction machine.

The parameters of this machine are:

- Number of phases, $N_{ph} = 3$;
- Number of stator slots, $N_s = 36$;
- Number of rotor bars (2 compared cases), $N_{bar} = 48, 49$;
- Number of pole-pairs under first current harmonic injection, $p_1 = 2$;
- Rated stator current, $I_s = 325 A_{peak}$;
- Rated torque, $T = 72 N.m$.

The Table A3 shows the harmonic winding factors for the 50 first space harmonics. The winding factor “ k_v ” divided by its range “ v ” reflects the importance of the harmonic in torque production. In fact, the coefficient $M_{sr,v}$ is proportional to the factor “ k_v/v ” (in the column 3). The lines containing the 15 most important space harmonics are colored in gray in the table.

Table A3. Harmonic winding factors of the three-phase winding.

v	k_v	k_v/v	v	k_v	k_v/v	v	k_v	k_v/v	v	k_v	k_v/v	v	k_v	k_v/v
1	0.945	0.945	11	0.061	0.006	21	0.577	0.027	31	0.140	0.005	41	0.140	0.003
2	0	0	12	0	0	22	0	0	32	0	0	42	0	0
3	0.577	0.192	13	0.140	0.011	23	0.140	0.006	33	0.577	0.017	43	0.061	0.001
4	0	0	14	0	0	24	0	0	34	0	0	44	0	0
5	0.140	0.028	15	0.577	0.038	25	0.061	0.002	35	0.945	0.027	45	0	0
6	0	0	16	0	0	26	0	0	36	0	0	46	0	0
7	0.061	0.009	17	0.945	0.056	27	0	0	37	0.945	0.026	47	0.061	0.001
8	0	0	18	0	0	28	0	0	38	0	0	48	0	0
9	0	0	19	0.945	0.050	29	0.061	0.002	39	0.577	0.015	49	0.140	0.003
10	0	0	20	0	0	30	0	0	40	0	0	50	0	0

Appendix D

The five-phase induction machine has a tooth-concentrated winding, its geometry is shown in Figure A2. The parameters of this machine are:

- Number of phases, $N_{ph} = 5$;
- Number of stator slots, $N_s = 20$;
- Number of rotor bars (2 compared cases), $N_{bar} = 64, 65$;
- Number of pole-pairs under first current harmonic injection, $p_1 = 4$;
- Rated stator current, $I_s = 400 A_{peak}$;
- Rated torque, $T = 250 N.m$.

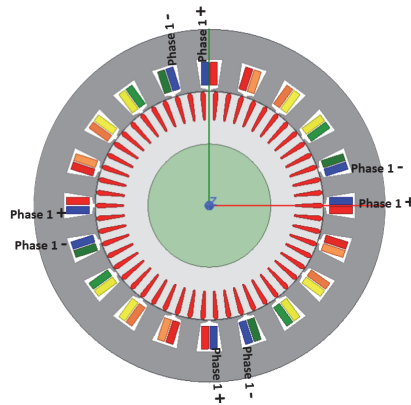


Figure A2. Geometry of the five-phase induction machine with double-layer tooth concentrated winding.

The Table A4 shows the harmonic winding factors for the 50 first space harmonics. The winding factor “ k_v ” divided by its range “ v ” reflect the importance of the harmonic in torque production. In fact, the coefficient M_{sr_v} is proportional to the factor “ k_v/v ” (in the column 3). The lines containing the 20 most important space harmonics are colored in gray in the table.

Table A4. Harmonic winding factors of the five-phase winding.

v	k_v	k_v/v	v	k_v	k_v/v	v	k_v	k_v/v	v	k_v	k_v/v	v	k_v	k_v/v
1	0.59	0.59	11	0.59	0.05	21	0.59	0.03	31	0.59	0.02	41	0.59	0.01
2	0.95	0.48	12	0.95	0.08	22	0.95	0.04	32	0.95	0.03	42	0.95	0.02
3	0.95	0.32	13	0.95	0.07	23	0.95	0.04	33	0.95	0.03	43	0.95	0.02
4	0.59	0.15	14	0.59	0.04	24	0.59	0.02	34	0.59	0.02	44	0.59	0.01
5	0	0	15	0	0	25	0	0	35	0	0	45	0	0
6	0.59	0.10	16	0.59	0.04	26	0.59	0.02	36	0.59	0.02	46	0.59	0.01
7	0.95	0.14	17	0.95	0.06	27	0.95	0.04	37	0.95	0.03	47	0.95	0.02
8	0.95	0.12	18	0.95	0.05	28	0.95	0.03	38	0.95	0.03	48	0.95	0.02
9	0.59	0.07	19	0.59	0.03	29	0.59	0.02	39	0.59	0.02	49	0.59	0.01
10	0	0	20	0	0	30	0	0	40	0	0	50	0	0

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