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Comparisons Between Numerical Calculations and Measurements in the Vaned Diffusor of SHF Impeller

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Abstract: The paper presents analysis of the performance and the internal flow behaviour in the vaned diffusor of a radial flow pump using PIV (Particle Image Velocimetry) and pressure probe traverses. PIV measurements have already been performed at mid height inside one diffusor channel passage for a given speed of rotation and various mass flow rates. These results have been already presented in several previous communications. New experiments have been performed using a 3 holes pressure probe traverses from hub to shroud diffusor width at different radial locations between the two diffusor geometrical throats. Numerical simulations are also realized with the commercial codes Star CCM+ 7.02.011 and CFX. Frozen rotor and fully unsteady calculations of the whole pump have been performed. Comparisons between numerical results, previous experimental PIV results and new probe traverses one’s are presented and discussed for one mass flow rate. In this respect, a first attempt to take into account fluid leakages between the rotating and fixed part of the pump has been checked since it may affect the real flow structure inside the diffuser.

Keywords: Vaned diffuser, Radial flow pump, unsteady flow, numerical calculations, Particle Image Velocimetry, three-holes probe.

1. Introduction

Knowledge improvement of rotor stator interactions in radial flow pumps has already been a favorite research theme for a large number of laboratories (references [1] to [5]). A great number of rotor stator interactions analyses based on modeling activities have been realized, but experimental analyses remain scarce. Thanks to nowadays computer performances and the evolution of PIV techniques (ref. [6] and [7]), it has been possible to measure velocity fields within the impeller and the diffuser of a radial flow pump. Tests have been performed in air with the so-called ‘SHF’ impeller for three kinds of diffusers: one vaneless and two differently vaned. Results obtained with the vaneless diffuser have already been presented in previous symposia and published (see ref. [8] to [17]).

In numerical simulation, two aspects have to be considered. The first one concerns the governing equations which are solved in the model: two kinds of numerical calculations are currently used in turbomachinery (i-frozen rotor calculations, ii-unsteady calculations). The second aspect concerns the geometrical model. Some geometrical simplifications are currently used. For example, the leakage flows are often neglected. It is obvious that a complex model (fully unsteady, with leakage flows) will be more time consuming but closer to the real physics. The aim of the paper is to test several more or less sophisticated models to evaluate their ability to predict properly the pump performance and the internal flow behaviour. To do so, numerical results are compared to experimental
Comparisons Between Numerical Calculations and Measurements in the Vaned Diffuser of SHF Impeller

The existing database has been completed by pressure probe measurements for a complete flow performance analysis. These measurements are compared to several numerical calculations in the present paper. The experimental set-up and the numerical models are presented in the first parts of the paper. The comparisons and analysis of their results (global performance as well as velocity profile) within the diffuser are then discussed.

2. Experiments

2.1 Test and apparatus

Test pump model and PIV measurements conditions have been already described [8-10] and main geometrical characteristics as well as operating conditions are given in table 1. This set-up allows the existence of a “positive” leakage flow going into the gap between impeller outlet section and vaned diffusor inlet section which is specific to the experimental set-up. This is due to the fact that the pump outlet pressure corresponds to atmospheric conditions.

Table 1  Pump parameters

<table>
<thead>
<tr>
<th>Entry duct</th>
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<tr>
<td>Diameter</td>
<td>0.14113 m</td>
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<tr>
<td>length</td>
<td>0.25 m</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th>Impeller</th>
<th></th>
</tr>
</thead>
<tbody>
<tr>
<td>Inlet radius</td>
<td>0.14113 m</td>
</tr>
<tr>
<td>Outlet radius</td>
<td>0.2549 m</td>
</tr>
<tr>
<td>Number of blades</td>
<td>Z_i=7</td>
</tr>
<tr>
<td>Outlet height</td>
<td>0.04 m</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th>Diffusor</th>
<th></th>
</tr>
</thead>
<tbody>
<tr>
<td>Inlet radius</td>
<td>0.2549 m</td>
</tr>
<tr>
<td>Outlet radius</td>
<td>0.44 m</td>
</tr>
<tr>
<td>Number of vanes</td>
<td>Z_d=8</td>
</tr>
<tr>
<td>Height</td>
<td>0.04 m</td>
</tr>
<tr>
<td>Rotation speed</td>
<td>1710 rpm</td>
</tr>
</tbody>
</table>

2.2 Three–holes probe

A three holes probe (see Fig. 1) has been used to make hub to shroud traverses in the same planes where wall static pressures are measured. Using a specific calibration one can get total pressure, static pressure, absolute velocity and its two components in radial and tangential direction.

2.3. General flow conditions

All types of measurements have been performed for several mass flow rates. But, results presented in this paper refer mainly to one mass flow rate corresponding to the design point of the vaned diffusor Q/Q_{ni}=0.766, which is different from impeller design mass flow rate Q/Q_{ni}=1. This vaned diffusor design was chosen in order to allow an enlarge pump performance characteristic curve for low mass flow rates. In order to well represent the flow field, twenty-three location positions are defined as it can be seen in Fig. 2. For each location, ten axial positions are registered (b*=0.125, 0.2, 0.25, 0.375, 0.5, 0.625, 0.75, 0.875, 0.925, 0.975). The present analysis focuses only on locations 19 to 23 in the blade to blade channel of the diffuser.

Fig. 1 Picture and sketch of the three-holes pressure probe

Fig. 2 Diffusor measurement locations for probe traverse and unsteady calculations
3. Calculations

Calculations were performed using Star CCM+ code and the results are compared to already published ones obtained with CFX [8-10].

3.1. Frozen rotor calculations (Star CCM+)

Three-dimensional steady incompressible Reynolds averaged Navier–Stokes equations are solved. The SST k-ω turbulence model is used [19]. The calculation domain was divided into three zones: the inlet zone, the impeller zone and the diffusor zone.

The boundary condition at the inlet consisted of a mass flow rate (0.305 kg/s; Q/Qni = 0.766). The boundary condition at the outlet was the atmospheric pressure (0 Pa). The fluid (air) was considered incompressible at a constant temperature of 20ºC.

A polyhedral mesh with prism layers is used for all calculations (5 prism layers for a total prism layer thickness of 1 mm). The target size is 3 mm and the minimum size 0.5 mm. The size of the grid is about 3.10⁶ cells. A sketch of local mesh is given in Fig. 3.

Line probes are plotted in the whole domain: Each probe is duplicated seven times in order to obtain all parameters (pressure, total pressure, radial, tangential and axial velocity and velocity magnitude) for seven different relative positions of the diffusor comparatively to the impeller (see Fig. 4). These ones are equal to 2πn/(1/Zi-1/Zd) with n=0,1, 2, 3, 4, 5, and 6.

3.2. Unsteady calculations

3.2.1. Star CCM+

The unsteady calculations were realized with the same mesh. The convergence criteria are less than 1.e-4. Maximum values of y+ are around 15, and mainly below 9 in the whole computational domain as can be seen in Fig.5.

A time step of 4.87e-5 s, corresponding to an angular rotation of 0.5°, has been chosen.

3.2.2. CFX

In order to investigate the influence of leakage flows, numerical results obtained using CFX are also presented. G. Cavazzini et al [8-10] proposed two kinds of unsteady simulations, with and without leakage flows.

As regards the simulations considering the leakage flows, the seal systems both at the impeller inlet and outlet were modelled (Figures 6, 7). At the labyrinth of the seal system close to the impeller inlet the mass flow rates determined from the experimental data was prescribed, assuming stochastic fluctuations of the velocities with 5 %
Comparisons Between Numerical Calculations and Measurements in the Vaned Diffusor of SHF Impeller

free-stream turbulence intensity. At the impeller outlet the leakage mass flow rate was controlled by the known experimental pressure in the large plenums upstream the labyrinth.

Much more details on model, mesh and numerical simulations are reported in ref [10].

Fig. 6 – Seal system at the impeller inlet

Fig. 7 – Seal system at the impeller outlet

4. Results and comparison

4.1. Global performances

Fig 8 presents global results of the performance of the pump. In this figure \( Q^* \) is defined as follows:

\[
Q^* = \frac{Q}{Q_*}
\]  

(1)

Non-dimensional calculated total head is defined by equation (2) and non-dimensional isentropic head by equation (3)

\[
\psi_n = \frac{C\Omega}{Q_* U_{2}^2/2}
\]  

(2)

\[
\psi_s = \frac{dp_r}{\rho U_{2}^2/2}
\]  

(3)

Global results of total theoretical head pump are in good agreement between one dimensional approach and frozen rotor calculation. Unsteady calculation results, performed for the non-dimensional mass flow rate \( Q^* = 0.766 \), give the same level, as it can be seen in Fig. 8.

Comparison between real total head curve obtained by calculations and experiments are also given in the same figure. One has to note that pump mass flow rates are higher than impeller mass flow rates due to “positive” leakage flow. This is the reason why the experimental head pump is higher than the numerical total pump head coefficient.

4.2. Hub-to-shroud local flow characteristics.

The results obtained at location 19 has been chosen to compare frozen rotor and unsteady calculations. Numerical results are shown in Fig. 9 and Fig. 10 respectively for frozen rotor and unsteady calculations. Radial velocity distributions depend on the relative impeller blade position compared with vaned diffusor ones as shown in Fig. 9, whereas only time acts for unsteady results. Unsteady calculations give smoother hub to shroud gradients and less difference between
Comparisons Between Numerical Calculations and Measurements in the Vaned Diffusor of SHF Impeller

blade to blade results.

Local numerical results of frozen rotor calculations without leakage (“CCM+ 1”), unsteady calculations without leakage (obtained from two different codes: “CCM+ 2” and “CFX 1”), and unsteady calculations with leakage (obtained only from CFX [10]: “CFX 2”) are presented in Fig.11 to 14. PIV measurements (“PIV”), wall static pressure (“ps”) and three-holes pressure probes ("probes") results are plotted on the same figures.

The influence of leakage can be easily seen when comparing numerical results obtained without and with leakage. Experimental results are in good agreement with calculations obtained with leakage. This can be mainly seen on radial velocity distribution shown in Fig. 11. These results must be more deeply analysed, in particular PIV ones; they strongly depend on impeller position and only time-averaged values are presented here. Pressure probe results are also depending on unsteady effects and this has to be taken into account for further data reduction analysis.

The influence of leakage can also be seen, looking at pressure distributions in the diffusor. Static pressure recovery is better with leakage (Fig.14). Total pressure probe results show higher levels compared with calculations without leakage and also show that leakage effects acts mainly near shroud section preventing separation in location 22 and 23.

5. Conclusions
Comparisons between numerical frozen rotor assumption and fully unsteady calculations show that global performances can be obtained with frozen rotor assumption but that local values must be analysed with unsteady results even in the vaned diffusor far from impeller outlet section.

Results issued from different numerical approaches and experiments coming from PIV and pressure probe have been presented and compared. The pump model configuration allows leakage to be an important parameter that has to be taken into account in order to make good comparisons between numerical and experiments.

Acknowledgements
The authors wish to thank Region Nord- Pas de Calais and CNRS for their financial support in the frame of the CISIT program.
Comparisons Between Numerical Calculations and Measurements in the Vaned Diffusor of SHF Impeller

Fig. 11a to 11e hub-to-shroud non-dimensional radial velocity $v_{r*}$ for different locations inside the diffusor channel

Fig. 12a to 12e hub-to-shroud non-dimensional radial velocity $v_{u*}$ for different locations inside the diffusor channel
Comparisons Between Numerical Calculations and Measurements in the Vaned Diffusor of SHF Impeller

Fig. 13a to 13e hub-to-shroud non-dimensional pressure $\Psi_s$ for different locations inside the diffusor channel

Fig. 14a to 14e hub-to-shroud non-dimensional total pressure $\Psi_t$ for different locations inside the diffusor channel
Nomenclature

b  Impeller or diffuser width (m)
C  Impeller moment (mN)
H  Total pump head (m)
N  Rotation speed (rpm)
p  Pressure (Pa)
\( p_t \)  Total pressure (Pa)
Q  Volume flow rate in impeller (m³/s)
\( Q_{in} \)  Mass flow rate in impeller (kg/s)
\( Q_k \)  Nominal volume flow rate (m³/s)
R_i  Radius of section i (m)
v_r  Radial velocity (m/s)
v_r=U_2  Non-dimensional radial velocity (-)
v_\theta  Tangential velocity (m/s)
v_\theta=U_2  Non-dimensional tangential velocity (-)
v_z  Axial velocity (m/s)
v_z=U_2  Non-dimensional axial velocity (-)
U_2=\Omega R_2  Frame velocity at impeller outlet (m/s)
Z_b  Number of impeller blades = 7 (-)
Z_d  Number of diffusor blades = 8 (-)
\rho  Air density (kg/m³)
\Theta  Azimuthal vane blade angle (°)
\Omega  Angular Velocity (rad/s)
\Psi_{ti}  Non-dimensional total pump head (-)
\Psi_{ti}  Non-dimensional isentropic head (-)

Index

1  Impeller inlet
2  Impeller outlet
3  Diffusor outlet
c  calculated
d  diffusor
i  Impeller or isentropic

References

Comparisons Between Numerical Calculations and Measurements in the Vaned Diffusor of SHF Impeller


