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# Parts internal structure definition using non-uniform patterned lattice optimization for mass reduction in additive manufacturing

Laurent Chougrani<sup>1</sup> · Jean-Philippe Pernot<sup>2</sup> · Philippe Véron<sup>2</sup> · Stéphane Abed<sup>1</sup>

## Abstract

Today, being able to generate and produce shapes that fit mechanical and functional requirements and having as low as possible mass is crucial for aerospace and automotive applications. Besides, the rise of new additive manufacturing technologies has widened the possibilities for designing and producing complex shapes and internal structures. However, current models, methods and tools still represent a limitation to that new horizon of printable shapes. This paper addresses the way internal lattice structures can be generated and optimized to reduce the mass of a product. A new framework is introduced that allows the modeling and optimization of non-uniform patterned lattice structures. Using non-uniform structures, additional degrees of freedom are introduced and allow the definition of a wide variety of shapes which can better fit the requirements. First, a non-uniform patterned lattice structure is generated using the results of an initial finite element analysis. This initial structure is then optimized while iteratively removing the beams considered as useless with respect to a user-specified mechanical criteria. At each iteration, the lattice structure is sent to a finite element solver that returns the von Mises stress map used to drive the simplification process. Here, the simulations are performed on the wireframe lattice structures to speed up the optimization loops. Once this process is completed, the final structure is no longer fully patterned, but it is re-organized to reduce the mass while satisfying the mechanical criteria. This approach is illustrated with examples coming from our prototype software.

**Keywords** Additive manufacturing · 3D modelling · Shape optimization · Lattice structures · Finite element analysis · Mass reduction · Bio-mimicry · Topology optimization · Variable neighborhood search

## 1 Introduction and motivations

Additive manufacturing (AM) is a relatively new but yet wildly developing industry which changes drastically the way of designing parts. Classical manufacturing processes like milling or machining tend to use physical tools to remove material, extracting the desired structure from it.

Differently, AM processes, especially Laser Beam Melting (LBM or DMLS) and Electron Beam Melting (EBM), tend to start from barely nothing, e.g., metal powders, and build the desired structure out of it, using energy beams (non-material tools) to locally fuse and solidify the powder [1]. Consequently, those new technologies introduce a totally different way of designing, but also new production constraints and limitations as well as new terminologies. Thus, the notion of lattice structures has become popular, and refer to meso-structures made out of inter-connected material beams. Lattice structures are an interesting way for AM to stand out from other production means which cannot generate such complex structures. Even if they still have mechanical properties yet to be investigated, lattice structures offer a very promising way for mass reduction. Any industrial sector (e.g., automotive, aerospace) that produces system which move tends to reduce mass for energy consumption purposes [2]. In this context, there is a consensus on the fact that AM is a good candidate for trying new designs and new

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numeric methods to produce lightweight products. To reduce the mass of an object, two main ways have been investigated: lattice structure and topology optimization. Both techniques lead to complex shapes that can be hard or impossible to manufacture without additive manufacturing. Today, there already exist examples of industrial success stories involving mass reduction through AM. For example, the Thales/ESA's deployment mechanism structure for which the mass has been reduced by 5, the costs reduced by 4, the number of components reduced by 10 thanks to a combination of topology optimization and lattice insertion (Fig. 1). Costs for such parts is high, it is still a niche market, essentially for aerospace, automotive and energy industrial fields. Moreover, as the final geometry might be convoluted, the computation time for optimization and the number of iterations can be very high.

In this paper, a new way to generate and optimize lattice structures is proposed. The contribution is threefold: (1) a non-uniform lattice structure generation technique is introduced, based on an analogy to crystallography; (2) a new optimization framework is developed to remove lattice beams according to mechanical criteria, thus decreasing the mass of the structure; (3) a new stress-based criterion is used to compare different results from the proposed algorithm. The resulting structure satisfies to user-specified requirements such as mass reduction and rigidity.

Even though researchers and industrials have investigated this research field, there still exist some limitations on the level of freedom the designers have when treating complex parts. The limitations of the existing approaches are discussed in Sect. 2. The proposed framework is introduced in Sect. 3 and the different steps are detailed in Sect. 4. The implementation and the results are presented in Sect. 5. Section 6 concludes this paper and highlights future perspectives of this work.

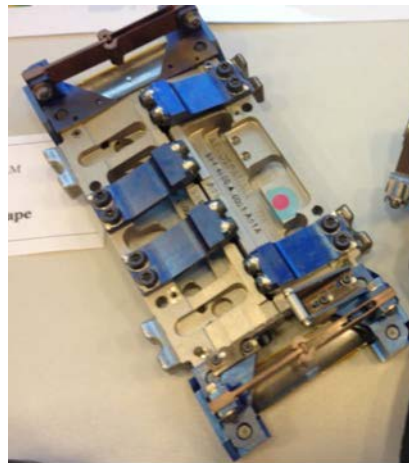
## 2 Related work

In recent years, different approaches have been introduced to define rigid, light and printable structures. They can be classified according to three main categories presented in the following sections.

### 2.1 Topology optimization

Continuous topology is a way of modifying structures to get lightweight parts while optimizing material layout within a given design space with the goal of maximizing the performance of the system [3]. Unfortunately, topology optimization can be time consuming when the required results needs to be of a high resolution (i.e., high number of voxels). Some other limitations can be discussed depending on the method. For instance, SIMP methods [4] may lead to noisy structures (Fig. 14a) and is the most used method in the industry due to its good computation time performances. Level set and/or BESO/ESO methods [5, 6] can significantly reduce noisy results but need a re-meshing step that can be time consuming in the long run. Topology optimization is also quite a sensitive analysis [7]. However, those limitations can be dealt with through seeing that topology results could be interpreted as inter-connected structures, plates or beams. This observation combined with AM capacities has also led to lattice structure studies [8]. In this paper, advantages of both techniques are combined in a unified framework which optimizes both the topology of the object and a non-uniform patterned lattice structure inserted in it. Such structures (inter-connected beams) can be found in nature in both organic and inorganic structures. Since AM allows to generate such structures with much more resolution than any other manufacturing process, it opens the field of biomimicry, i.e bio-inspired designs.

**Fig. 1** Satellite solar panel deployment mechanism (courtesy of Thales Alenia Space, ESA and Poly-Shape): **a** non-additive manufacturing design, **b** additive manufacturing redesign



(a)



(b)

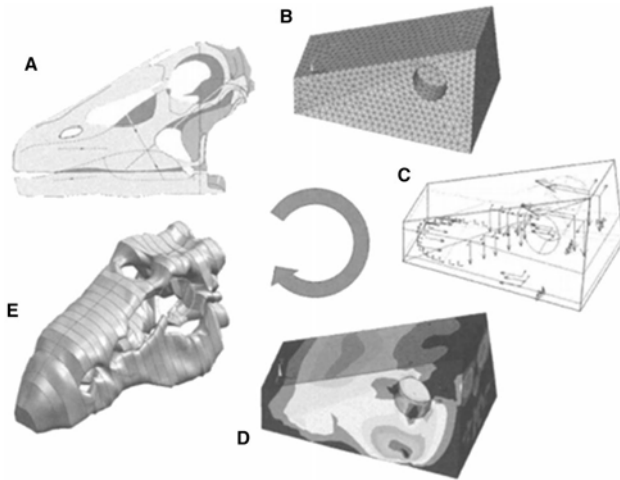
## 2.2 Bio-mimicry

Nature has always been a source of esthetic experience for the mankind and of inspiration for artists and designers [9]. Actually, natural shapes created by natural forces resemble topology optimization results. In 2005, Witzel [10, 11] managed to artificially compute a reptile skull through topology optimization. Through this study, they have shown that seeking for shape ideas in nature can be valuable (Fig. 2).

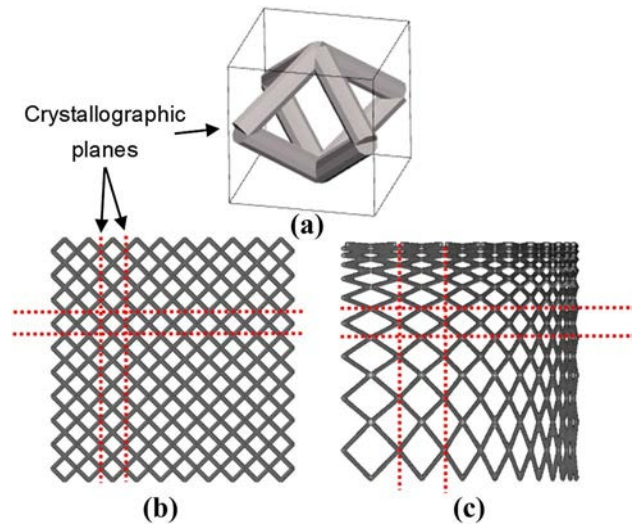
This seek for nature comparison and inspiration is called bio-mimicry and it is getting more and more investigated [12]. Some approaches tend to look for patterns that could be used to define lattice structures. Among them, Voronoi diagrams, Delaunay triangulation or fractals are examples of patterns that can be found in tortoise shells, soap bubble agglomerations, giraffe's skin patterns, and many others. Those schemes can be of interest in the AM context because they can be easily generated and propagated through a volume. However, they

represent a difficulty in terms of mechanical models as they require too many elements to be described properly both in terms of geometry and behavior. Non-living structures using lattice can be found in nature, the most famous being crystals [13]. Those lattices have been historically the first to be used in 3D printing for the simple reason that they consist of the duplication of a unitary cell (Fig. 3a)

Patterned lattice often refers to crystallographic structures, replacing atoms with nodes and inter-atomic links with beams (Fig. 3b). Such structure is commonly used due to its symmetry and its defined topology, which allows to quickly compute a mechanical response through some finite element models. Crystallographic lattice are defined in terms of inter-atomic planes. Following this analogy, it is then possible to stretch or compress the distances between those planes to increase the set of accessible shapes (Fig. 4). This can be implemented in different ways

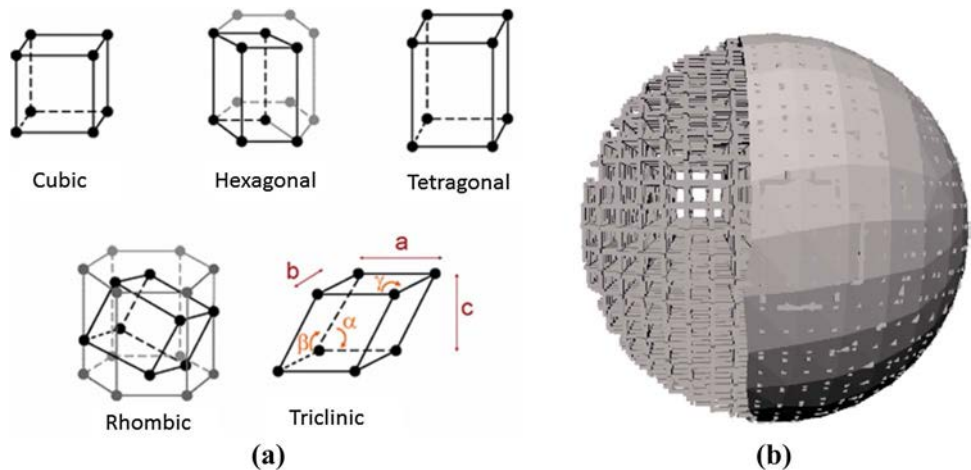


**Fig. 2** Topology optimization for bio-mimetic simulation of a skeleton [10, 11]



**Fig. 4** Example of lattice structures: **a** elementary cell, **b** uniform lattice deployment, **c** adaptive lattice deployment

**Fig. 3 a** Examples of atomic lattice structures [13]; **b** insertion of a crystallographic quadratic type lattice within a sphere





[14]. In this paper, as the structure needs to be optimized with respect to mechanical stress, the distances between the crystallographic planes will be defined so as to follow as much as possible the von Mises stress distribution. Thus, high stress density will contract the nearby lattice elements and low stress density will stretch it.

Nevertheless, once an elementary cell has been selected (Fig. 4a), and even though the distances between the crystallographic planes can be adjusted, the topology of the structure is fixed. This reduces the degree of freedom in the definition of the lattice. In some case, it would be interesting to let more freedom also in the definition of the topology of the lattice. The framework proposed in this paper allows such local modifications to optimize the topology as well as the geometry of the lattice structures.

### 2.3 Mechanical model

Another approach to define and optimize lattice structures is to make use of so-called homogenized materials [15]. This consists in looking for a fully filled material with similar properties to the lattice structure that is studied, and use this new material for finite element analysis as a solid material. However, this type of approach owns several shortcomings. For instance, as this equivalent material is set for a given lattice structure, the lattice topology (i.e., its connectivity matrix) cannot be changed throughout the optimization process (Fig. 5). This would mean that the equivalent material needs to be known for the new topology, thus leading to the need of an exhaustive equivalent material library. Anyhow, such an equivalent material would smooth the mechanical properties and the influence of the geometric structure.

Furthermore, this method would also not be directly reliable for un-patterned lattice structures, as the equivalent material would become a set of equivalent materials connected to each other. Such an approach has not yet been tested in the literature, but would likely lead to some difficulties.



**Fig. 5** Example of an optimized lattice structure using homogenized materials and preserving the topology (source Autodesk)

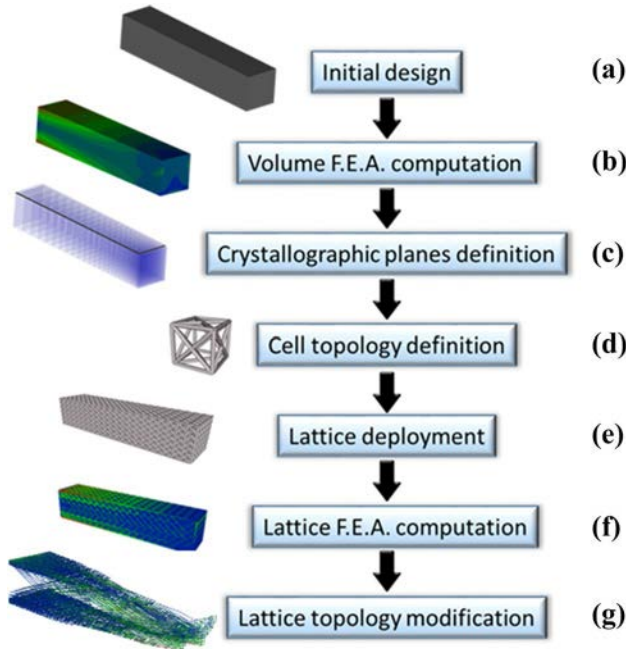
### 2.4 Non-uniform patterned lattice

As a conclusion, patterned lattice structures are widespread in current AM processes and are generally preferred to other lattice structures. Unfortunately, there are still gaps between the needs for generating non-uniform patterned lattice structures (Fig. 4c), the actual generation capabilities of 3D modelling software (Fig. 4b) and the capacities of modelling the mechanical response of those structures. This paper introduces a new framework to generate an initial non-uniform patterned lattice structure and optimize its topology while removing beams according to mechanical criteria. The generation and optimization modules are decoupled in the sense that the optimization could also be performed on a user-specified lattice structure that would not have been generated with our approach. The generation module defines the topology and geometry of the lattice, whereas the optimization module acts on its topology to try to reduce the mass while smoothing the stress distribution. During the optimization loops, the mechanical simulations are performed using a technic that derives from the methods of Alzahrani and Wang [8, 16]. A specific criterion is used to compare the optimality of the resulting structures in terms of stress distribution. The new framework is introduced in Sect. 3 and the details of the different modules are given in Sect. 4. Section 5 discusses the implementation and results.

## 3 Bio-inspired optimization framework

Going further into our bio-mimetic comparison, nature's ultimate issue is minimizing energy consumption. Any organ or body part needs to be fed with nutrients, oxygen and other elements. So, considering that the amount of energy generated per amount of time is fixed for a creature, it needs to distribute this energy among all of its body parts. The percentage of energy distributed to a particular body part depends on the contribution of the body part to the whole body capacity to sustain itself. In this paper, the beams of the lattice can be compared to bones of a skeleton, as their main purpose is to give rigidity to the creature's body (all other purposes of bones, such as fluid circulation, will be neglected to make it simpler). Inspired by this analysis, an optimization framework is defined, that looks for every beam (comparatively to bones in a skeleton), and sorts the contribution of each beam to the structure's stress contribution to decide if "it should be fed", meaning that it should remain, or "starve", meaning that the beam should be suppressed. This is close to what topology optimization does but applied to lattice structures. This principle has been experimented on human bones by manipulating the amount of stress received by a bone and looking at the evolution of the bone's shape and mass over time [17]. This is based both on Darwin's

theories of evolution [18] and bone structure observations [19]. Those studies infer that in terms of density (that can be related to beams diameter), bones tend to reduce their thickness and re-organize their inner structure as their stress load varies. Of course, such an evolution is done over a relatively long period of time. In this paper, this bio-inspired framework does not modify the beam's radii, but it will be considered that if the beam contribution is under a certain amount, this beam can directly be removed. In the future, further studies could act on radii evolution, thus, leading to a more accurate bio-inspired method. The proposed optimization framework consists of different steps presented in Fig. 6 and detailed in the next section. As previously introduced, to explore a wider range of available lattice structures, the inter-reticular planes can be stretched. To do so, a spatial scalar field is needed. von Mises field will be used here to position the crystallographic planes according to the initial stress distribution (Fig. 6b, c). Then, a lattice pattern is selected based on its mechanical behavior and its printability (Fig. 6d), it is then propagated within the volume (Fig. 6e). A first calculation is performed to ensure that the limit criterion (von Mises maximum stress) is not already reached (Fig. 6f). Finally, the optimization loop is run until the limit criterion is reached (Fig. 6g). As explained, this loop consists of suppressing beams that contribute the less to the mechanical "stress absorption" (this will be discussed later on) of the structure, according to our bio-mimicry statements as well as to a statistical evaluation of the considerate beam contribution among the whole beam population. This



**Fig. 6** Framework for patterned lattice structure optimization

process stops when a certain limit is reached (here, a von Mises maximum value) and the structure is saved.

## 4 Algorithms and methods

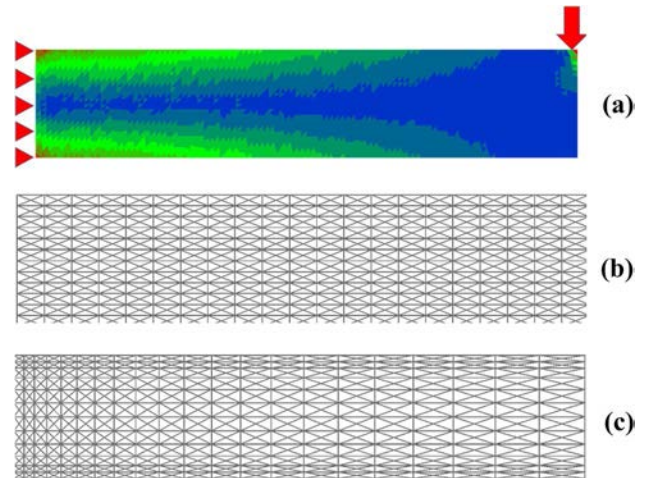
### 4.1 Crystallographic plane definition

The planes that define the cell's boundaries can be manipulated in such ways that their relative positions can vary non-uniformly in three dimensions. Thus, new designs can be tackled with a wider range of lattice fibers and possible curvature evolutions (Figs. 4, 7). To drive the position of the crystallographic planes, the von Mises scalar field is used as an attractive or repulsive field. Therefore, the more the von Mises stress is important, the more the distance between the planes is small. Of course, many other distribution laws could be defined and used. Figure 7b, c shows two distributions which can be obtained through that proposed method.

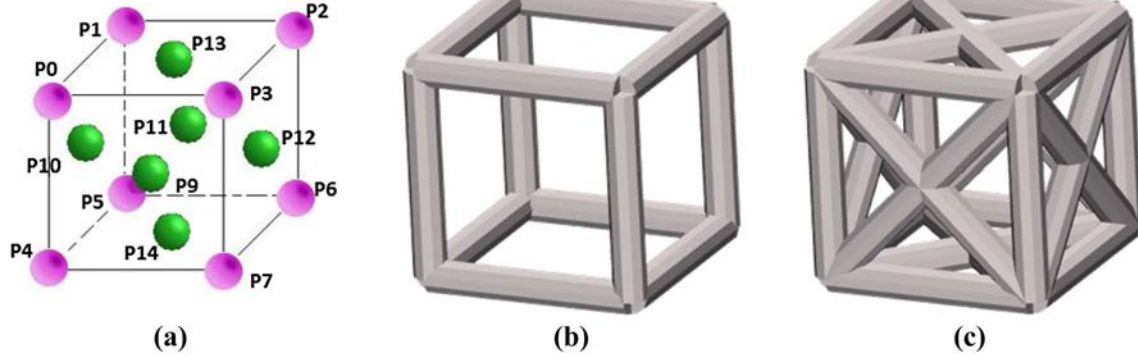
### 4.2 Cell topology definition

Lattice can be created out of crystallographic structures by replacing atoms with nodes and atom links with beams. Thus, several patterns can be imagined and compared in terms of mechanical behavior. To quickly illustrate the way an initial pattern can be identified, results for two extreme cases are presented: the so-called grid pattern (Fig. 8b) and face-centered cubic (FCC) pattern (Fig. 8c).

The topology of the grid pattern (resp. FCC pattern) can be encoded using a connectivity matrix  $C_{\text{Grid}}$  (resp.  $C_{\text{FCC}}$ ) coming from the Graph theory [20]. Considering a set of



**Fig. 7** Insertion of lattice structure: **a** von Mises field. **b** Uniform lattice insertion. **c** Non-uniform lattice insertion



**Fig. 8** Cell topology definition: **a** FCC crystallographic atoms and links. The grid pattern only uses corner nodes P0–P7 **(b)**, and the FCC pattern uses both corner and centered nodes P0–P14 **(c)**

$(n+1)$  nodes  $P_0-P_n$ , one can create a set of  $C_n^2$  potential beams, denoted  $B_{ij}$ , ranging from  $B_{00}$  to  $B_{nn}$  and containing all the combinations of node connections. The effective lattice structure  $\mathcal{L}$  does not necessarily contain all the elements of  $B_{ij}$  (it generally does not). An element of  $B_{ij}$  contained in  $\mathcal{L}$  is said to be “activated”. Thus, the elements  $C_{ij}$  of the connectivity matrix  $\mathbf{C}$  of dimension  $((n+1) \times C_n^2)$  are defined as stated in the following equation:

$$C_{ij} = \begin{cases} 1 & \text{if } B_{ij} \in \mathcal{L} \\ 0 & \text{if } B_{ij} \notin \mathcal{L} \end{cases} \quad (1)$$

Even if the Graph theory also allows for  $-1$  values, it is not needed for this purpose. Using this equation, the following two elementary matrices  $\mathbf{C}_{\text{Grid}}$  and  $\mathbf{C}_{\text{FCC}}$  can be obtained and used to define the whole lattice structure by looping through the volume. For the sake of readability, rows containing only 0 are removed.

$$\mathbf{C}_{\text{Grid}} = \begin{bmatrix} 1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix} \quad (2)$$

Using our approach, the topology of the initial lattice structure is defined by the connectivity matrix whereas the geometry of the lattice structure is given by the position of the crystallographic planes. Of course, the choice of a pattern impacts the mechanical behavior of the entire structure. To go further in this understanding, those two patterns have

been applied on a unitary cell of size 5 mm and have been loaded with a shear stress of 10 N. The radius of the beams is 1 mm. Results are displayed in Table 1.

One can notice that the FCC pattern offers a good compromise between a great reduction of the max displacement (about ten times less) and a relatively low increase of the mass (about two times more). So, even if our objective is to minimize as much as possible the mass, starting with a FCC-based lattice structure is a good idea. Such a choice can allow a great variety of orientations and a high possibility of combinations while securing the rigidity aspects. Anyhow, since our optimization process removes beams, it can converge to a solution close to the one that would have been obtained while using the Grid pattern. Finally, one can mention other researches that follow mechanical criteria to insert multiple cell types at different places within the structure [21]. In this case, the difficulty lies on the inter-connectivity between cells. Even if this can be an interesting way of improving computation times, it will not be studied in this paper.

### 4.3 Lattice deployment

Given the bounding box of the volume to fill in with cells, the lattice is fully known through both the number of needed cells per dimension and the deformation fields. Thus, two cases can be distinguished:

**Table 1** Estimated result values for both structures (for a same given  $n \times n \times n$  lattice cell unitary size)

	Grid	FCC	FCC/grid ratio
Max display (mm)	4.98E-03	5.26E-04	0.1056
Max VM (MPa)	27.8	5.01	0.1802
Volume (Cmm)	1.04E-07	2.28E-07	2.192



- *Uniform lattice* In this case, a simple loop based on the  $(dx, dy, dz)$  unit cell's dimensions, and the  $(X, Y, Z)$  bounding box's dimensions is used. The user gives the number  $N_i$  of cells per axis  $i \in \{x, y, z\}$  through which the unit cell's dimensions can be computed:

$$\begin{cases} dx = \frac{X}{N_x} \\ dy = \frac{Y}{N_y} \\ dz = \frac{Z}{N_z} \end{cases} \quad (3)$$

- *Non-uniform lattice* Here, a deformation field is needed to drive the crystallographic plane stretching. In the proposed approach, Bezier curves are used to define this deformation field along the three directions of the bounding box. For a given direction of the bounding box  $i \in \{x, y, z\}$ , the Bézier curve is defined as follows [22]:

$$\mathbf{P}_i(t) = \sum_{j=0}^{n_i} B_{j,n_i}(t) \times \mathbf{s}_j, \quad \forall t \in [0;1] \quad (4)$$

With  $(n_i + 1)$  control points positioned with respect to the local maximums of the von Mises stress that have been computed before. Those control points act like attractors on the curve, contracting planes around positions of high von Mises stress. The number of control points can be adjusted depending on the stress distribution. Adding control points would result in a more contracted lattice structure, even though their number generally remains below four; otherwise, other control curves could be useful such as NURBS. The initial and final control points are positioned on the boundary of the bounding box to cover the entire space. Finally, the user specifies the number  $N_i$  of cells per direction which is used to position the crystallographic planes according to points  $\mathbf{P}_i(t_k)$  of the Bezier curve located using a uniform law:

$$t_k = k \times \Delta_i, \quad \forall k \in \{0, \dots, N_i\} \quad \text{and} \quad \Delta_i = \frac{1}{N_i} \quad (5)$$

Since the volume to fill in is defined by a triangle mesh whose shape may be complex, a trimming step is required to trim the cells which go outside. Thus, once the planes have been positioned using one of the two techniques, it is possible to sort out the intersections between the planes and the triangles. Cells are then created at the intersection between six planes, they are filled in with a given pattern (Sect. 4.2) and potentially trimmed if they are intersecting with triangles. Follow-

ing this strategy, it is ensured that every beam is inside the triangle mesh.

## 4.4 Optimization loop

In this paper, the mass of the structure is reduced step by step while removing beams of an initially patterned lattice as defined in Sect. 4.2. The removal is performed according to a mechanical criterion, i.e., the distribution of von Mises stress. Manufacturing constraints are not fully addressed here, as the FCC pattern is printable regarding its stretched dimensions

in  $x, y$  and  $z$ . Further work would be needed to consider this printability constraints. The method is discussed in the next sections. The overall idea is to sequentially run a FE analysis (Sect. 4.4.1) to define the beams with a low stress, and then remove those beams (Sect. 4.4.2). The removal step is performed while following an optimization criterion (Sect. 4.4.3). Finally, the simplified structure is cleaned to remove self-hanging beams (Sect. 4.4.4).

### 4.4.1 FE analysis

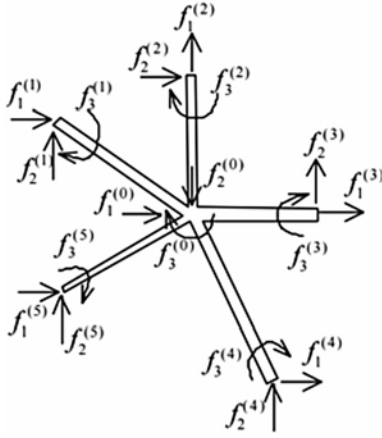
Once the lattice structure is generated and a load case is defined, the simulation model is sent to a finite element solver that returns the average von Mises stress in each beam. This stress distribution provides the contribution of each beam to the stress absorption. Here, the beams are modeled as a wire-frame structure (1D elements) to reduce computation times. This simplification brings up some limitations that are not discussed in this paper.

### 4.4.2 Beam deletion

To optimize the topology of the lattice structure, it is needed to understand how the structure behaves when beams are removed. For that matter, one can look at the structure in terms of inter-connected unit cells. This concept can adapt to any type of pattern, so this discussion will address a random beam neighborhood to remain general. It will only depend on the previously introduced connectivity matrix  $C$ . A lattice node is defined as a spatial point where beams interconnect with each other. It can be represented in terms of its neighborhood such as in Fig. 9 where the different loads are represented per beam [8].

Using the energy conservation principle, assuming the hypotheses of linear static mechanics and that beams have the same radii, one can go back to a simplified approximation of the virtual work principle. In 2D, one can write that the isolated system is such as (Fig. 9):

$$\begin{cases} \sum_{i=0}^5 f_i = 0 \\ \sum_{j=0}^2 f_j = 0 \\ \sum_{i=0}^1 f_i = 0 \end{cases} \quad (6)$$



**Fig. 9** Example of a node neighborhood and associated loads [8]

For each beam removal, Eq. 6 leads to a new stress repartition within the cells of the whole structure. The way the stress distribution evolves strongly relies on the lattice pattern and the load case. In the proposed approach, only low-stressed beams are removed so that the amount of stress that was circulating within the suppressed beam is redirected into the other surrounding beams. Consequently, the stress levels of the surrounding beams rise.

Let  $\delta(i)$  be the variation of stress on node connected to a removed beam. Then, each node in the remaining structure will see its stress level varying and become  $f_j^{\Phi(i)}$  such that Eq. 6 still remains and becomes :

$$\forall k \in \{0, \dots, 5\} \quad \begin{cases} \sum_{i=0}^n f_i^{\Phi(i)} = 0 \\ \sum_{i=0}^n \ddot{f}_k^{\Phi(i)} = 0 \end{cases} \quad (7)$$

$$\text{VMIS} = \frac{\overline{\text{EFFX}}^2}{\text{SECT}} + \frac{\text{MOMX} * \text{DX}}{\text{TORS}} + \frac{\text{MOMY} * \text{DY}}{\text{INRY}} + \frac{\text{MOMZ} * \text{DZ}}{\text{INRZ}} \quad (10)$$

With  $n$  the number of beams on a neighborhood. This stands for a 3D structure (i.e., 3D neighborhood). Then, for each node in the remaining structure:

$$f_j^{\Phi(i)} = f_j^i + c(i) \times \delta(i), \quad (8)$$

where  $\delta(i)$  can be discussed regarding the selected pattern (the different incoming angles, radius of each beams, and the applied load case). Thus, its value will not be defined here.  $c(i)$  represents the percentage of stress that was circulating through the node connected to the deleted beam that is now

circulating through another node. Then, through conservation of energy, the total amount of redirected stress needs to be of 100% and the local redirected amount of stress cannot overcome the deleted amount of stress. This statement leads to the following Eq. 9:

$$\sum_{i=0}^n e(i) = 1 \quad \text{with } e(i) \in [-1; 1]. \quad (9)$$

Thus,  $c(i)$  represents the percentage of stress that each remaining beam will have to sustain. This phenomenon leads to stress homogenization within the structure. It appears clearly that removing one low-stressed beam will not highly disturb the structure. However, either removing one high-stressed beam or too many low-stressed beams could lead to a high amount of stress redistribution, which could lead to the structure disruption.

It is then chosen to remove beams according to a criterion that will select a certain amount of elements based on their stress level and remove them. This criterion will not be discussed here, and so is the overall algorithm.

#### 4.4.3 Optimization criterion

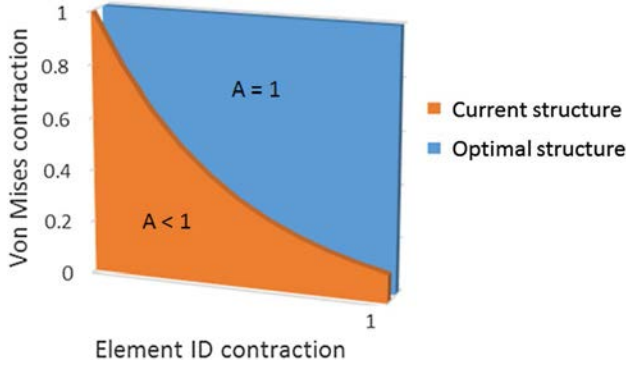
After the iterative removal process, a lattice structure is considered as optimal if all the remaining beams are at the same level of the von Mises stress (this level being defined by the user as the von Mises limit criterion). On that particular matter, as 1D beams are used, one could argue that von Mises criterion does not apply. The finite element solver used: Cast3m [23], proposes the same operator called "VMIS" to

compute an equivalent constraint criterion, no matter if elements are 3D, 2D or 1D. In the 1D case (beams) the criterion is computed such as:

with respect to the Cast3m notations which can be found in [23, 24]. In the proposed approach, this stress level directly corresponds to a user-specified limit. From that point, one can define a scalar value  $A$  that will be used as a performance

indicator for the algorithm. This scalar value is the area under the curve representing the distribution of the stress level in the lattice structure (Fig. 10). For an optimal structure, i.e., a structure for which the remaining beams are at the same level of the von Mises stress, this area is simply equal to:

$$A_{\text{optimal}} = N_B \cdot L_{\text{VM}} \quad (11)$$



**Fig. 10** Rescaled distribution of the stress level for the current and optimal structures

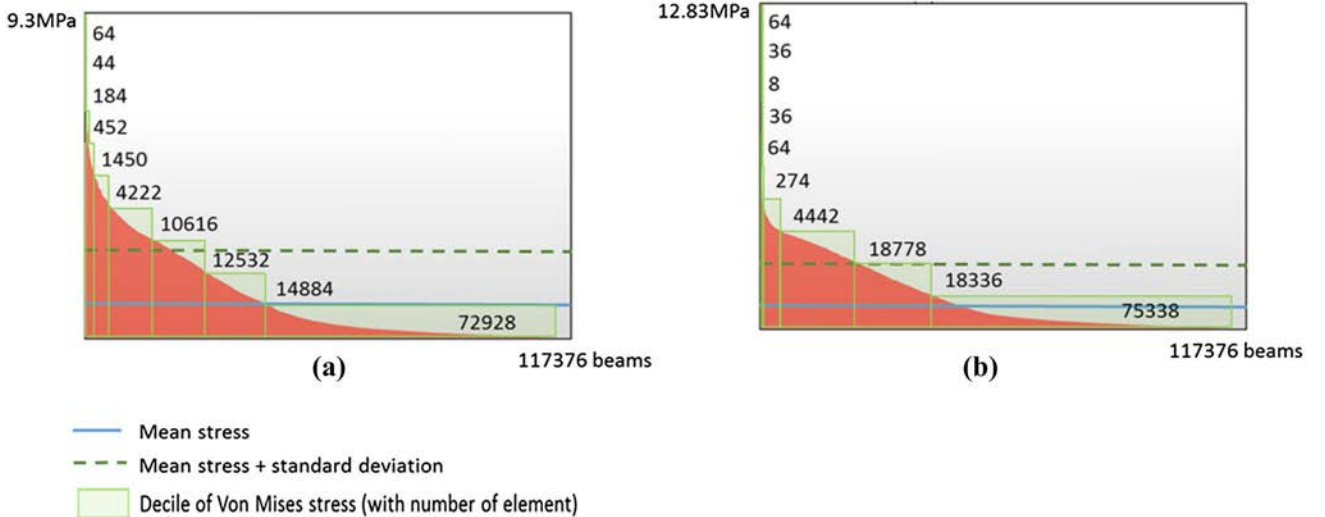
where  $N_B$  is the number of beams and  $L_{VM}$  is the limit von Mises criterion. To compare lattice structures composed of a different number of beams, it is preferable to rescale the von Mises stress and the number of beams over the interval  $[0; 1]$  as shown in Fig. 10.

This criterion helps evaluating how much the stress is equally distributed in the lattice structure during the optimization loop. Consequently, the closest  $A$  is from 1, the closest the structure is to the optimal one (Fig. 10). Thus, at each step of the optimization process, knowing the stress repartition in the structure, one can remove the lattice elements that are unnecessary for the structure rigidity. However, removing the elements one by one would lead to high computation times making it non-suitable for industrial applications. It is thus needed to use a criterion to remove sets of beams and not the beams one by one. Ideally, one

could use a criterion that mimics those used in topology optimization, i.e., based on either a level set method or a homogenization of the compliance matrix [25]. However, an adaptation has to be performed to use this criterion on a lattice structure. A result of such a method on 3D elements is discussed in Sect. 5. In this paper, a statistical data criterion are used to sort low-stressed beams. To be effective, the adopted criterion must be acceptable in terms of computation times and lead to a non-degenerated structure, i.e., a structure which can still sustain the load case. Here, the so-called cut criterion  $C_{cut}$  corresponds to the difference between the von Mises mean stress  $M_{VM}$  and the von Mises stress standard deviation  $SD_{VM}$ :

$$C_{cut} = M_{VM} - SD_{VM} \quad (12)$$

Using a statistical approach leads to a non-absolute criterion, each beam gets a notation relatively to other beams and to the maximum stress required. It means that even if the absolute simulated stress values deviate from the experiment, it is still accurate in terms of relative stress distribution, and one can still sort out which beams need to be removed. Moreover, since wireframe simulation models are far faster to compute than 3D models, this method can compute dense lattice structures quickly and as the criterion is relative, one only needs to calibrate a safety coefficient based on expert knowledge to stop the optimization process. Then, if mandatory, a 3D finite element analysis can be used to get more accurate results. The proposed criterion gives both the position of the beam within the whole population in term of von Mises stress (through standard deviation) and its distance to mean stress. For an homogeneous population of beams, this leads to a removal of less than 10% of



**Fig. 11** Initial  $A$  criterion for both classic and stretched lattice structures (Fig. 7), von Mises stress value per element sorted from highest to lowest. **a** Uniform lattice, **b** non-uniform lattice

the total population at each iteration [26]. Furthermore, the number of beams that will be removed is proportional to the total number of beams remaining within the structure. This allows to auto-regulate the beam population. This process directly uses the von Mises (approximated von Mises due to wireframe models) distribution. It affects the connectivity matrix  $C$  and the final lattice structure is not anymore a pattern. Figure 11 illustrates two stress distributions for the two crystallographic plane arrangements of Fig. 7. A pure shear load case of 10 N is applied at the extremity and no beams are removed. Even for the same FCC elementary cells, and for the same number of beam elements, the stress distribution can vary depending on how the crystallographic are distributed. Figure 11b shows that the non-uniform adaptive lattice has a more homogeneous stress repartition (see decile of von Mises stress). Applying our performance criterion would lead to a value  $A = 0.027$  for the uniform lattice and  $A = 0.030$  for the non-uniform lattice (11% gain) (Fig. 7).

#### 4.4.4 Cleaning

Once the beams have been iteratively removed and the stop criterion has been reached, self-hanging beams are removed. They correspond to beams that sustain no stress. Thus, beams that have a stress level of 0 MPa ( $\pm c$ ) are removed. As a remark, one can understand that as 1D beams are used, the stress level of a beam  $B_{ij}$  is computed as:

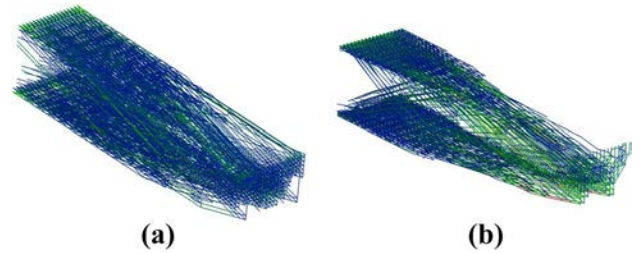
$$VMIS(B_{ij}) = \frac{1}{2} \cdot (VMIS(N_i) + VMIS(N_j)) \quad (13)$$

Then, for an overhanging beam, the stress level on one node would be 0, and on the other node would it have a certain value, leading to a beam stress level superior to 0. The

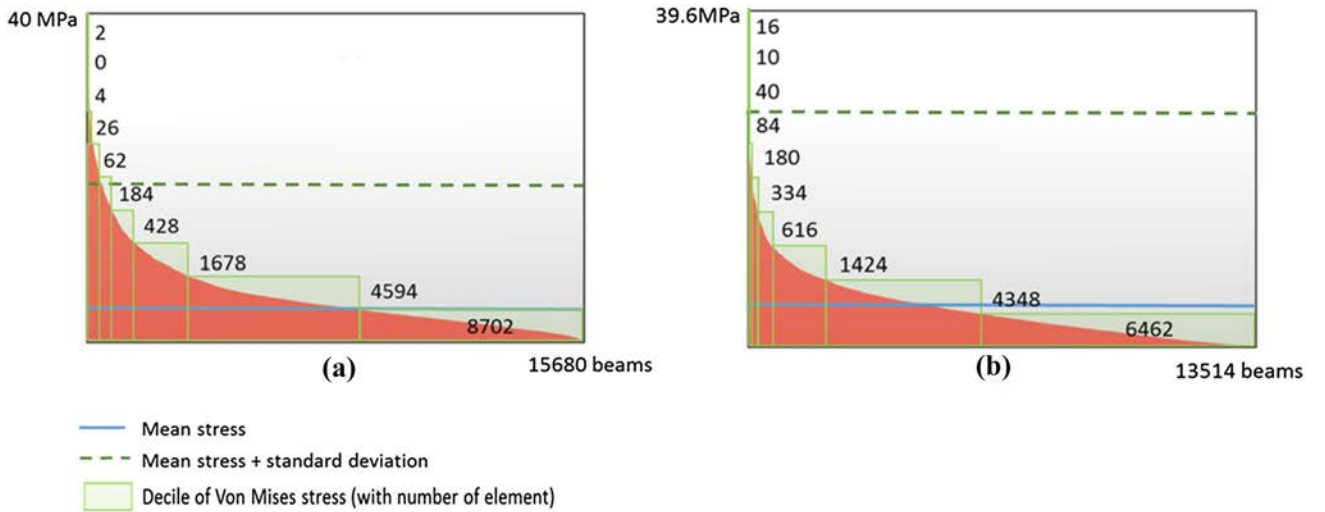
removing of unnecessary beams can then not be based on their stress level but on the fact that one of their nodes is not stressed. This issue could be address more effectively using the  $C$  matrix to sort out overhanging beams (nodes with only one beam connected) and delete them.

## 5 Implementation and results

The proposed framework has been implanted on our prototype platform developed in C# and based on MVVM WPF technology. It creates non-uniform patterned lattice structures that are then optimized using Cast3m as a finite element solver, integrated in our optimization loop. To validate the proposed approach, the case study of Fig. 7 will be used. A load case of 10 N perpendicular to the structure main axis is applied at the extremity, and the user-specified maximum von Mises criterion  $L_{VM}$  is set to 40 MPa. After few iterations, the structure has evolved towards an optimized structure that suits the technical requirements and allows an effective mass reduction. The stress has been homogenized within the structure that is not degenerated (Figs. 13, 14).



**Fig. 13** Isometric views of the final structures. **a** Uniform lattice. **b** non-uniform lattice

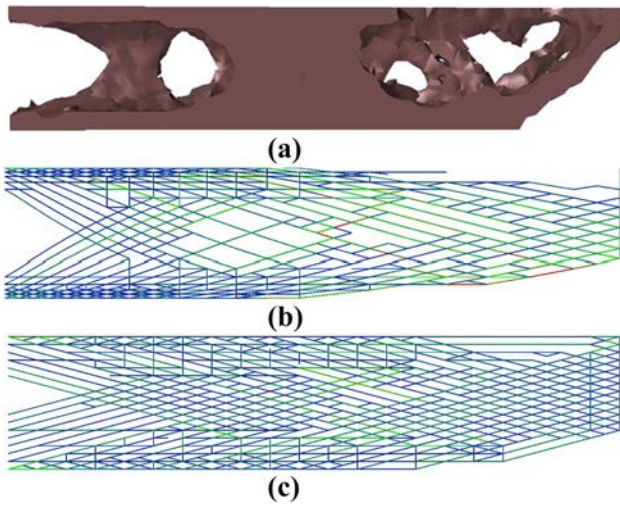


**Fig. 12** Results after the optimization, von Mises stress value per element sorted from high to low. **a** Uniform lattice, **b** non-uniform lattice



**Table 2** Comparison of the maximum stress and associated volume for the initial/final and uniform/non-uniform structures

	Uniform	Non-uniform	Ratio (%)
Nber of beam			
Initial	117,376	117,376	0.00
Final	15,680	13,514	16.03
Max VM (MPa)			
Initial	9.3	12.8	+27.3
Final	40	39.6	1.01
Volume (mm <sup>3</sup> )			
Initial	5092	4998	1.88
Final	960.6	721.1	33.2
mean VM (MPa)			
Initial	1.11	1.21	+8.26
Final	4.4	5.6	+21.43
std. deviation			
Initial	1.8	1.88	+4.26
Final	15.2	24.2	+37.23



**Fig. 14** Comparison between **a** SIMP continuous topology, **b** non-uniform lattice structure and **c** uniform lattice structure

Here, our performance indicator is  $A = 0.11$  for the uniform lattice and  $A = 0.14$  for the non-uniform one with a relative gain of 27% (Fig. 12). This illustrates a better von Mises stress distribution throughout the beam population.

As one can see comparing Figs. 11, 12, the optimization process tends to smooth the curve of stress repartition, illustrating the homogenization process. Figure 12 also contains the decile repartition of von Mises stress for the uniform and non-uniform lattice structures. Comparing Figs. 11, 12 shows that through the overall diminution of beam elements, the stress level in each of those elements has raised

up toward the user-specified von Mises limit criteria. It corresponds to the homogenization phenomenon.

Figure 13 shows that the denser parts of the resulting structures are located on the upper and lower planes near the fixations, and that the center area's density has been significantly lowered. Such a behavior is consistent with what usual optimization leads to, like IPN structures for example.

Table 2 gives output values, the ratio column provides the percentage of relative values from non-uniform to uniform structure. The gain from initial to final volume is of 85.6% for the non-uniform structure and 81.1% for the uniform, showing that with the same initial number of beam, their spatial repartition is of importance for the optimization process. Finally, Table 2 compares all the results and sorts the obtained gain for each value (gain from non-uniform to uniform lattices). Mass was the main objective of the optimization process and can be improved, for this test case, by 33% only by stretching the lattice structure. One can also see that mean stress is higher (and so closer to the limit criterion) when using a non-uniform lattice.

Figure 12 shows that the stress distribution is better than for the initial lattice structure (Fig. 11) as most of the beams have a stress level under [mean stress + standard deviation]. The performance indicator  $A$  is still smaller than 1, but its value has been multiplied by a factor 5 compared to the initial structure (from 0.03 to 0.14). In this work, only the removal of beams was authorized, but allowing other transformations could lead to an even better distribution and could help raising the  $A$  value further more. Among those foreseen transformations, one could cite beam radius evolution, beam switching and beam to plate transformation.

To further analyze the results, a classic topology optimization based on SIMP method has been run on the initial fulfilled volume. The results are shown in Fig. 14. One can see that the general shape of our optimized lattice structure tends to keep material in the same place as the classic topology optimization does.

Finally, one can see that the way the crystallographic planes are distributed is important for the optimization process. Thus, alternative plane distribution rules still have to be studied so as to improve even more the performance indicator and reduce the mass.

## 6 Conclusions and future works

In this paper, a non-uniform patterned lattice structure optimization framework has been proposed and validated. The optimization process starts with an initially non-uniform structure and iteratively removes beams according to the von Mises stress distribution within the whole structure. At



the end, the structure becomes un-patterned, the stress is better distributed and the mass is reduced. Compared to Fig. 2, results of Fig. 14a, b illustrate that the proposed approach also tends to generate nature-like structures.

A first remark one would have would be to use directly the constraints tensor field instead of the von Mises scalar field. This would allow to differentiate tension from compression, and to enable a better space redistribution of the crystallographic planes. For the sake of simplicity, and as a first attempt, von Mises criterion has been chosen precisely for being a simple scalar instead of nine scalars embedded in a tensor, but future works will extend this particular point. Such method has already been tested and has shown interesting results [27].

Similarly, using 1D finite elements does not allow to go up to a complete optimization, but to get a lattice shape containing as less beams as possible. Then, it is necessary to run a 3D finite element calculation to ensure that the technical requirements are met. This is an alternative to common approaches that still needs to be improved but allows to change lattice topology during optimization contrary to homogenized material methods (Fig. 5). This widens the design possibilities.

Finally, through Figs. 12, 14 one can see that there is still potential for pushing the von Mises front up to the limit criterion and get the  $A$  criterion as close as possible to 1. As already said, this improvement could be done by allowing more transformations. Instead of only removing beams, beams could be switched, new beams could be added, beam's radii could be changed and so on [26]. Moreover, Fig. 14 and Table 2 illustrate that working on non-uniformly distributed crystallographic planes improves significantly the performance indicator  $A$  and leads to a significant mass reduction when compared to a uniform lattice structure. Thus, the way those planes are located throughout space still needs to be thought. It can be seen as a pre-optimization process. This method could give access to a wider range of reachable geometries and topologies for complex structures, improving material space repartition in a continuous topology-related fashion, thus reducing significantly the mass of systems through lattice structures.

Finally, the comparison between the adopted simulation models and the corresponding experiments was not tackled and still remains an issue for this method to be efficient.

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