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A new analytical model for predicting the radio-thermal oxidation kinetics and the lifetime of electric cable insulation in nuclear power plants. application to silane cross-linked polyethylene

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A B S T R A C T

The radio-thermal oxidation of silane cross-linked polyethylene (Si-XLPE) was studied in air under different γ dose rates (6.0, 8.5, 77.8, and 400 Gy.h⁻¹) at different temperatures (21, 47, and 86°C). The changes in the physico-chemical and mechanical properties of Si-XLPE throughout its exposure were determined by FTIR spectroscopy, differential scanning calorimetry (DSC), swelling measurements, rheometry in rubbery (DMTA) and in molten states, and uniaxial tensile testing. It was found that oxidation leads to the build-up of a wide variety of carbonyl and hydroxyl products (mostly carboxylic acids and hydroperoxides) and an efficient chain scission process that catastrophically reduces the concentration in elastically active chains and the elongation at break from the early periods of exposure. A new analytical model was derived from the current radio-thermal mechanistic scheme without making the usual assumption of thermal stability of hydroperoxides. After an initial period where the oxidation kinetics occurs with a constant rate, this model allows also predicting the auto-acceleration of the oxidation kinetics when the hydroperoxide concentration reaches a critical value of about 1.6×10^{-1} mol.L⁻¹. Choosing this critical value as a structural end-of-life criterion allows a more direct assessment of the lifetime of Si-XLPE in the various radio-thermal environments under study, except at the highest temperature (i.e. 86°C) where the kinetic model can still be noticeably improved.

Keywords:

Silane cross-linked polyethylene
Radio-thermal oxidation
Analytical kinetic model
Hydroperoxide decomposition
Carbonyl build-up
Chain scissions
Structural end-of-life criterion
Lifetime prediction

1. Introduction

Since the early 1980s, in order to meet the urgent need of nuclear power plant operators, the polymer aging community is looking for a non-empirical approach for lifetime prediction of polymeric materials exposed in nuclear environments. A practical case that mobilized renowned research teams around the world is the radio-thermal ageing of electric cable insulations under normal operating conditions in a nuclear reactor. According to Khelidj et al. [1], at room temperature in air, the lifetime t_F of electric cable insulations made of unfilled and unstabilised polyethylene (PE) varies with dose rate I according to the general trend depicted in Fig. 1. This almost universal long-term behavior of PE seems to be completely independent of its macromolecular structure (i.e. linear, branched or chemically cross-linked) and its crystallinity.

This curve displays a negative curvature in the domain of low dose rates where it tends towards a horizontal asymptote (when

$I \rightarrow 0$) corresponding to the “ceiling lifetime” t_{F0} due to “pure” thermal oxidation. Schematically, it can be decomposed into three distinct domains, each one corresponding to a different kinetic regime:

- “Pure” thermal oxidation predominates at very low dose rates (typically for $I < I_T \approx 1.6 \times 10^{-7}$ Gy.s⁻¹) and is characterised by very long lifetimes:

$$t_{F0} \approx 18.5 \text{ years} \quad (1)$$

- “Pure” radio-oxidation predominates at high dose rates (typically for $I > I_R \approx 5.0 \times 10^{-1}$ Gy.s⁻¹) and is characterized by a lethal dose independent of the dose rate. As a result, the lifetime is proportional to the reciprocal dose rate:

$$t_F = 1.6 \times 10^5 I^{-1} \quad (2)$$

- A combination of the two previous reactions occurs at intermediary dose rates (i.e. for $I_T < I < I_R$), which is the range of practical interest for nuclear power plant operators. This radio-thermal oxidation is initiated by two competitive sources of radicals, neither of which can be totally neglected compared to

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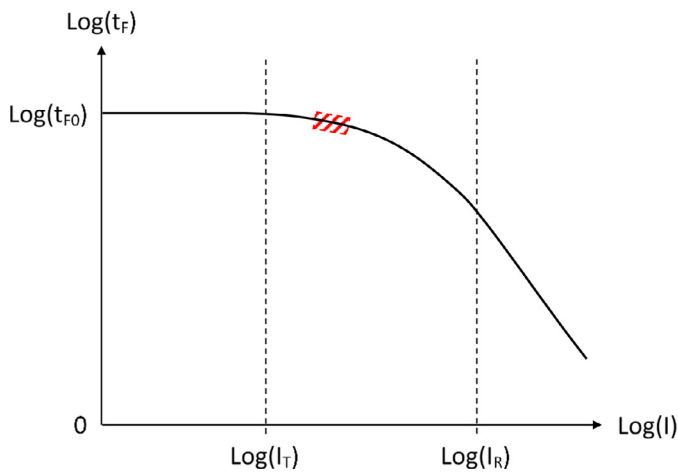
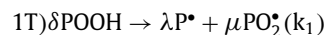
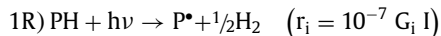


Fig. 1. Shape in logarithm co-ordinates of the lifetime t_F versus γ dose rate I in air at room temperature for unstabilised and unfilled PE [1]. The shaded area (in red) corresponds to the normal operating conditions of electric cables in a nuclear reactor.

the other during the exposure [2]: polymer radiolysis (1R) and thermal decomposition of hydroperoxides (1T):



where r_i is the rate and G_i is the radical yield (expressed in number of radicals P^{\bullet} per 100 eV absorbed) of the radiochemical initiation. k_1 is the rate constant and δ , λ , and μ are stoichiometric coefficients describing the molecularity of thermal initiation such as:

$k_1 = k_{1u}$, $\delta = 1$, $\lambda = 2$ and $\mu = 0$ for unimolecular decomposition,
and $k_1 = k_{1b}$, $\delta = 2$ and $\lambda = \mu = 1$ for bimolecular decomposition.

To our knowledge, there is still no analytical equation for describing the entire shape of the curve $t_F = f(I)$ reported in Fig. 1.

In the second past half century, the vast majority of studies were devoted to the rigorous identification and understanding of the oxidation mechanisms, but exclusively at high dose rates where the duration and cost of aging experiments under irradiation can be significantly reduced [3-8]. When $I > I_R$, it was noted that the induction period is really too short to be measured. It was thus concluded that the oxidation kinetics already starts in steady regime from the early periods of exposure. In addition, it was also observed that the effect of dose rate is a simple kinetic consequence of the radio-oxidation reaction. Schematically, the formation rate of oxidation products is proportional to $I^{1/2}$ when they result from the propagation step (of the mechanistic scheme), whereas it is proportional to I when they result from the initiation or termination step.

Based on only two assumptions: steady-state (for radical species) and thermal stability of hydroperoxides throughout the exposure, analytical equations were proposed in the literature to account for these experimental trends. As an example, the two following general equations were determined from a kinetic analysis of the mechanistic scheme for the formation rate r_Y (expressed in mol.L^{-1}) and the radiochemical yield G_Y (dimensionless) of a given chemical event Y , respectively [1]:

$$r_Y = 10^{-7} G_Y I \quad (3)$$

$$\text{with } G_Y = 10^7 a_Y + 10^7 b_Y I^{-1/2} \quad (4)$$

where a_Y and $b_Y I^{-1/2}$ are initiation/termination and propagation terms, respectively.

However, even at low temperature close to ambient, the second assumption (thermal stability of hydroperoxides) can only be valid for sufficiently short durations. Indeed, it should be recalled that the rate of the thermal decomposition of hydroperoxides is an increasing function of their concentration [9]. Consequently, thermal initiation is expected to gradually gain in importance with the accumulation of hydroperoxides throughout the exposure, to finally become competitive with the radiochemical initiation. It is obvious that the global oxidation kinetics should be even more impacted in the case of an industrial application for which a very long lifetime of several decades is expected.

To our knowledge, the pioneering researchers in this field, to have proposed a kinetic model in which irradiation generates hydroperoxides which decompose thermally to give radicals initiating new oxidation chains, are Gillen & Clough [10]. Their model predicts an auto-acceleration of the oxidation kinetics during the exposure whereas, in all the previous models (for instance, see Eq. (3)), oxidation occurred at a constant rate. Unfortunately, Gillen & Clough chosen questionable initial conditions: $[POOH] = [POOH]_0 = 0$ at $t = 0$. However, polymers (especially industrial grades having undergone processing operations) contain always thermally unstable species such as hydroperoxides, otherwise it would be difficult to explain their thermal or photochemical oxidizability. As a consequence of these unrealistic initial conditions, the model developed by Gillen & Clough diverges in the domain of very low dose rates. Indeed, it predicts that $t_F \rightarrow \infty$ when $I \rightarrow 0$.

In most recent publications [11,12], Gillen tried to get around this problem by relying on the research works of Ito [13-16] and Mares et al. [17]. He considered an extremely simplified radio-thermal oxidation mechanistic scheme composed of only three first-order reactions: "pure" radio-oxidation, "pure" thermal oxidation, and their coupling, which is assumed to lead to synergistic effects. To take into account these synergistic effects, he defined a global time-dependent rate constant k_{R+T} as follows:

$$k_{R+T}(t) = k_R(t) + k_T(t) + k_S(t) \quad (5)$$

where k_R and k_T are effective rate constants to be determined experimentally from the changes in mechanical properties (conventionally, elongation at break) in pure radiochemical (i.e. under high dose rate at low temperature) and pure thermal environments (i.e. in the absence of irradiation), respectively. The effective rate constant k_S of the synergistic reaction is deduced from the changes in mechanical properties measured in radio-thermal environment.

Finally, Gillen defined the synergistic effects as follows:

$$q(t) = \frac{k_{R+T}(t)}{k_R(t) + k_T(t)} = 1 + \frac{k_S(t)}{k_R(t) + k_T(t)} \quad (6)$$

If the ratio $q > 1$, the synergistic effects are clearly indicated. However, as the rate constants k_R , k_T and k_S are time dependent, q varies non-monotonically throughout the exposure [11].

However, at its current state, this new approach is not yet completely convincing for the following reasons. From a practical point of view, it does not allow avoiding the issue of the choice of initial conditions. Instead of an initial concentration ($[POOH]_0$), it is now necessary to determine an initial rate constant for thermal oxidation (k_{T0}), but also for radio-oxidation (k_{R0}) and for synergistic reaction (k_{S0}).

From the perspective of a non-empirical lifetime prediction, this new approach raises a series of fundamental issues that have never been addressed by its author, for instance:

- i) Eq. (5) was postulated but never justified. Thus, the assumption of a simple linear combination of the contributions of the

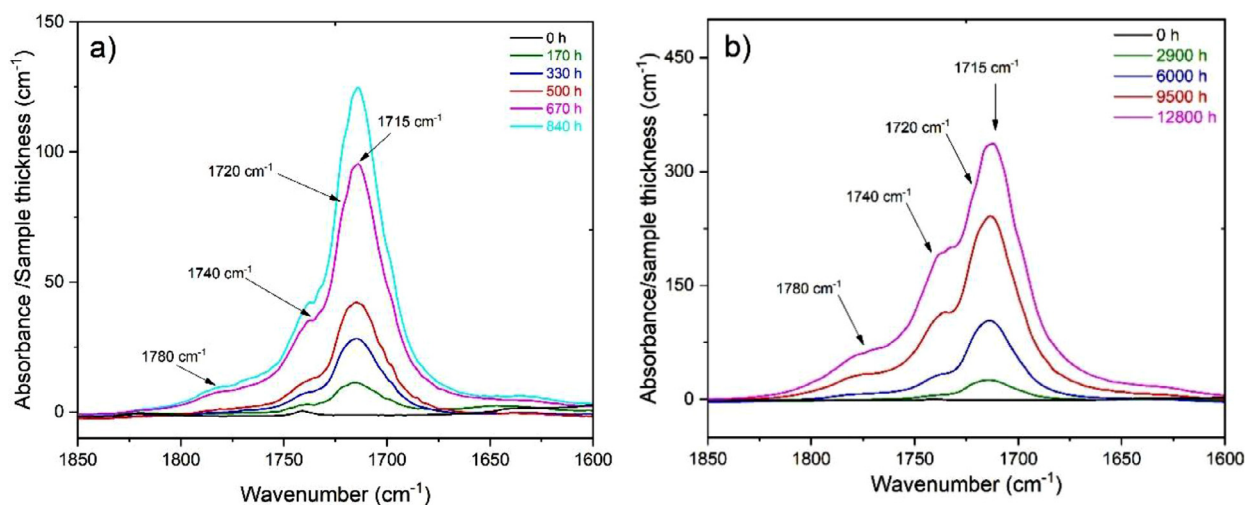


Fig. 3. Changes in carbonyl region of the FTIR spectrum of Si-XLPE before and after radio-thermal ageing in air under $400 \text{ Gy}\cdot\text{h}^{-1}$ at 21°C (a) and under $8.5 \text{ Gy}\cdot\text{h}^{-1}$ at 47°C (b).

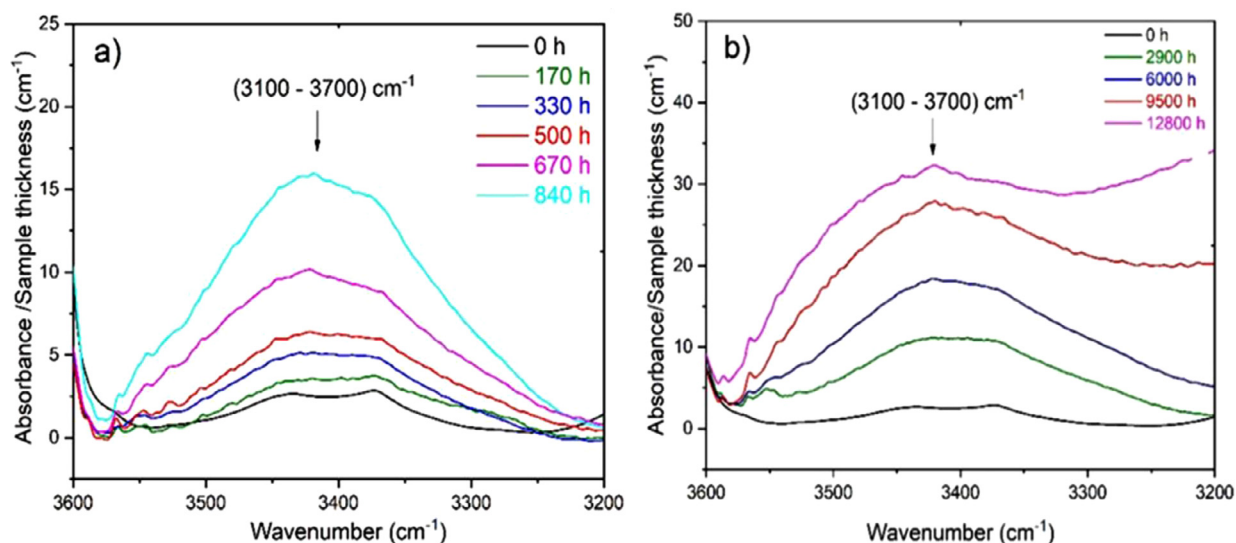


Fig. 4. Changes in the hydroxyl region of the FTIR spectrum of Si-XLPE before and after radio-thermal ageing in air under $400 \text{ Gy}\cdot\text{h}^{-1}$ at 21°C (a) and under $8.5 \text{ Gy}\cdot\text{h}^{-1}$ at 47°C (b).

3700 cm^{-1} , see Fig. 4) also quickly grows. According to literature [25,27-30], this band is attributed to the stretching vibration of the O-H bond of hydroperoxides, alcohols, and carboxylic acids.

The concentrations of carboxylic acids and hydroxyl products [Prod] were determined by applying the classical Beer-Lambert's law:

$$[\text{Prod}] = \frac{\text{OD}}{\text{ep} \varepsilon} \quad (7)$$

where OD is the optical density of the IR absorption band (dimensionless), ε is the molar extinction coefficient for the corresponding oxidation product (expressed in $\text{L}\cdot\text{mol}^{-1}\cdot\text{cm}^{-1}$), and ep is the film thickness (in cm).

The orders of magnitude of ε were determined in a previous publication [28]. It was found that $\varepsilon = 680$ and $70 \text{ L}\cdot\text{mol}^{-1}\cdot\text{cm}^{-1}$ for carboxylic acids and hydroxyls, respectively.

2.3.3. Differential scanning calorimetry

Differential scanning calorimetry (DSC) was used under pure nitrogen to monitor and titrate the hydroperoxides (POOH), but also to highlight the changes in crystalline morphology during the exposure. DSC thermograms were recorded with a TA instruments

DSC Q1000 calorimeter beforehand calibrated with an indium reference. Film samples with a mass ranged between 5 and 10 mg were introduced in a closed standard aluminum pan to be analyzed between -50°C and 250°C , with a heating rate of $10^\circ\text{C}\cdot\text{min}^{-1}$ under a nitrogen flow of $50 \text{ mL}\cdot\text{min}^{-1}$.

As an example, Fig. 5 shows the changes in the DSC thermogram in air under the highest (i.e. $400 \text{ Gy}\cdot\text{h}^{-1}$) and lowest dose rates ($8.5 \text{ Gy}\cdot\text{h}^{-1}$) at low temperature close to ambient. It can be clearly observed that an exothermic peak starts to grow above the melting zone of the crystalline phase. This new peak is typically ranged between 130 and 230°C . As shown in a previous publication [28], it results from the thermal decomposition of POOH in the DSC cavity.

The concentration of hydroperoxides [POOH] was determined as follows:

$$[\text{POOH}] = \frac{\Delta H_{\text{POOH}}}{\Delta H_{\text{theory}}} \quad (8)$$

where ΔH_{POOH} is the area under the exothermic peak between 130 and 230°C on the DSC thermogram and ΔH_{theory} is the theoretical value of the decomposition enthalpy of POOH, which can be

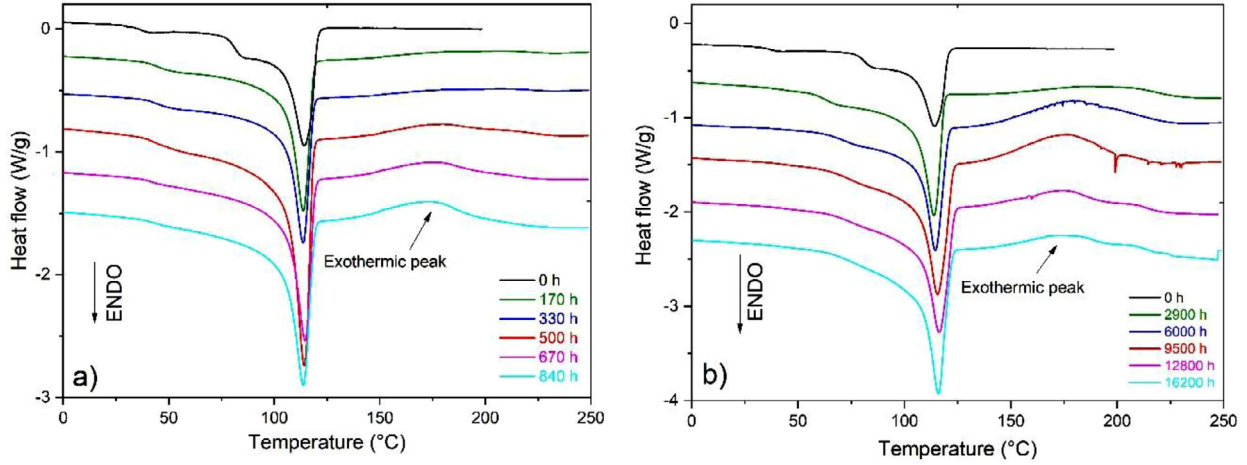


Fig. 5. Changes in the DSC thermogram of Si-XLPE before and after radio-thermal ageing in air under $400 \text{ Gy}\cdot\text{h}^{-1}$ at 21°C (a) and under $8.5 \text{ Gy}\cdot\text{h}^{-1}$ at 47°C (b).

calculated from the classical theoretical concepts of thermochemistry or determined experimentally from POOH model compounds: $\Delta H_{\text{theory}} = 291 \text{ kJ}\cdot\text{mol}^{-1}$ [28].

In addition, the global crystallinity ratio X_C of the polymer was determined with the common equation:

$$X_C = \frac{\Delta H_m}{\Delta H_{m0}} \times 100 \quad (9)$$

where ΔH_m is the sum of the areas under the endothermic peaks observed between 50 and 125°C on the DSC thermogram, and ΔH_{m0} is the melting enthalpy of the PE crystal: $\Delta H_{m0} = 292 \text{ J}\cdot\text{g}^{-1}$ [31].

2.3.4. Swelling measurements

Film samples were swelled with a good solvent to determine their swelling ratio and deduce the concentration of elastically active chains in the macromolecular network. A sample having an initial mass m_i was introduced into xylene previously heated at 130°C for 24 hours until reaching its equilibrium mass m_s after complete swelling. This sample was then dried under vacuum at 80°C for 48 hours in order to determine the gel mass m_d . The swelling ratio Q was calculated as follows:

$$Q = 1 + \frac{\rho_{\text{pol}}}{\rho_{\text{sol}}} \cdot \frac{(m_s - m_d)}{m_d} \quad (10)$$

where ρ_{pol} and ρ_{sol} are the respective densities of Si-XLPE ($\rho_{\text{pol}} = 0.912$) and xylene ($\rho_{\text{sol}} = 0.761$).

The Flory-Rhener's theory [32] was then used to determine the concentration of elastically active chains ν between two consecutive chemical cross-links in the macromolecular network:

$$\nu = \frac{[\ln(1 - V_{r0}) + V_{r0} + \chi V_{r0}^2]}{V_{\text{sol}} \left[V_{r0}^{\frac{1}{f}} - \frac{2V_{r0}}{f} \right]} \quad (11)$$

where V_{sol} is the molar volume of xylene ($V_{\text{sol}} = 139.3 \text{ cm}^3\cdot\text{mol}^{-1}$), χ is the Flory-Huggins' interaction parameter between xylene and Si-XLPE ($\chi = 0.31$ [33,34]), f is the functionality of chemical cross-links ($f = 4$ for Si-XLPE) and V_{r0} is the volumic fraction of polymer in the swollen macromolecular network, of which expression is given below:

$$V_{r0} = \frac{1}{1 + \frac{(Q-1)\rho_{\text{pol}}}{\rho_{\text{sol}}}} \quad (12)$$

2.3.5. Rheometry in solid state (DMTA)

Rheometry in solid state (DMTA) was used in tension mode to determine the storage modulus E' at the rubbery plateau and

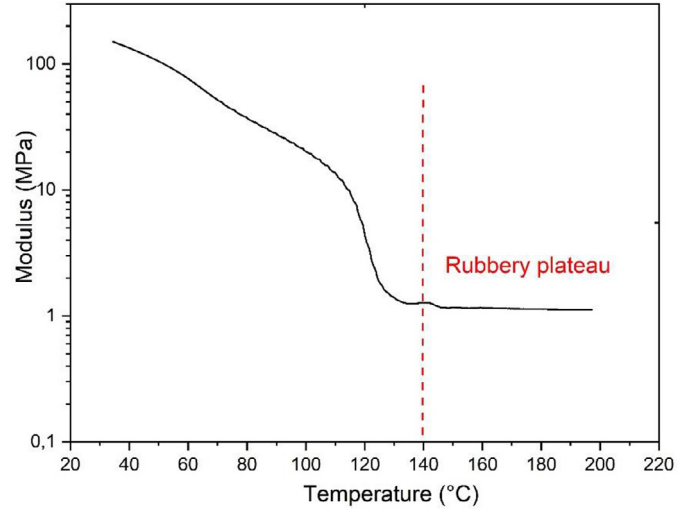


Fig. 6. DMTA thermogram of unaged Si-XLPE.

check the concentrations of elastically active chains determined by swelling in xylene. DMTA thermograms were recorded with a TA instruments DMA Q800 viscoelastic analyzer. Parallelepipedic film samples of about 30.3 mm long and 6.3 mm large were attached to the device clamps to be analyzed between 25 and 200°C , with a heating rate of $3^\circ\text{C}\cdot\text{min}^{-1}$, a frequency of 1 Hz , and a strain amplitude of 0.2% under nitrogen.

As an example, Fig. 6 shows the DMTA thermogram of Si-XLPE before its radio-thermal ageing. It can be clearly observed that the melting of the crystalline phase leads to sudden decrease of the storage modulus of more than one decade around 120°C . The rubbery plateau of the completely amorphous Si-XLPE is clearly visible for temperatures higher than 140°C .

The Flory's theory [35] was used to determine the total concentration of elastically active chains ν_t from the values of E' measured at 140°C :

$$E' = 3\nu_t\rho_{\text{pol}}RT \quad (13)$$

where ρ_{pol} is the polymer density at 140°C ($\rho_{\text{pol}} = 897 \text{ kg}\cdot\text{m}^{-3}$), R is the perfect gas constant ($R = 8.314 \text{ Pa}\cdot\text{m}^3\cdot\text{mol}^{-1}\cdot\text{K}^{-1}$), and T is the absolute temperature (expressed in K).

When the macromolecular network is not swollen by a solvent, in addition to chemical cross-links, it is expected that E' depends also on a second type of cross-links, which are not of a chemical

but physical nature, namely entanglements. In a first approach, ν_t can be written as the sum of these two contributions:

$$\nu_t = \nu + \nu_e \quad (14)$$

where ν_e would be the concentration of elastically active chains resulting only from entanglements in the macromolecular network (i.e. in the total absence of chemical cross-links). The value of ν_e can be simply deduced by comparing the rheometry and swelling data.

For the unaged Si-XLPE, it was found that: $\nu_t = 0.14 \pm 0.01 \text{ mol.L}^{-1}$ and $\nu = 0.12 \pm 0.02 \text{ mol.L}^{-1}$, so that: $\nu_e = 0.03 \pm 0.02 \text{ mol.L}^{-1}$. In other words, it was found that ν_e is very low and almost of the same order of magnitude than the experimental scattering of ν , i.e. almost undetectable. It was thus concluded that there are too few entanglements to significantly influence the concentration of elastically active chains. Consequently, entanglements were neglected in Si-XLPE and the concentration of elastically active chains ν resulting only from chemical cross-links was simply deduced as follows:

$$\nu = \frac{E'}{3\rho_{\text{pol}}RT} \quad (15)$$

2.3.6. Rheometry in molten state

For the radio-thermal ageing conditions under consideration, chain scissions were found very efficient in Si-XLPE. Indeed, a very advanced damage of the macromolecular network was observed for moderate to high conversion ratios of the oxidation reaction, as the film samples were almost completely dissolved in xylene and their rubbery plateau was no longer detected by DMTA (the polymer passed directly from semi-crystalline to molten state). Therefore, when swelling and DMTA measurements were no longer possible, rheometry in molten state was performed at 140°C to determine the shear modulus G' and thus, access the concentration of elastically active chains. Strain sweep curves were recorded with a TA Instruments ARES viscoelastic analyzer. Disc film samples of 25 mm diameter and 1 mm thick were introduced in a coaxial parallel plate geometry to be analyzed at 140°C with an angular frequency of 1 rad.s^{-1} and a strain amplitude ranging between 0.01 and 100% under nitrogen flow.

As an example, Fig. 7 shows the strain sweep curve after 167 h of radio-thermal exposure in air under the highest dose rate (i.e. 400 Gy.h^{-1}) at room temperature. As expected, it can be observed that G' reaches a constant value in the domain of low strains (typically for $\varepsilon < 10\%$) where linear viscoelasticity is valid. This value

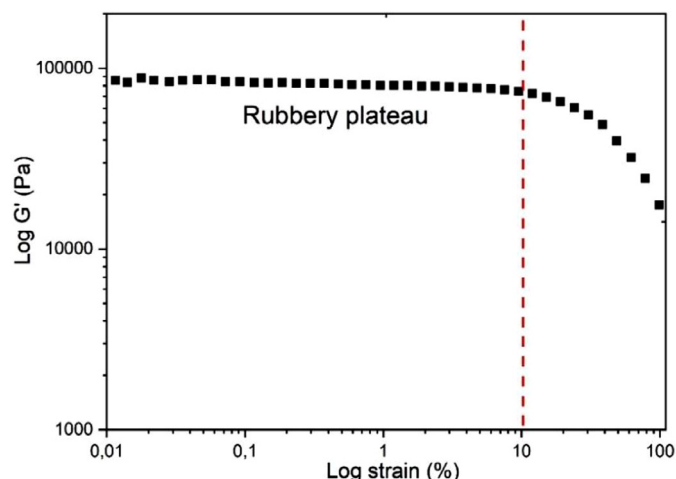


Fig. 7. Sweep strain curve of Si-XLPE after 167 h of exposure in air under 400 Gy.h^{-1} at 21°C.

was thus chosen for calculating the concentration of elastically active chains with Eq. (15). Knowing that: $E' = 3G'$, it comes finally:

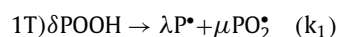
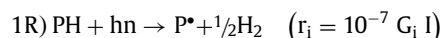
$$\nu = \frac{G'}{\rho_{\text{pol}}RT} \quad (16)$$

3. Theory

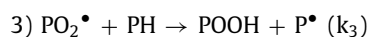
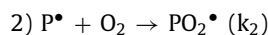
3.1. Radio-thermal oxidation mechanistic scheme

As said in introduction, in the domain of practical interest for nuclear power plant operators (i.e. for $1.6 \times 10^{-7} < I < 5.0 \times 10^{-1} \text{ Gy.s}^{-1}$), oxidation is initiated by both polymer radiolysis (1R) and thermal decomposition of hydroperoxides (1T). The corresponding radio-thermal oxidation mechanistic scheme can be thus written as follows:

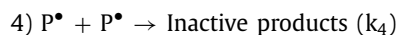
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Propagation:



Terminations:

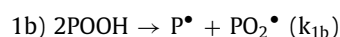


where PH, POOH, P^{\bullet} , and PO_2^{\bullet} designate an oxidation site, an hydroperoxide, alkyl and peroxy radicals, respectively. δ , λ , and μ are stoichiometric coefficients. r_i , G_i , and k_j (with $j = 1, \dots, 6$) are a rate, the radical yield, and rate constants, respectively.

In most cases of "pure" thermal oxidation, POOH decomposition is exclusively unimolecular at high temperature (i.e. $k_1 = k_{1u}$, $\delta = 1$, $\lambda = 2$, and $\mu = 0$) or mainly bimolecular at moderate to low temperature ($k_1 = k_{1b}$, $\delta = 2$, and $\lambda = \mu = 1$), but rarely a combination of both modes. In fact, it can easily be shown that there is a hyperbolic shape boundary delimiting the two domains in which each thermal initiation predominates over each other in a map (T , t). This boundary corresponds to a critical hydroperoxide concentration $[\text{POOH}]_C$ reached when the respective rates of the two thermal initiation modes are equal [3]. Since the activation energy is higher for the unimolecular than bimolecular mode: $E_{1u} > E_{1b}$, $[\text{POOH}]_C$ is an increasing function of temperature. In PE, it writes [36]:

$$[\text{POOH}]_C = \frac{k_{1u}}{k_{1b}} = 2.9 \times 10^3 \text{ Exp}\left(-\frac{35000}{RT}\right) \quad (17)$$

The numerical application of Eq. (17) gives that $[\text{POOH}]_C$ would be ranged between 10^{-3} and $10^{-2} \text{ mol.L}^{-1}$ at low temperature close to ambient for PE. This is the same order of magnitude of the initial concentration of hydroperoxides $[\text{POOH}]_0$ usually reported in the literature for many hydrocarbon polymers [28,37]. That is the reason why, in the present study, it will be considered that "pure" thermal oxidation is directly initiated in bimolecular mode:



3.2. System of differential equations

At the low conversion ratios of the oxidation reaction (i.e. for $[\text{PH}] \approx [\text{PH}]_0 \approx \text{constant}$), the following system of differential equations can be derived from the previous radio-thermal oxidation mechanistic scheme:

$$\frac{d[\text{P}^\bullet]}{dt} = r_i + k_{1b}[\text{POOH}]^2 - k_2C[\text{P}^\bullet] + k_3[\text{PO}_2^\bullet][\text{PH}] - 2k_4[\text{P}^\bullet]^2 - k_5[\text{P}^\bullet][\text{PO}_2^\bullet] \quad (18)$$

$$\frac{d[\text{PO}_2^\bullet]}{dt} = k_{1b}[\text{POOH}]^2 + k_2C[\text{P}^\bullet] - k_3[\text{PO}_2^\bullet][\text{PH}] - k_5[\text{P}^\bullet][\text{PO}_2^\bullet] - 2k_6[\text{PO}_2^\bullet]^2 \quad (19)$$

$$\frac{d[\text{POOH}]}{dt} = -2k_{1b}[\text{POOH}]^2 + k_3[\text{PO}_2^\bullet][\text{PH}] \quad (20)$$

where the oxygen concentration C is related to the oxygen partial pressure P_{O_2} in the exposure environment by the classical Henry's law:

$$C = S \cdot P_{\text{O}_2} \quad (21)$$

and S is the coefficient of oxygen solubility into the polymer. Typical values of S reported for low density polyethylene (LDPE) in the literature are about $1.8 \times 10^{-8} \text{ mol.L}^{-1}.\text{Pa}^{-1}$ regardless the temperature [38]. Knowing that in air under atmospheric pressure: $P_{\text{O}_2} = 0.21 \times 10^5 \text{ Pa}$, it comes finally: $C = 3.8 \times 10^{-4} \text{ mol.L}^{-1}$.

There is no rigorous analytical solution for this system of differential equations, which can be solved numerically without any simplifying assumption [2]. However, an approached analytical solution would be very useful because it would allow appreciating the behavioral trends of the reactive species $[\text{P}^\bullet]$, $[\text{PO}_2^\bullet]$, $[\text{POOH}]$ with key exposure parameters, in particular: γ dose rate I , temperature T , and oxygen partial pressure P_{O_2} . In addition, this analytical solution would then allow calculating key physico-chemical properties, from a practical point of view, because they can be checked experimentally, such as: oxygen consumption (Q), carbonyl build-up ($[\text{P}=\text{O}]$), damage state of macromolecular network (ν), etc. (see section 3.4). Finally, if a structural end-of-life criterion was identified, it would be then possible to propose an equation for predicting the lifetime of the insulating material (see section 3.5).

Analytical solving is only possible through the use of simplifying assumptions. As said in introduction, the two most widely used simplifying assumptions are: steady-state (for radical species) and thermal stability of hydroperoxides. If the first assumption seems to be well founded and allows explaining the constancy of the oxidation rate in the first periods of exposure in radio-thermal environments, in contrast, the second one prevents accounting for the auto-acceleration of the oxidation kinetics in the rest of the exposure. This acceleration is due to the decomposition of hydroperoxides above a critical concentration $[\text{POOH}]_F$, which is assumed to be a fraction of the steady concentration of hydroperoxides $[\text{POOH}]_\infty$ reached in the case of "pure" thermal oxidation [2].

Next sections are devoted to the proposal and validation of a sound analytical kinetic model for radio-thermal oxidation taking into account the hydroperoxide decomposition.

3.3. Analytical solving

Denote r_i the global initiation rate of the oxidation reaction. It writes as the sum of the radiochemical and thermal initiation rates:

$$r_i = r_i + 2k_1[\text{POOH}]^2 \quad (22)$$

In all the developments which follow, it will be assumed that oxidation is mainly initiated by the polymer radiolysis throughout

the exposure (i.e. $r_i \gg 2k_1[\text{POOH}]^2$) even if, at long-term, hydroperoxide decomposition is added as additional (but secondary) source of radicals. In other words, only the right part of the domain of intermediary dose rates will be the subject of a kinetic analysis (see Fig. 1).

Applying the steady-state assumption to radical species leads to:

$$\frac{d[\text{P}^\bullet]}{dt} + \frac{d[\text{PO}_2^\bullet]}{dt} = 0 \quad (23)$$

$$\text{i.e. } k_4[\text{P}^\bullet]^2 + k_5[\text{P}^\bullet][\text{PO}_2^\bullet] + k_6[\text{PO}_2^\bullet]^2 - \frac{r_i}{2} = 0 \quad (24)$$

whose positive root is:

$$[\text{P}^\bullet] = \frac{k_5[\text{PO}_2^\bullet]}{2k_4} \left\{ -1 + \left[1 + \psi \left(\frac{[\text{PO}_2^\bullet]_0^2}{[\text{PO}_2^\bullet]^2} - 1 \right) \right]^{1/2} \right\} \quad (25)$$

where ψ is a kinetic parameter quantifying the competition between the different terminations of radical species:

$$\psi = \frac{4k_4k_6}{k_5^2} \quad (26)$$

and $[\text{PO}_2^\bullet]_0$ is the steady concentration of peroxy radicals in oxygen excess:

$$[\text{PO}_2^\bullet]_0 = \left(\frac{r_i}{2k_6} \right)^{1/2} \quad (27)$$

$$\text{i.e. } [\text{PO}_2^\bullet]_0 = \left(\frac{r_i}{2k_6} \right)^{1/2} \left(1 + \frac{2k_1[\text{POOH}]^2}{r_i} \right)^{1/2} \quad (28)$$

$$\text{i.e. } [\text{PO}_2^\bullet]_0 \approx \left(\frac{r_i}{2k_6} \right)^{1/2} \quad (29)$$

As expected, $[\text{PO}_2^\bullet]_0$ takes approximatively the same value as in a "pure" radiochemical environment [39].

If, as shown by Gillen et al. [40], $\psi \ll 1$, then there is a more or less wide range of oxygen concentrations for which:

$$1 \geq \frac{[\text{PO}_2^\bullet]^2}{[\text{PO}_2^\bullet]_0^2} > \frac{5\psi}{1+5\psi} \quad (30)$$

In this range of oxygen concentrations, Eq. (25) can be simplified as follows:

$$[\text{P}^\bullet] = \frac{2k_6}{k_5} ([\text{PO}_2^\bullet]_0 - [\text{PO}_2^\bullet]) \quad (31)$$

If introducing Eq. (31) into Eq. (19) and if applying the steady-state assumption to peroxy radicals, it comes finally:

$$[\text{PO}_2^\bullet] = [\text{PO}_2^\bullet]_0 \frac{\beta C}{1 + \beta C} \quad (32)$$

$$\text{i.e. } [\text{PO}_2^\bullet] = \left(\frac{r_i}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \quad (33)$$

where β^{-1} corresponds to the critical value of the oxygen concentration C_c above which oxygen excess is reached:

$$\beta = \frac{1}{C_c} = \frac{2k_6k_2}{k_5[k_3[\text{PH}] + (2r_1k_6)^{1/2}]} \quad (34)$$

$$\text{i.e. } \beta = \frac{2k_6k_2}{k_5 \left[k_3[\text{PH}] + (2r_1k_6)^{1/2} \left(1 + \frac{2k_1[\text{POOH}]^2}{r_i} \right)^{1/2} \right]} \quad (35)$$

$$\text{i.e. } \beta \approx \frac{2k_6k_2}{k_5[k_3[\text{PH}] + (2r_1k_6)^{1/2}]} \quad (36)$$

As also expected, β takes approximatively the same value as in a "pure" radiochemical environment.

Finally, Eq. (31) leads to:

$$[P^*] = \frac{2k_6}{k_5} [PO_2^*]_0 \left(1 - \frac{\beta C}{1 + \beta C} \right) \quad (37)$$

$$\text{i.e. } [P^*] = \frac{(2r_1 k_6)^{1/2}}{k_5} \left(1 - \frac{\beta C}{1 + \beta C} \right) \quad (38)$$

It is now possible to integrate Eq. (20). The introduction of Eq. (33) into Eq. (20) leads to:

$$\frac{d[POOH]}{dt} = -2k_{1b}[POOH]^2 + k_3[PH] \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \quad (39)$$

This equation can be rewritten such as:

$$\frac{d[POOH]}{[POOH]^2 - \frac{k_3[PH]}{2k_{1b}} \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C}} = -2k_{1b} dt \quad (40)$$

$$\text{i.e. } \text{Ln} \left\{ \frac{\left(\frac{k_3[PH]}{2k_{1b}} \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \right)^{1/2} - [POOH]}{\left(\frac{k_3[PH]}{2k_{1b}} \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \right)^{1/2} + [POOH]} \right\} = -Kt + a \quad (41)$$

$$\text{i.e. } \frac{\left(\frac{k_3[PH]}{2k_{1b}} \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \right)^{1/2} - [POOH]}{\left(\frac{k_3[PH]}{2k_{1b}} \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \right)^{1/2} + [POOH]} = b \text{Exp}(-Kt) \quad (42)$$

$$\text{i.e. } [POOH] = \left(\frac{k_3[PH]}{2k_{1b}} \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \right)^{1/2} \frac{1 - b \text{Exp}(-Kt)}{1 + b \text{Exp}(-Kt)} \quad (43)$$

where a and b are constants to be determined from boundary conditions, and K is a first-order rate constant expressed by:

$$K = 2 \left(2k_3[PH]k_{1b} \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \right)^{1/2} \quad (44)$$

Imposing initial conditions (i.e. $[POOH] = [POOH]_0$ when $t = 0$) to Eq. (43) gives finally:

$$[POOH] = [POOH]_\infty \frac{1 - b \text{Exp}(-Kt)}{1 + b \text{Exp}(-Kt)} \quad (45)$$

$$\text{with } [POOH]_\infty = \left(\frac{k_3[PH]}{2k_{1b}} \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \right)^{1/2} \quad (46)$$

$$\text{and } b = \frac{[POOH]_\infty - [POOH]_0}{[POOH]_\infty + [POOH]_0} \quad (47)$$

For weakly pre-oxidized samples, $[POOH]_\infty \gg [POOH]_0$. Experimental measurements [28,37] show that: $[POOH]_\infty > 10 [POOH]_0$, therefore it can be reasonably considered that: $b \approx 1$.

According to Eq. (45), triggering thermal initiation will completely slows down hydroperoxide build-up during the exposure to finally lead to a maximum concentration of hydroperoxides $[POOH]_\infty$, in the same way as for "pure" thermal oxidation [41]. If this is really the case, it would be interesting to consider this asymptotic value in order to define a more relevant end-of-life criterion in this study, given that this latter now depends on all the exposure conditions (i.e. γ irradiation dose rate, temperature, and oxygen partial pressure).

3.4. Calculation of physico-chemical properties

From these analytical solutions for $[POOH]$, $[PO_2^*]$, and $[P^*]$, it is now possible to calculate several key physico-chemical properties, from a practical point of view, because these properties can be checked experimentally. Undoubtedly, hydroperoxide concentration $[POOH]$ and oxygen consumption Q are the most relevant properties because their mathematical expressions can be derived from the radio-thermal oxidation mechanistic scheme without using any additional adjustable parameter.

Recall that the rate of oxygen consumption writes:

$$r(C) = -\frac{d[O_2]}{dt} = k_2 C [P^*] - k_6 [PO_2^*]^2 \quad (48)$$

$$\text{i.e. } r(C) = \frac{2k_6 k_2 C}{k_5} ([PO_2^*]_0 - [PO_2^*]) - k_6 [PO_2^*]^2 \quad (49)$$

$$\text{i.e. } r(C) = k_3 [PH] \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} + r_1 \frac{\beta C}{1 + \beta C} \left(1 - \frac{\beta C}{2(1 + \beta C)} \right) \quad (50)$$

It can be easily checked that, when $\beta C \gg 1$, Eq. (50) leads to the expression of the rate of oxygen consumption in oxygen excess [1]:

$$r_0 = k_3 [PH] \left(\frac{r_1}{2k_6} \right)^{1/2} + \frac{r_1}{2} \quad (51)$$

The integration of Eq. (50) between 0 and t gives access to the concentration of oxygen consumed by the chemical reaction Q:

$$Q = \int_0^t r(C) dt \quad (52)$$

$$\text{i.e. } Q = \left[k_3 [PH] \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} + r_1 \frac{\beta C}{1 + \beta C} \left(1 - \frac{\beta C}{2(1 + \beta C)} \right) \right] t \quad (53)$$

Other key physico-chemical property in the literature are the concentrations of carbonyl products $[P=O]$ and chain scissions S. In the radio-thermal oxidation mechanistic scheme, carbonyl products can only be formed in thermal initiation (1T) and termination (6) through specific chemical events (e.g. β scission, disproportionation, etc.) which generally compete with many other chemical events, in particular with hydrogen abstraction which gives alcohols. Consequently, the calculation of the formation rate of these oxidation products requires the use of adjustable parameters:

$$\frac{d[P=O]}{dt} = \gamma_{1CO} k_{1b} [POOH]^2 + \gamma_{6CO} k_6 [PO_2^*]^2 \quad (54)$$

where γ_{1CO} and γ_{6CO} are the respective formation yields of carbonyl products in thermal initiation (1T) and termination (6), see section 3.1.

If replacing $[POOH]$ and $[PO_2^*]$ by their analytical solutions in Eq. (54), it comes:

$$\frac{d[P=O]}{dt} = \gamma_{1CO} k_{1b} [POOH]_\infty^2 \left(\frac{1 - b \text{Exp}(-Kt)}{1 + b \text{Exp}(-Kt)} \right)^2 + \gamma_{6CO} \frac{r_1}{2} \left(\frac{\beta C}{1 + \beta C} \right)^2 \quad (55)$$

$$\text{i.e. } \frac{d[P=O]}{dt} = \gamma_{1CO} \frac{k_3 [PH]}{2} \left(\frac{r_1}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \left(\frac{1 - b \text{Exp}(-Kt)}{1 + b \text{Exp}(-Kt)} \right)^2 + \gamma_{6CO} \frac{r_1}{2} \left(\frac{\beta C}{1 + \beta C} \right)^2 \quad (56)$$

$$\begin{aligned} \text{i.e. } \frac{d[P=O]}{dt} &= \gamma_{1CO} \frac{k_3[PH]}{2} \left(\frac{r_i}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \\ &+ \gamma_{6CO} \frac{r_i}{2} \left(\frac{\beta C}{1 + \beta C} \right)^2 - \gamma_{1CO} k_3 [PH] \left(\frac{r_i}{2k_6} \right)^{1/2} \\ &\times \frac{\beta C}{1 + \beta C} \frac{2b \text{Exp}(-Kt)}{(1 + b \text{Exp}(-Kt))^2} \end{aligned} \quad (57)$$

The integration of Eq. (57) between 0 and t leads finally to:

$$\begin{aligned} [P=O] &= \left[\gamma_{1CO} \frac{k_3[PH]}{2} \left(\frac{r_i}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} + \gamma_{6CO} \frac{r_i}{2} \left(\frac{\beta C}{1 + \beta C} \right)^2 \right] t \\ &+ 2\gamma_{1CO} \frac{k_3[PH]}{K} \left(\frac{r_i}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \\ &\times \left(\frac{1}{1 + b \text{Exp}(-Kt)} - \frac{1}{1 + b} \right) \end{aligned} \quad (58)$$

Similarly, chain scissions can only be formed in thermal initiation (1T) and termination (6) through β scissions. Consequently, their mathematical expressions are identical in form to those previously established for carbonyl products, only their formation yields are different:

$$\frac{dS}{dt} = \gamma_{1S} k_{1b} [POOH]^2 + \gamma_{6S} k_6 [PO_2]^2 \quad (59)$$

$$\begin{aligned} \text{i.e. } \frac{dS}{dt} &= \gamma_{1S} \frac{k_3[PH]}{2} \left(\frac{r_i}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} + \gamma_{6S} \frac{r_i}{2} \left(\frac{\beta C}{1 + \beta C} \right)^2 \\ &- \gamma_{1S} k_3 [PH] \left(\frac{r_i}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \frac{2b \text{Exp}(-Kt)}{(1 + b \text{Exp}(-Kt))^2} \end{aligned} \quad (60)$$

and

$$\begin{aligned} S &= \left[\gamma_{1S} \frac{k_3[PH]}{2} \left(\frac{r_i}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} + \gamma_{6S} \frac{r_i}{2} \left(\frac{\beta C}{1 + \beta C} \right)^2 \right] t \\ &+ 2\gamma_{1S} \frac{k_3[PH]}{K} \left(\frac{r_i}{2k_6} \right)^{1/2} \frac{\beta C}{1 + \beta C} \left(\frac{1}{1 + b \text{Exp}(-Kt)} - \frac{1}{1 + b} \right) \end{aligned} \quad (61)$$

where γ_{1S} and γ_{6S} are the respective yields of chain scissions in thermal initiation (1T) and termination (6), see section 3.1.

In the case of PE oxidation, it is clear that embrittlement results essentially from chain scissions [18,19]. Neglecting, in a first approach, an eventual chemical cross-linking process, the changes in the concentration of elastically active chains in Si-XLPE throughout the exposure can simply be calculated as follows [42]:

$$\nu = \nu_0 - S \quad (62)$$

where ν_0 and ν are the respective concentrations of elastically active chains before and after radio-thermal ageing. Swelling measurements and rheometry in solid state (DMTA) allowed showing that $\nu_0 = 0.12 \pm 0.02 \text{ mol.L}^{-1}$ for Si-XLPE.

3.5. Lifetime prediction

Embrittlement is expected to occur at very low conversion ratio of the chain scission process. On one side, the corresponding critical ratio S_F is related through Eq. (62) to a critical concentration of elastically active chains ν_F to be determined experimentally. On the other side, as explained at the end of section 3.2, S_F would also be related to a critical hydroperoxide concentration $[POOH]_F$ corresponding to the onset of the rapid auto-acceleration of the oxidation reaction when triggering thermal initiation. As shown in

a previous publication [2], the definition of this last structural end-of-life criterion is quite simple:

$$[POOH]_F = \frac{[POOH]_\infty}{q} \quad (63)$$

where q is an arbitrary criterion, higher than unity but not so far, to be determined experimentally.

In a first approach, it was decided to use $[POOH]_F$ for predicting the lifetime t_F of the insulating material. It was thus assumed that $t = t_F$ when $[POOH] = [POOH]_F$.

In this case, introducing Eq. (63) into Eq. (45) leads to:

$$q = \frac{1 + b \text{Exp}(-Kt_F)}{1 - b \text{Exp}(-Kt_F)} \quad (64)$$

$$\text{i.e. } \text{Exp}(-Kt_F) = \frac{q - 1}{b(1 + q)} \quad (65)$$

$$\text{i.e. } t_F = \frac{1}{K} \text{Ln} \left[\frac{b(1 + q)}{q - 1} \right] \quad (66)$$

4. Results and discussion

4.1. Long-term behavior of PE insulating materials

Fig. 8 shows the changes in the elongation at break ε_R of Si-XLPE in air under the three dose rates under study at low temperature close to ambient. As expected, γ irradiation induces a catastrophic embrittlement of Si-XLPE. The horizontal dashed line plotted for $\varepsilon_R = \varepsilon_F = 50\%$ was used to graphically determine the lifetime under these various exposure conditions. It was found that t_F is about 30, 50, and 260 days under the highest (i.e. 400 Gy.h⁻¹), medium (77.8 Gy.h⁻¹), and lowest dose rates (8.5 Gy.h⁻¹), respectively.

These t_F values were plotted as a function of dose rate in Fig. 9 where they compared to those previously reported in the literature for other types of PE [1]. It should be note that the t_F value determined at 25°C for “pure” thermal oxidation is also plotted, but for a dose rate of $1.5 \times 10^{-10} \text{ Gy.s}^{-1}$ corresponding to natural radioactivity.

The almost universal character of the long-term behavior of PE is confirmed. It appears to be completely independent of the chemistry (grafted with chemical group or not), macromolecular

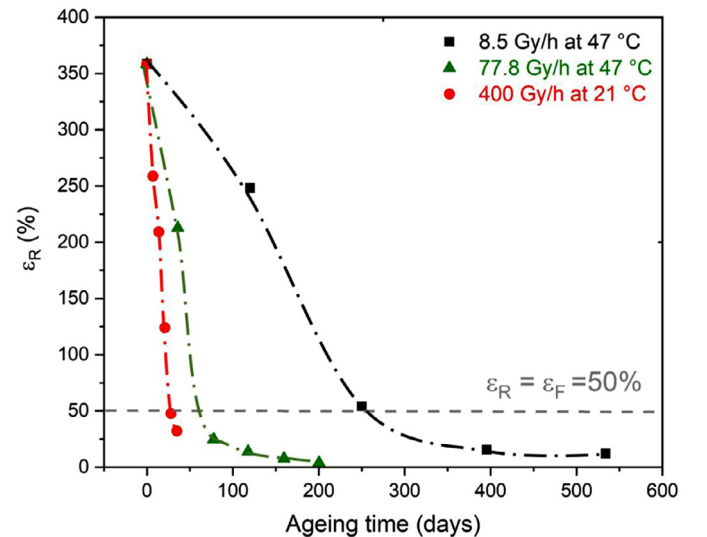


Fig. 8. Changes in the elongation at break of Si-XLPE in air under the three dose rates under study at low temperature close to ambient.

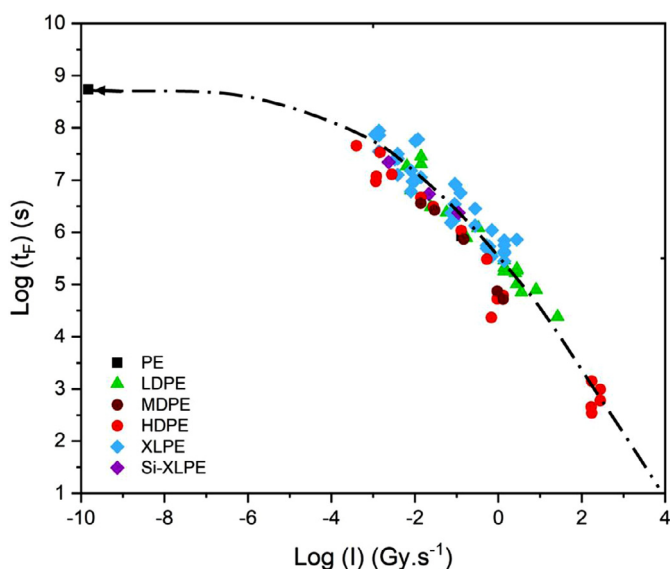


Fig. 9. Plot in logarithm co-ordinates of lifetime t_F versus γ dose rate I in air at room temperature for unstabilised and unfilled linear PE. Comparison of the t_F values determined for Si-XLPE in this study with those reported for various other types of linear (i.e. LDPE, MDPE and HDPE) and chemically cross-linked LDPE (i.e. XLPE) in the literature [1]. The general trend of the curve is shown using a dotted line.

Table 2

Experimental values of the lifetime and steady and critical hydroperoxide concentrations determined experimentally for Si-XLPE in air in the various radio-thermal environments under study.

I (Gy.h ⁻¹)	400	77.8	8.5	6.0
T (°C)	21	47	47	86
t_F (days)	30	50	260	–
$[POOH]_{\infty}$ (mol.L ⁻¹)	2.4	3.1×10^{-1}	1.7×10^{-1}	4.1×10^{-2}
$[POOH]_F$ (mol.L ⁻¹)	1.5×10^{-1}	1.8×10^{-1}	1.6×10^{-1}	–

structure (i.e. linear, branched or chemically cross-linked) and crystallinity of PE. It can be clearly seen that all the t_F values determined for Si-XLPE are well positioned on the master curve. In addition, these values are located in the right part of the domain of intermediary dose rates where radio-thermal oxidation proceeds. Consequently, all the physicochemical data determined in this study will be used to check the reliability of the analytical equations established in previous sections.

4.2. Reliability of the analytical kinetic model

In a first stage, Eqs. (45), (58), and (62) were used to respectively simulate the changes in concentration of hydroperoxides [POOH], carbonyls [P=O], and elastically active chains ν of Si-XLPE in air under the three dose rates under study (i.e. 8.5, 77.8, and 400 Gy.h⁻¹) at low temperature close to ambient. Simulations are reported in Fig. 10, Fig. 11, and Fig. 12, respectively. These results call for the following comments.

At short term, [POOH] increases with time until reaching an asymptotic value $[POOH]_{\infty}$ which depends on the radio-thermal exposure conditions. All this first part of the kinetic curves is satisfyingly simulated with Eq. (45), the resulting values of $[POOH]_{\infty}$ have been reported in Table 2. As expected, $[POOH]_{\infty}$ is an increasing function of irradiation dose I , but a decreasing function of temperature T . Indeed, this asymptote is the consequence of the sudden acceleration of the thermal initiation (1T), which counteracts completely the build-up of hydroperoxides formed in the propagation step of the oxidation reaction.

At longer term, for certain radio-thermal exposure conditions (e.g. under 77.8 Gy.h⁻¹ at 47°C), a beginning of decrease in [POOH]

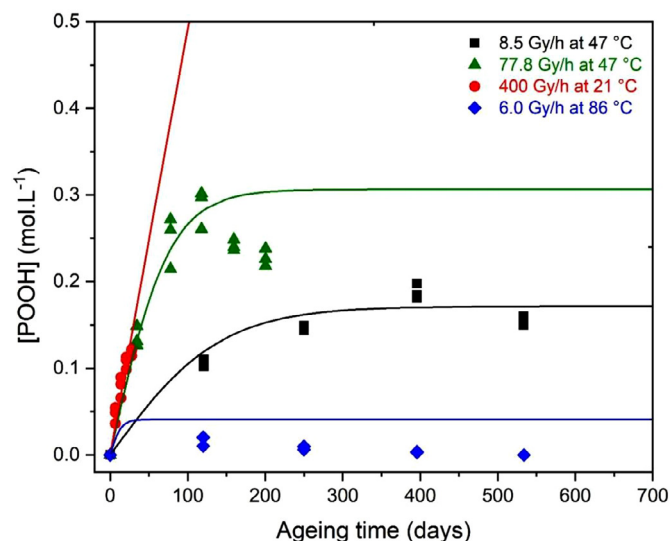


Fig. 10. Changes in the hydroperoxide concentration of XLPE in air in the various radio-thermal environments under study. Comparison between simulation with Eq. (45) (solid lines) and experimental data (symbols).

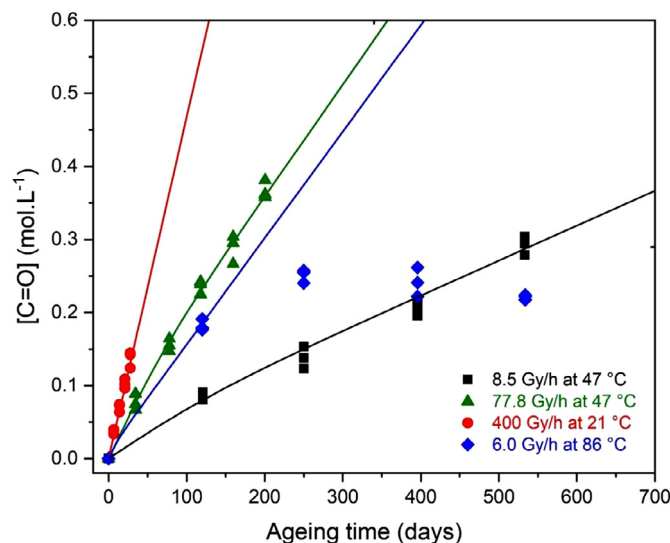


Fig. 11. Changes in the carbonyl concentration of XLPE in air in the various radio-thermal environments under study. Comparison between simulation with Eq. (58) (solid lines) and experimental data (symbols).

due to the depletion of oxidation sites (PH) can be observed, but this second part of the kinetic curves cannot be correctly described by Eq. (45) because this latter is based on the assumption of low conversion ratios.

$[POOH]_F$ was determined graphically in Fig. 10 (i.e. when $t = t_F$) for the various radio-thermal exposure conditions under study. Its values are also reported in Table 2. It was found that $[POOH]_F$ is independent of the radio-thermal exposure conditions, which completely validates the choice of this critical quantity as structural end-of-life criterion in this study. The average value of $\approx 1.6 \times 10^{-1}$ mol.L⁻¹ was chosen for predicting the lifetime of Si-XLPE. This is exactly the same $[POOH]_F$ value that was retained in a previous publication for other types of (linear and chemically cross-linked) PE, but without making real demonstration, just based on mechanistic considerations [2].

In the early periods of exposure, $[P=O]$ and ν vary linearly with time (absence of induction period). Then, when thermal initiation gains in importance, a small loss in linearity (progressive slow-

Table 3

Values of the kinetic parameters used for modeling the oxidation kinetics of Si-XLPE in the various radio-thermal environments under study. Papers in which the same orders of magnitude were previously reported are referenced in the last column on the right.

T (°C)	21	47	47	86	
I (Gy.h ⁻¹)	400	77.8	8.5	6.0	References
G _i	8	8	8	8	[1,2,39]
k _{1b} (L.mol ⁻¹ .s ⁻¹)	5.0 × 10 ⁻⁹	2.5 × 10 ⁻⁷	2.4 × 10 ⁻⁷	1.0 × 10 ⁻⁵	[2,3,20,36,39,43,44]
k ₂ (L.mol ⁻¹ .s ⁻¹)	10 ⁸	10 ⁸	10 ⁸	10 ⁸	[2,3,20,36,39,43,44]
k ₃ (L.mol ⁻¹ .s ⁻¹)	1.6 × 10 ⁻³	1.9 × 10 ⁻²	1.9 × 10 ⁻²	3.6 × 10 ⁻¹	[2,3,20,36,39,43,44]
k ₄ (L.mol ⁻¹ .s ⁻¹)	8.0 × 10 ¹¹	8.0 × 10 ¹¹	8.0 × 10 ¹¹	8.0 × 10 ¹¹	[2,3,20,43,44]
k ₅ (L.mol ⁻¹ .s ⁻¹)	1.2 × 10 ¹⁰	7.0 × 10 ¹⁰	9.0 × 10 ¹⁰	2.4 × 10 ¹¹	This study
k ₆ (L.mol ⁻¹ .s ⁻¹)	5.0 × 10 ⁴	1.0 × 10 ⁶	2.0 × 10 ⁶	6.0 × 10 ⁷	[2,3,20,36,39,43,44]
γ ₁₅ (%)	90	42	52	51	[20,43,44]
γ ₆₅ (%)	90	42	52	51	[20,43,44]
γ _{1CO} (%)	90	70	75	100	[20,43,44]
γ _{6CO} (%)	90	70	75	100	[20,43,44]

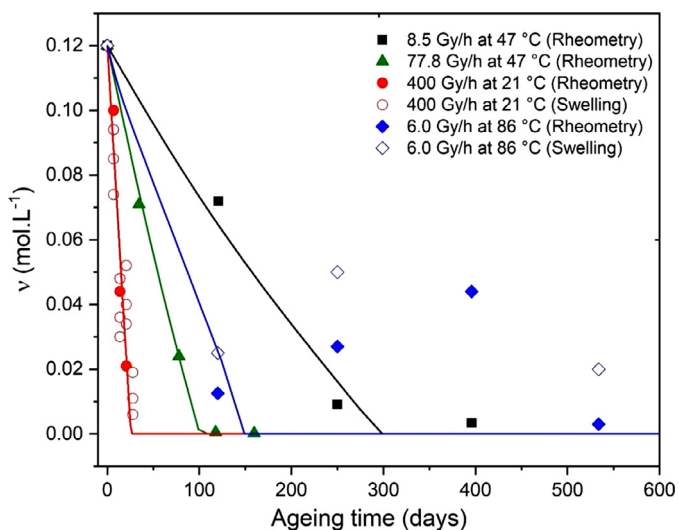


Fig. 12. Changes in the concentration of elastically active chains of XLPE in air in the various radio-thermal environments under study. Comparison between simulation with Eq. (62) (solid lines) and experimental data determined by swelling (solid symbols) and rheometry (empty symbols).

down) can be observed for both physico-chemical quantities. The whole kinetic curves are satisfyingly simulated with Eqs. (58) and (62), respectively.

In a second stage, Eqs. (45), (58), and (62) were used to respectively simulate the physico-chemical changes of Si-XLPE in air under 6.0 Gy.h⁻¹ at 86°C. Simulations are also reported in Fig. 10, Fig. 11, and Fig. 12, respectively. It can be clearly seen that the impact of the thermal initiation on the global oxidation kinetics is much more important at 86°C than at 47°C. Unfortunately, the three analytical equations are not able to accurately account for the whole kinetic curves, presumably due to the starting assumption, which was deduced from Eq. (22). In order to reach this goal, it would be necessary to take into account that thermal initiation could be of the same order of magnitude as (if not greater than) radiochemical initiation. This further kinetic analysis will be the subject of a future publication. At this stage of investigations, it can be seen that the three analytical equations allow accounting for the general trends of the oxidation kinetics under these critical exposure conditions (i.e. under low dose rate at moderate temperature) where, obviously, a stronger coupling between radiochemical and thermal oxidations operates.

The values of the kinetic parameters used for all these simulations, i.e. radiochemical yield G_i, rate constants k_j (with j = 1, ..., 6) and the formation yields γ_{1CO} and γ_{6CO} of carbonyls products,

Table 4

Values of the rate constant K and criterion q used for predicting the lifetime of Si-XLPE in air under the three dose rates under study at low temperature close to ambient.

I (Gy.h ⁻¹)	400	77.8	8.5
T (°C)	21	47	47
K (s ⁻¹)	4.8 × 10 ⁻⁸	3.1 × 10 ⁻⁷	1.6 × 10 ⁻⁷
q	15.0	1.9	1.1
t _f (days)	32	43	246

and γ₁₅ and γ₆₅ of chain scissions, have been reported in Table 3. These values call for the following comments.

The orders of magnitude previously determined on linear PE (i.e. LDPE, MDPE and HDPE) for almost all these kinetic parameters were kept for these simulations [1-3,20,36,39,43,44]. Only k₅ was slightly reduced at low temperature close to ambient (i.e. at 47 and 21°C) because it had already been proposed to allocate a small activation energy to this termination rate constant in order to better simulate the oxidation kinetics initiated by chlorine disinfectants in PE pipes carrying drinking water at ambient temperature [20,43,44]. Thus, the importance of this activation energy is completely checked in this study. In addition, if considering all the k₅ values published to date for PE [2,3,20,43,44], it is found that its activation energy would be around 14 kJ.mol⁻¹.

As there was no notable change in the values of all the kinetic parameters, it can be concluded that the reliability of the analytical kinetic model was checked in this study. Its great heuristic value has been demonstrated in Fig. 10, Fig. 11, and Fig. 12, even if some improvements are still possible, in particular at the highest temperature (86°C).

The last stage of this study was to check the reliability of the last Eq. (66). Rate constant K and criterion q were determined with Eq. (44) and (63), respectively. Then, their values were introduced into Eq. (66) in order to predict the lifetime t_f of Si-XLPE in air under the three dose rates under study at low temperature close to ambient. The resulting t_f values are reported in Table 4. It can be seen that they are of the same order of magnitude than those determined experimentally, thus confirming in turn the great heuristic value of Eq. (66).

5. Conclusion

A new analytical model was developed for predicting the radio-thermal oxidation kinetics and the lifetime of electric cable insulations made of unfilled and unstabilised Si-XLPE in nuclear power plants. This model was derived from a mechanistic scheme in which the oxidation reaction is initiated both by the polymer radiolysis and the thermal decomposition of hydroperoxides, with-

out making the usual assumption of thermal stability of hydroperoxides. After an initial period where the oxidation kinetics occurs with a constant rate, this model allows also predicting the auto-acceleration of the oxidation kinetics when hydroperoxide decomposition is no longer negligible. A structural end-of-life criterion: $[POOH]_F \approx 1.6 \times 10^{-1} \text{ mol. L}^{-1}$, corresponding to the onset of this auto-acceleration, was proposed for predicting the lifetime of Si-XLPE in the various radio-thermal environments under study.

A satisfying agreement was obtained between theory and experiments as long as thermal initiation remains a secondary source of radicals relatively to radiochemical initiation, i.e. under the three dose rates under study (from 8.5 to 400 Gy.h⁻¹) at low temperature close to ambient (i.e. 47 and 21°C). However, under the lowest dose rate (i.e. 6.0 Gy.h⁻¹) at the highest temperature (86°C), the analytical model does not correctly simulate the changes in concentration of the three physico-chemical properties investigated experimentally, i.e. hydroperoxides, carbonyls products and elastically active chains. In fact, in these critical radio-thermal exposure conditions, it would be necessary to take into account that thermal initiation could be of the same order of magnitude as (if not greater than) the radiochemical initiation. This further kinetic analysis will be the subject of a future publication. The final objective will be the proposal of an analytical equation allowing to simulate the whole master curve: $t_F = f(I)$ reported in Fig. 9, whatever the dose rate.

However, industrial polymers are rarely used pure, in particular they contain antioxidants. That is the reason why the next longer term objective will be to extend the kinetic model to Si-XLPE stabilized by the most common families of antioxidants, namely hydroperoxide decomposers (i.e. sulphides and phosphites) and radical scavengers (i.e. hindered phenols and secondary aromatic amines). To reach this objective, the main action mechanisms of these antioxidants will be introduced into the radio-thermal oxidation mechanistic scheme in order to derive a new system of differential equations to be solved. It should be pointed out that the main role of these antioxidants is to increase the induction period. Consequently, when these latter have been completely consumed, the Si-XLPE matrix should oxidize according to almost the same kinetics and thus, the structural end-of-life criterion established in this study should remain valid for predicting its lifetime.

Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

CRediT authorship contribution statement

Sarah Hettal: Investigation, Writing - original draft. **Sébastien Roland:** Validation, Writing - review & editing, Supervision. **Konsta Sipilä:** Writing - review & editing, Supervision, Project administration. **Harri Joki:** Investigation. **Xavier Colin:** Conceptualization, Methodology, Software, Validation, Writing - review & editing, Supervision.

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