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# Optimal dual-PZT sizing and network design for baseline-free SHM of complex anisotropic composite structures

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## Abstract

Structural health monitoring (SHM) processes for aeronautic composite structures are generally based on the comparison of healthy and unknown databases. The need for prior baseline signals is one of the barriers to industrial deployment and can be avoided with *baseline-free* SHM (BF-SHM) methods, based on the attenuations and reflections of symmetric and antisymmetric Lamb wave modes attributable to damage. A promising mode decomposition method is based on the use of dual-PZTs (concentric disc and ring electrodes lying on a single piezoelectric transducer of lead zirconate titanate). However, the performance of such methods highly depends on the Lamb wave mode properties (propagation speed and attenuation, which vary with material orientation and inter-PZT distance), the number and the sensitivity of the dual-PZT to each mode (which depends on the frequency and element size). Considering these constraints, an original three-step process able to design a full dual-PZT network and the optimal range of excitation frequencies to consider on a highly anisotropic and arbitrarily complex aeronautic structure is presented. First, the dispersion curves of Lamb waves in the investigated material together with the minimal size of the damage to detect are used to estimate the size of the dual-PZT, as well as convenient excitation frequencies. A local finite element model representative of the full-scale structure is then used to estimate the optimal distance and orientation between neighboring PZT elements. Finally, a network optimization solver applies these parameters to place dual-PZTs on the fan cowl of an aircraft nacelle and provides a candidate network covering the whole structure.

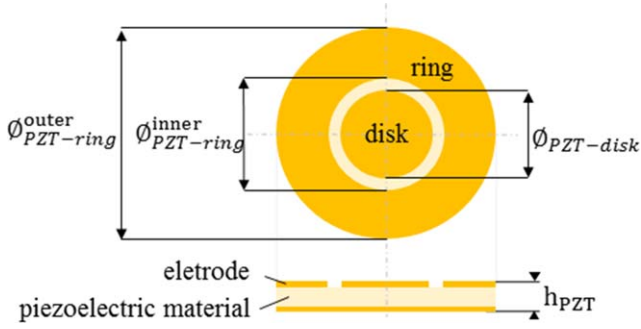
Keywords: structural health monitoring, Lamb waves, dual-PZT, anisotropic composite structure, mode decomposition, baseline-free methods

(Some figures may appear in colour only in the online journal)

## 1. Introduction

The real-time and autonomous monitoring of the health state of structures is referred to as structural health monitoring (SHM) [1]. Damage detection using Lamb waves is one of the most common SHM methods for aeronautic composite structures. Among all the methods employed to trigger and catch Lamb waves, the use of a piezoelectric transducer (PZT)

is one of the cheapest and easiest to settle. A common approach is based on the comparison of signals recorded on a pristine structure—the baseline—against those obtained on a possibly damaged one. The resulting signals, which will be used to elaborate the damage indexes, strongly depend on the structure's mechanical and geometrical properties as well as on environmental conditions. Some solutions have been proposed to compensate for the effect of temperature [2–5] or



**Figure 1.** A simplified view of a dual-PZT composed of a concentric ring and a disk.

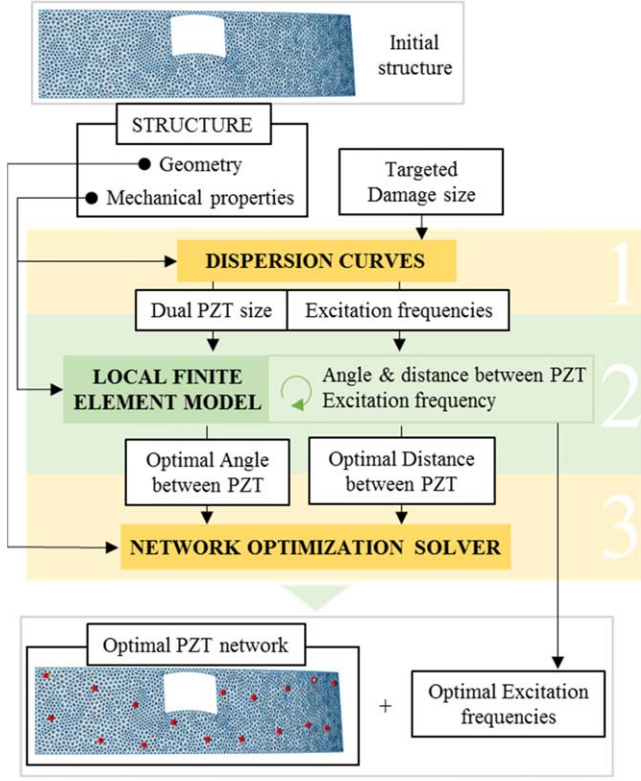
to extend baseline data to a wider range of environmental conditions [6, 7]. This drawback of existing methods has also fed the interest for baseline-free methods. Baseline-free is an abusive word here (see axiom II in [1]) since data, models or physical assumptions are always needed to decide whether a structure is damaged or not. This expression thus refers to methods where prior reference signals recorded from the pristine condition of the monitored structure are not necessary for damage detection.

Symmetric and antisymmetric Lamb wave modes are very useful for SHM as they convey complementary information when interacting with damage [8, 9, ch. 19.4]. Lamb wave modes are related to Lamb solutions for wave propagations in isotropic plates but are also commonly (and abusively) used to designate plate-like waves that propagate in anisotropic materials. Early studies on damage detection in composite structures state that the  $S_0$  mode is more appropriate than the  $A_0$  mode for delamination detection, since the former exhibits higher sensitivity to the structural thickness [8, para. 3.3]. In most SHM applications on composite structures, the  $S_0$  mode is actually preferred since it has lower attenuation, allowing for larger distances between neighboring PZTs in the designed network. However, the  $A_0$  mode could be more appropriate since it allows interaction with delamination at smaller excitation frequencies than the  $S_0$  mode and has been favored in most recent works [10–12]. Actually, both modes can be used as long as the excitation frequency chosen (and the resulting wavelength of each mode) allow interaction with the damage [9, para. 19.5.2.2]. Common methods used to study the behavior of a given mode ( $A_0$  or  $S_0$ ) are based on the response signal obtained with collocated transducers [9, ch. 17], or the use of sweet spot frequencies [13, ch. 8] where only one mode is generated. The former method requires access to both sides of the instrumented structure as well as the perfect collocation of the transducers, and the latter restrains the application to specific excitation frequencies that may not correspond to the amount of damage to detect. These limitations are overcome by using dual-PZTs. dual-PZTs are made of a concentric ring and disk, as shown in figure 1. Both parts can be used as actuators and sensors resulting in more signal combinations than conventional PZTs and allowing for the isolation of the first antisymmetric  $A_0$  or symmetric  $S_0$  mode. With dual-PZTs, damage can then be detected in a baseline-free manner by

using the damage-introduced mode conversion and attenuation (which is not the same for each mode) of the propagating waves [8]. The theoretical aspect justifying the use of these PZTs for baseline-free SHM (BF-SHM) has been well developed in [14], and experiments are reported in [15] on aluminum and in [16] on composite structures. However, the interaction of each mode with the damage is conditioned by the wavelength of the waves generated by the dual-PZT and the damage dimension. As a consequence, the transducer's dimensions as well as the excitation frequencies must be chosen with care. The first objective of the present work is to propose a method for dual-PZT sizing and excitation frequency selection according to the targeted application, i.e. BF-SHM. Furthermore, previous works with dual-PZTs have been carried out on quasi-isotropic structures but, to our knowledge, no studies have been reported on anisotropic structures. In such materials, the amplitude as well as the propagation speed of the  $A_0$  and  $S_0$  modes on a path between two PZTs depends on the angle between the path direction and the orientation of the carbon fibers composing the structure [9]. Since the mode decomposition method depends on the amplitude ratio between both modes, the PZT positions (in terms of distance and orientation) need to be chosen carefully to guarantee the sensitivity of both modes in any configuration. The second objective of the present work is then to propose a method to determine the optimal distance and orientation between neighboring dual-PZTs. Another underlying objective is to optimize the number of active elements, which is a key point for industrial deployment, especially in aeronautical applications.

Some of the BF-SHM methods, where mode decomposition using a dual-PZT is of interest, are quickly reviewed here. The instantaneous baseline method is based on a transducer network with the paths having the same dimensions and orientations. Instead of comparing the signals obtained in unknown conditions with a pristine baseline, they are compared instantaneously between paths even under varying environmental conditions [17, 18–21]. Drawbacks concern the examined paths, which must have identical PZT positioning and bonding and avoid elastic wave reflections from the boundaries. Another technique, called the *time reversal* process, assumes that for a pristine structure, by sending back the reverse signal received by a given PZT, the input signal used in the first place should be recovered. Damage potentially breaks this relation, allowing for its detection. This method has been validated on composite structures [11, 22–25] and its limitations have been identified [26], namely, the high dependency on the excitation frequency, mode dispersion and the reflections from the boundaries. In the *reciprocity principle* method, two signals obtained from both possible directions of a given path are directly compared. This method has been used successfully in aluminum [15, 17, 27] and quasi-isotropic composite [12] structures. Its drawbacks are that the implemented PZTs must be identically bonded to the structure, and that damage localized far from the investigated path is difficult to detect.

BF-SHM studies focus on a description of the damage detection methods, but do not address the problem of network design when dealing with large structures. However, the PZT



**Figure 2.** The different steps for the optimal design of a dual-PZT network on an anisotropic composite structure.

network (size, number, distance and orientation of transducers as well as excitation frequency) to be used highly depends on the considered damage detection method, the environmental conditions, and the geometry of the structure. This step must thus be considered with great care. There have been studies on optimal sensor placement (OSP) for reference-based Lamb-wave-based SHM applications [28–33] with recent focus on composite structures [33–35]. These studies show that the relevant criteria for PZT positioning are the distance between two neighboring PZTs and the presence of complex geometries [34] (large holes, stiffeners, etc) on the paths between PZT. In BF-SHM applications, additional and very restrictive constraints on the PZT networks exist. Based on the observations described above, the current work proposes an original three-step process to design an optimal dual-PZT network for a structure to be monitored by means of BF-SHM methods. As depicted in figure 2, the different steps of this process are:

1. The mechanical properties of the structure are used to compute the dispersion curves of Lamb waves in the investigated material. Dispersion curves and the minimal amount of damage to be detected allow the dual-PZTs to be sized and the excitation frequencies of interest to be determined.
2. A local FEM model (a plate corresponding to a small part of the monitored structure) equipped with two dual-PZTs is used. Lamb wave attenuation is extracted with a distance  $r$  and orientation  $\theta$  and for different excitation frequencies  $f_0$ . Two original indexes quantifying the ability of the set  $(f_0, r, \theta)$  for mode decomposition are

then computed. This provides the optimal parameters  $(\hat{f}_0, \hat{r}, \hat{\theta})$  for the neighboring dual-PZTs of the network.

3. The network optimization solver applies the optimal parameters obtained from the local FEM model to the real structure, to propose an optimal dual-PZT configuration candidate.

The rest of the paper is organized as follows: in section 2, the material and geometry of the targeted aeronautical structure under study are presented. Then in section 3, each step of the process is described and illustrated with results obtained on a highly anisotropic composite material. Finally, in sections 4 and 5, the PZT networks obtained for different geometries and different composite materials are displayed, analyzed and discussed.

## 2. Material and structure under study

Composite materials are of great interest in an aeronautical context since they provide lighter structures with the same mechanical properties as previously used materials. The composite material under study is a highly anisotropic composite-fiber-reinforced polymer (CFRP) with 16 plies oriented at  $(0^\circ/90^\circ)$ . The mechanical properties of one ply are listed in table 1.

A part of a nacelle from an A380 plane (the fan cowl as shown in figure 3) is considered to illustrate the three-step process proposed to design an optimal dual-PZT network for BF-SHM methods. This geometrically complex structure is 2.20 m in height for a semi circumference of 5.80 m and is made of a composite monolithic carbon epoxy material. Piezoelectric elements manufactured by NOLIAAC (see table 2) are used to trigger and collect Lamb waves in the structure.

## 3. Optimal dual-PZT sizing and placement

### 3.1. PZT dimension and excitation frequency

In the first step of the proposed method, the dimensions of the PZT to be used as actuators and sensors and the excitation frequencies are determined based on the material properties and on the minimal size of the damage to be detected. This approach is developed to meet three main goals (see figure 4):

- (i) Only the first symmetric and antisymmetric modes  $S_0$  and  $A_0$  must be generated.
- (ii) The dual-PZT used as an actuator and the frequencies of the signal used must guarantee the excitation of the  $S_0$  and/or  $A_0$  modes with a wavelength small enough to interact with the targeted minimal damage.
- (iii) The dual-PZT used as a sensor must be sensitive to both  $S_0$  and  $A_0$  modes over the selected range of excitation frequencies.

The dispersion curves are computed numerically by finding all the possible frequencies that satisfy the Lamb wave dispersion

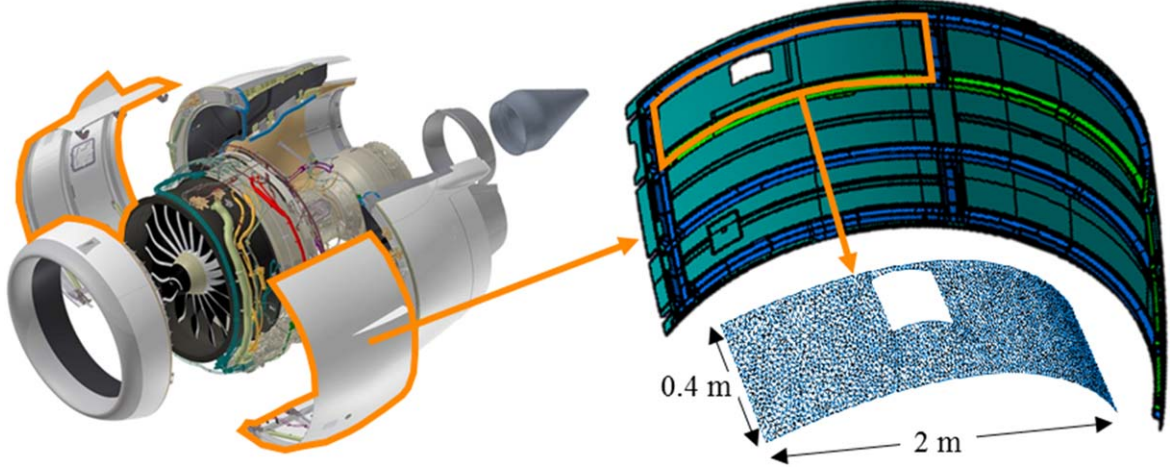


**Table 1.** The properties of a layer of the CFRP composite laminate.

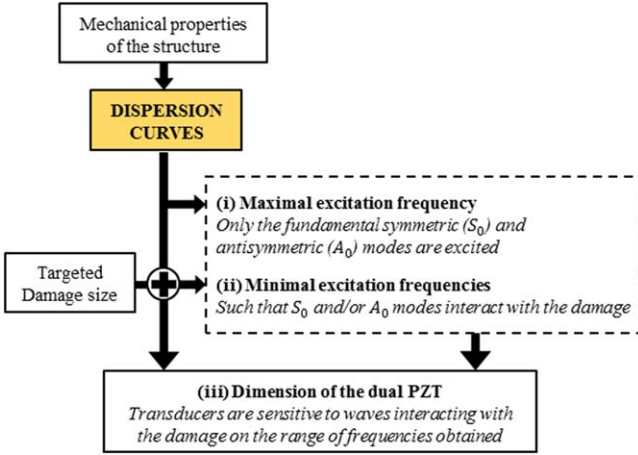
$E_1$ (GPa)	$E_2 = E_3$ (GPa)	$\nu_{31}$	$\nu_{23} = \nu_{12}$	$G_{23} = G_{31}$ (GPa)	$G_{12}$ (GPa)	$\rho$ (kg.m <sup>-3</sup> )
121	8.6	0.06	0.31	4.3	5.9	1520

**Table 2.** The electro-mechanical properties of the PZT material NCE51 from Noliac.

$E_{11}$ (GPa)	$E_{33}$ (GPa)	$a$ (1/K)	$\nu$	$d_{31}$ (pC N <sup>-1</sup> )	$d_{33}$ (pC N <sup>-1</sup> )	$s_{11}^E \cdot 10^{-12}$	$s_{33}^E \cdot 10^{-12}$	$\rho$ (kg.m <sup>-3</sup> )
62.50	52.63	$2.10 \cdot 10^{-6}$	0.38	-195.00	460.00	16.00	19.00	7600



**Figure 3.** The aircraft nacelle (left), the original fan cowl geometry and the simplified 3D geometry of the subpart being studied (right).



**Figure 4.** A schematic illustration of the three steps followed to select the dimensions of the dual-PZT from the dispersion curves. This corresponds to step 1 in figure 2.

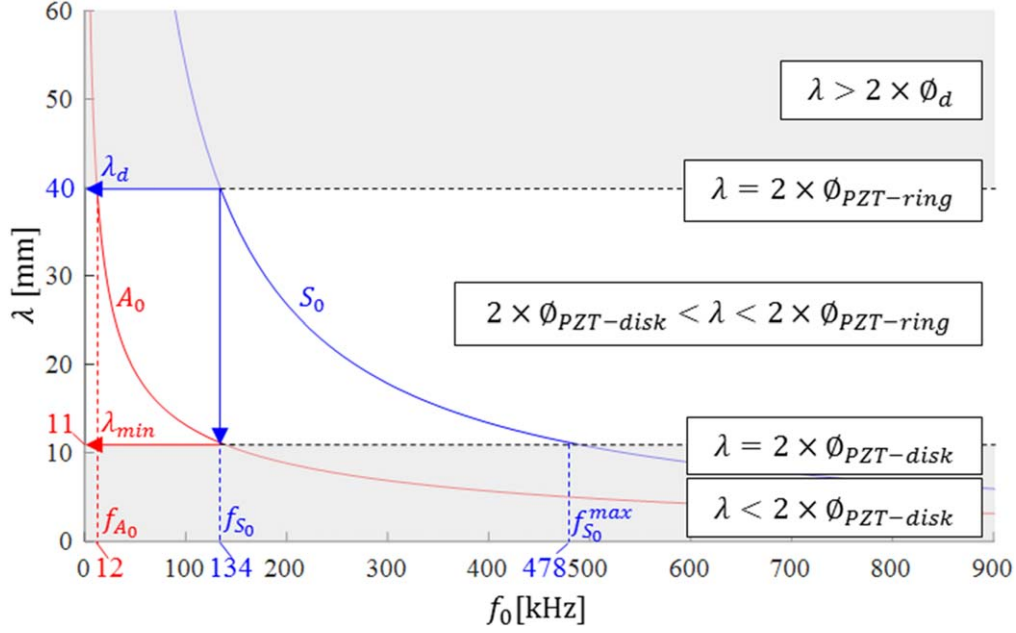
relations with iterations on the wave number  $k$  [36]. Once those values are computed, the dispersion curves can be displayed as a variation of the wavelength  $\lambda = 2\pi/k$  of each symmetric and antisymmetric mode with the frequency. The dispersion curves can also be computed for some rheological materials (please refer to [37] for some open software). The dispersion curves displayed in figure 5 are computed considering plane waves propagating in a homogenized material with the properties given in table 3 corresponding to the  $0^\circ$  direction of a  $(0^\circ, 90^\circ)_{16}$  CFRP plate.

The  $0^\circ$  orientation is chosen since it is the direction in which the wavelengths are larger; hence a dual-PZT with the right dimensions for these wavelengths will respect the criteria described in figure 4 in all orientations.

The BF-SHM methods targeted here consist of tracking the  $S_0$  and  $A_0$  mode conversion and attenuation. It is thus necessary to restrain the domain of applications where only those first two modes propagate. The dispersion curves for higher frequencies (not displayed here) show that the excitation frequency should not exceed  $f_{max} = 906$  kHz (corresponding in our case to the minimal excitation frequency for which mode  $A_1$  is generated) to meet this condition. A second point to consider is that theoretically half the wavelength  $\lambda$  of the wave used to detect damage must be shorter than or equal to the size of the damage  $\varnothing_d$  in order to allow interaction with it [38]:

$$\frac{\lambda}{2} \leq \varnothing_d. \quad (1)$$

This leads us to bound the frequency range under which  $S_0$  and  $A_0$  can be exploited: the largest wavelength validating equation (1) is  $\lambda_d = 2 \times \varnothing_d$ . According to this criterion, the excitation frequency used should be greater than the frequency  $f_{A0}$  for the  $A_0$  mode and  $f_{S0}$  for the  $S_0$  mode, corresponding to excitation frequencies for which the wavelengths  $\lambda_{A0}$  and  $\lambda_{S0}$  are lower than  $\lambda_d$  (see figure 5). As advised by the aeronautical manufacturer who provided the samples, the minimal amount of



**Figure 5.** The dispersion curves for the material described in table 3 showing the minimum frequency to use ( $f_{A0} = 12$  kHz for the  $A_0$  mode and  $f_{S0} = 134$  kHz for the  $S_0$  mode) to allow for interaction between the propagating wave and damage of  $\varnothing_d = 20$  mm. It also shows the range of frequencies to use to guarantee the sensitivity of the dual-PZT to the  $A_0$  mode and  $S_0$  mode for frequencies up to  $f_{S0} = 134$  kHz. The disk part of the dual-PZT is also sensitive to an  $S_0$  mode up to  $f_{S0}^{max} = 478$  kHz;  $\varnothing_{PZT-disk} = 5.5$  mm and  $\varnothing_{PZT-ring} = \varnothing_d = 20$  mm.

damage that needs to be detected (BVID for barely visible impact damage) is 20 mm, hence in our case,  $\lambda_d = 40$  mm,  $f_{A0} = 12$  kHz and  $f_{S0} = 134$  kHz (see figure 5).

Furthermore, it is mandatory to ensure that at these frequencies the dual-PZT element can actuate and receive the desired modes. This then leads us to select optimal dual-PZT dimensions. Two criteria are important in the choice of the PZT according to [39]:

- As an actuator, an optimal PZT size  $\varnothing_{PZT}$  is obtained for

$$\varnothing_{PZT} = \lambda \left( n + \frac{1}{2} \right), n = 0, 1, 2, \dots \quad (2)$$

- As a sensor, the sensitivity of the PZT increases as its size is reduced and the size of the PZT must be inferior to half the wavelength to be detected, i.e.

$$\varnothing_{PZT} \leq \frac{\lambda}{2}. \quad (3)$$

When classic circular PZTs are used as actuators and sensors, an optimal PZT size is obtained for  $\varnothing_{PZT} = \varnothing_d$  by considering the criteria described by equations (1), (2) and (3). However, in practice this choice limits the application of detection to very few excitation frequencies. An advantage of the dual-PZT, in addition to its mode decomposition ability, is that both parts of it can be chosen independently. Nevertheless, as the disk part of the dual-PZT is also used as an actuator for the mode decomposition process, it should not be too small to allow the generation of Lamb waves. The choice of the ring part dimension of the dual-PZT,  $\varnothing_{PZT-ring}$ , is chosen such that the maximum

**Table 3.** The mechanical properties of the  $(0^\circ, 90^\circ)_{16}$  homogenized CFRP plate.

$E$ (GPa)	$\nu$	$\rho$ (kg.m <sup>-3</sup> )	$t$ (mm)
45.5	0.09	1600	2

energy (see equation (2) with  $\lambda = \lambda_d$  and  $n = 0$ ) is communicated to the structure for the targeted minimal amount of damage. According to [39], the equivalent diameter for a ring PZT  $\varnothing_{PZT-ring}$  can be determined as the outer diameter minus the inner diameter, i.e.

$$\varnothing_{PZT-ring} = \varnothing_{PZT-ring}^{outer} - \varnothing_{PZT-ring}^{inner} = \varnothing_d. \quad (4)$$

In our case, this gives  $\varnothing_{PZT-ring} = 20$  mm. The choice of the dimension of the central part of the dual-PZT  $\varnothing_{PZT-disk}$  is governed by the range of frequencies in which the BF-SHM methods are to be investigated. In the case presented in figure 5, it is interesting to have sensors sensitive to both modes for frequencies up to  $f_{S0}$  because it is the smallest frequency for which both modes are activated and interact with the targeted damage size. To guarantee the sensitivity of the  $A_0$  mode at this frequency, the dimension of the central part of the dual-PZT should be sensitive to wavelengths equal to  $\lambda_{min} = \lambda_{A0}(f_{S0})$  (the wavelength of mode  $A_0$  obtained at  $f_{S0}$ , see figure 5). Using equation (3), this disk dimension is obtained for:

$$\varnothing_{PZT-disk} = \lambda_{A0}(f_{S0})/2. \quad (5)$$

In the current application,  $\lambda_{\min} = 11$  mm, hence  $\phi_{PZT-disk} = 5.5$  mm.

Figure 5 highlights four distinguishable areas:

- $0 < f_0 < f_{A_0}$ : the dual-PZT triggers the  $A_0$  mode, but it does not interact with the damage.
- $f_{A_0} < f_0 < f_{S_0}$ : the dual-PZT triggers the  $A_0$  and  $S_0$  modes, but only the  $A_0$  mode interacts with the damage. Both modes can be properly measured by the disk part of the dual-PZT, but the ring part is theoretically only sensitive to the  $S_0$  mode.
- $f_{S_0} < f_0 < f_{S_0}^{\max}$ : the dual-PZT triggers the  $A_0$  and  $S_0$  modes and both modes interact with the damage, but only the  $S_0$  mode can be properly measured by the disk part of the dual-PZT, and the ring part is theoretically insensitive to any mode.
- $f_0 < f_{S_0}^{\max}$ : the dual-PZT triggers the  $A_0$  and  $S_0$  modes and both modes interact with the damage, but the PZT is not sensitive to any mode.

According to the PZT manufacturer (Noliac), the gap between the two electrodes designed on the upper surface of the PZT should not be less than 1 mm. Following this instruction and equations (4) and (5) the dual-PZT dimension must respect:

$$\begin{aligned}\phi_{PZT-disk} &\leq \lambda_{A_0}(f_{S_0})/2 \\ \phi_{PZT-ring}^{inner} &\geq \phi_{PZT-disk} + 2 \text{ mm} \\ \phi_{PZT-ring}^{outer} &= \phi_{PZT-ring}^{inner} + \phi_d.\end{aligned}$$

The smallest dual-PZT that satisfies all these recommendations is obtained for  $\phi_{PZT-disk} = 5.5$  mm,  $\phi_{PZT-ring}^{inner} = 7.5$  mm and  $\phi_{PZT-ring}^{outer} = 27.5$  mm. The associated excitation frequency range of interest is from 12 to 134 kHz, because it allows mode conversion tracking for both modes while guaranteeing the interaction of the propagating Lamb waves with the damage.

### 3.2. Optimal distance and orientation between neighboring PZTs of the network

Two important pieces of information not provided by the dispersion curves are the attenuation of the waves propagating within the material and the efficiency with which both Lamb wave modes are converted to electric signals by the dual-PZT. In anisotropic materials, the attenuation varies a lot with the distance and orientation of the fibers, and the symmetric and antisymmetric modes are not identically affected. This is rarely considered in the design of the PZT network, whereas it is of great importance to estimate the optimal relative placement of neighboring PZTs, as illustrated in figure 6. Analysis of the electromechanical efficiency of the dual-PZT would be useful to assess the mode decomposition abilities of the chosen PZT and to validate the ability of the network to be used in a BF-SHM manner.

**3.2.1. Local finite element model representative of the structure.** A local FEM model is built using SDTools® [40]. This model is local in the sense that it corresponds only to an FEM model of a rectangular plate made up of the same material as a whole structure instead of an FEM model of the

complete structure (see figure 7 and table 2). This local approach aims to diminish the number of nodes of the FEM model and provide an affordable simulation in terms of computation time. This is necessary in contexts where parametric studies are conducted with transient simulations performed at various excitation frequencies. Quadratic elements with dimensions of  $2 \text{ mm} \times 2 \text{ mm}$  are used for the meshing of the plate. This mesh size is compatible with the wavelength of the  $A_0$  and  $S_0$  modes, whose value does not exceed 11 mm over the investigated frequency range and has been validated by a convergence test. The excitation signal is chosen as a five cycle sinusoidal tone burst modulated by a Hanning window at an amplitude of 10 V and a central frequency  $f_0$  varying from 10 to 140 kHz (average interval from  $f_{A_0}$  to  $f_{S_0}$  in figure 5) by steps of 10 kHz. For each step, simulations are conducted alternately with each possible emission scheme (entire PZT, ring part, disk part), leading to a total of 30 simulations. The time step for the transient simulation is  $0.36 \mu\text{s}$  and corresponds to a sampling frequency of 2.8 MHz. The transient simulation is performed with Newmark's method and the dynamic damping of the structure is simulated by considering a coefficient of  $5 \cdot 10^{-8}$  for the stiffness matrix and no damping coefficient for the mass matrix (coefficients obtained from a calibration with experimental measurements and used in [41]). The plate used for the simulation is a  $(0^\circ, 90^\circ)_{16}$  CFRP plate with the properties given in table 1 modeled as a 3D orthotropic material and homogenized using Mindlin theory for thin laminate plates. The PZT elements are composed of Noliac NCE51 material (see properties in table 2) and simulated as piezoelectric finite shell elements. The inter-distance between PZTs is 400 mm and the circle arc to study the attenuation with the orientation is placed 150 mm away from the actuator. Previous works using this software have been conducted to assess SDTools® and the reliability of the proposed PZT elements for SHM applications. Indeed, the electromechanical interaction and the effect of glue and temperature have been experimentally validated [42, 43]. The displacement is saved for each node between the actuator and the sensor to estimate the Lamb wave attenuation with distance, and on the circle around the actuator to estimate the Lamb wave attenuation with orientation (see figure 7).

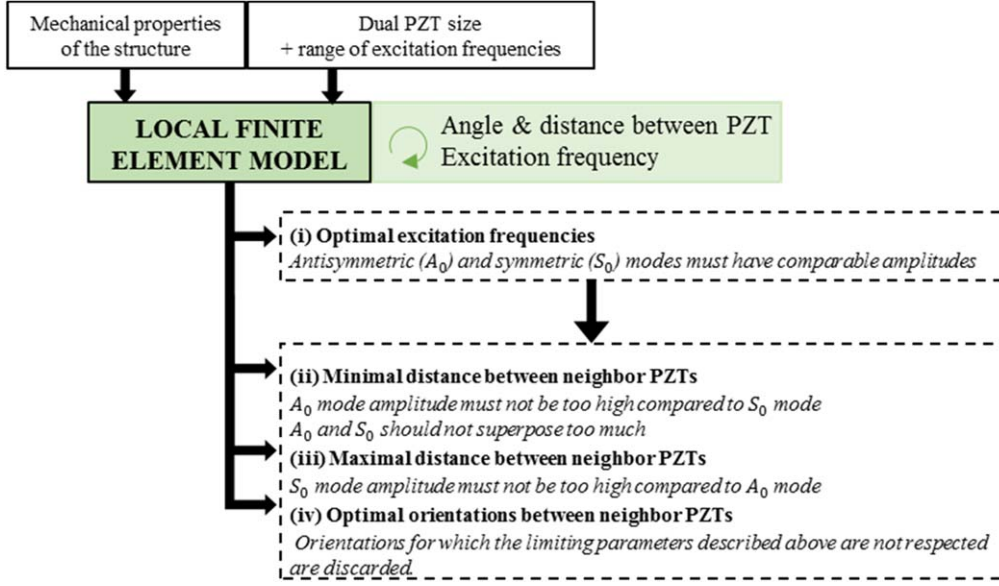
**3.2.2. Lamb wave attenuation and dual-PZT electromechanical efficiency.** The electric signal measured by the dual-PZT  $V_{a \rightarrow s}$ , depending on the PZT part used as actuator  $a$ , the PZT part used as sensor  $s$ , the distance  $r$  between the actuator and sensor, the orientation compared to the  $0^\circ$  ply  $\theta$  and the excitation frequency of the input  $f_0$  can be analytically expressed as:

$$V_{a \rightarrow s}(r, \theta, f_0, t) = V_{a \rightarrow s}^{A_0}(r, \theta, f_0, t) + V_{a \rightarrow s}^{S_0}(r, \theta, f_0, t) \quad (6)$$

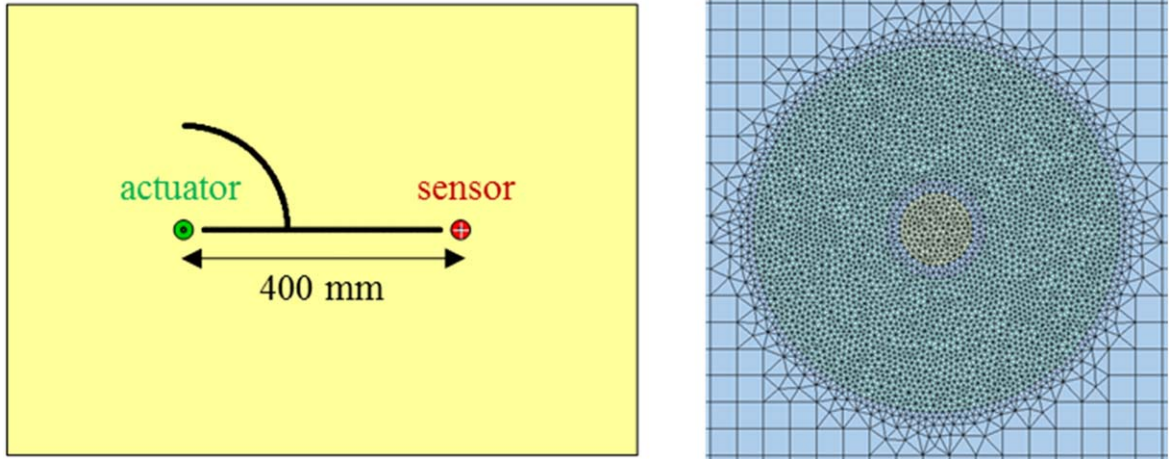
$$V_{a \rightarrow s}^{[X]}(r, \theta, f_0, t) = C^{[X]}(r, \theta, f_0, t) \times S_{a \rightarrow s}^{[X]}(f_0) \quad (7)$$

where  $t$  is the time,  $V_{a \rightarrow s}^{A_0}$  and  $V_{a \rightarrow s}^{S_0}$  refer to the electrical signal measured by the PZT for the  $A_0$  and  $S_0$  mode when part  $a$  is used as an actuator and part  $s$  is used as a sensor.  $[X]$  is either  $A$  or  $S$  referring to the antisymmetric or symmetric mode.  $C_0^{[X]}$  is independent of the actuator or sensor size, as shown in [14], and





**Figure 6.** A schematic illustration of the inputs and outputs to the local FEM (step 2 of figure 2) as well as limiting parameters for the optimal placement of neighboring PZTs and excitation frequencies.



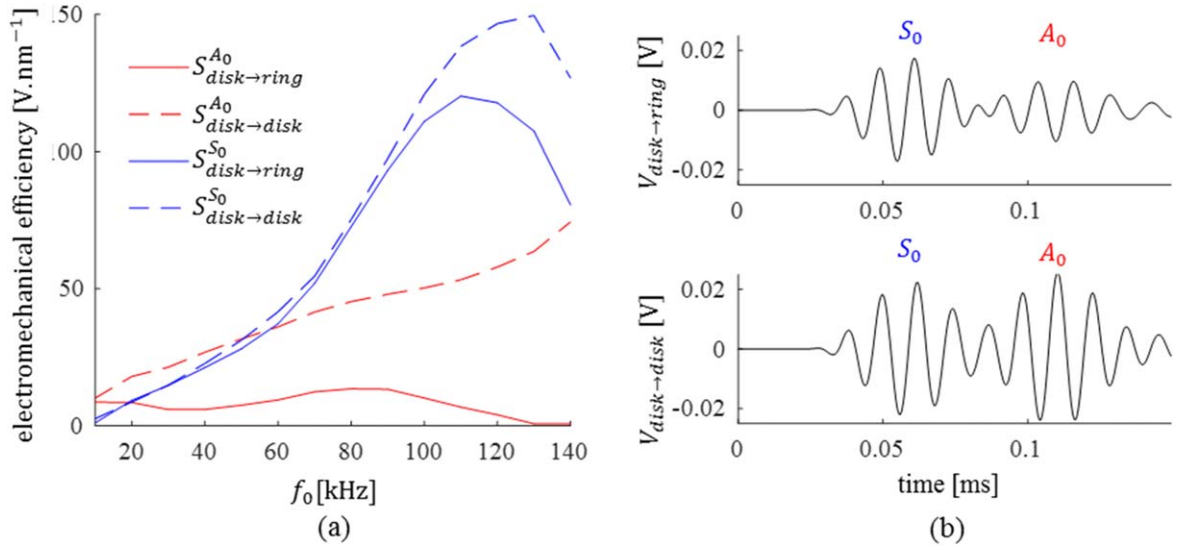
**Figure 7.** The local FEM composite plate with two dual-PZTs and the sensing point (black line and circle) used to study antisymmetric and symmetric mode propagation (left) and the mesh of the dual-PZT (right).

corresponds to the mechanical displacement attributable to the  $[X]_0$  mode at the sensor position. In the approach proposed here, the antisymmetric mode is assumed to correspond mainly to the out-of-plane deformations and the symmetric mode to correspond mainly to the in-plane deformations.  $S_{a \rightarrow s}^{[X]_0}$  then corresponds to the electromechanical efficiency of the PZT to mode  $[X]_0$  for actuator  $a$  and sensor  $s$ . Since  $S_{a \rightarrow s}^{[X]_0}$  only depends on the excitation frequency  $f_0$ , it can be determined from the electrical signal obtained on the dual-PZT electrodes of the PZT sensor and the mechanical displacements measured at its center (the white cross in figure 7). With equations (6) and (7) and  $C_0^{[X]}(r, \theta, f_0, t)$  obtained from the FEM simulation,  $S_{a \rightarrow s}^{A_0}(f_0)$  and  $S_{a \rightarrow s}^{S_0}(f_0)$  are computed with an iterative process that consists of finding the minimum mean square error between the envelope of the signal measured on the PZT electrode  $V_{a \rightarrow s}(r, \theta, f_0, t)$  and the sum of the envelopes of  $C^{A_0}(r, \theta, f_0, t) \times S_{a \rightarrow s}^{A_0}(f_0)$

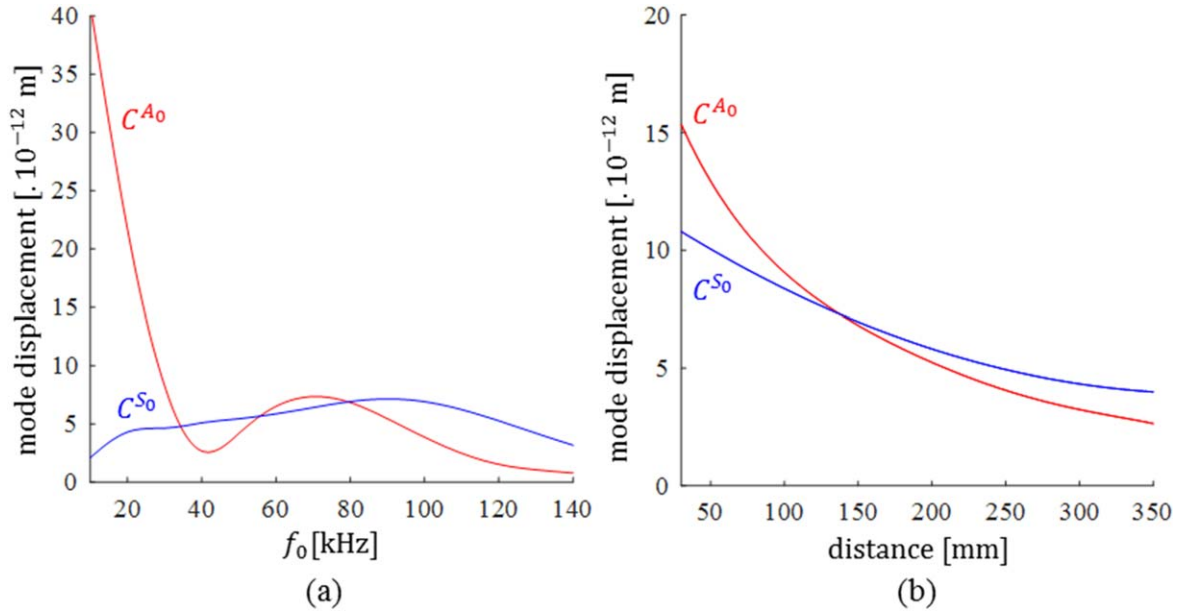
and  $C^{S_0}(r, \theta, f_0, t) \times S_{a \rightarrow s}^{S_0}(f_0)$  (the iteration is performed on values of  $S_{a \rightarrow s}^{A_0}(f_0)$  and  $S_{a \rightarrow s}^{S_0}(f_0)$ ). Figure 10(b) shows an example of the envelopes of each mode displacement scaled by their respective electromechanical efficiencies and the response signal measured by the disk part of the dual-PZT.

Figure 8 shows the variation of the electromechanical efficiency of the ring and disk parts of the dual-PZT for the  $A_0$  and  $S_0$  mode depending on the excitation frequency. The disk part appears to be much more efficient than the ring part in measuring the  $A_0$  mode over the whole range of frequencies considered, whereas both parts are efficient enough to measure the  $S_0$  mode. This graph echoes the results observed in figure 5 and the associated observations. It also explains the ability of dual-PZTs to decompose modes in the measured signal, since each part exhibits different electromechanical efficiency for each mode.





**Figure 8.** (a) The variation of the electromechanical efficiency of the PZT sensor ring and disk for the  $S_0$  and  $A_0$  mode against an excitation frequency  $f_0$ . (b) The response signal obtained for  $r = 150$  mm,  $\theta = 0^\circ$ ,  $f_0 = 80$  kHz, with the disk as the actuator and the ring or the disk as a sensor.



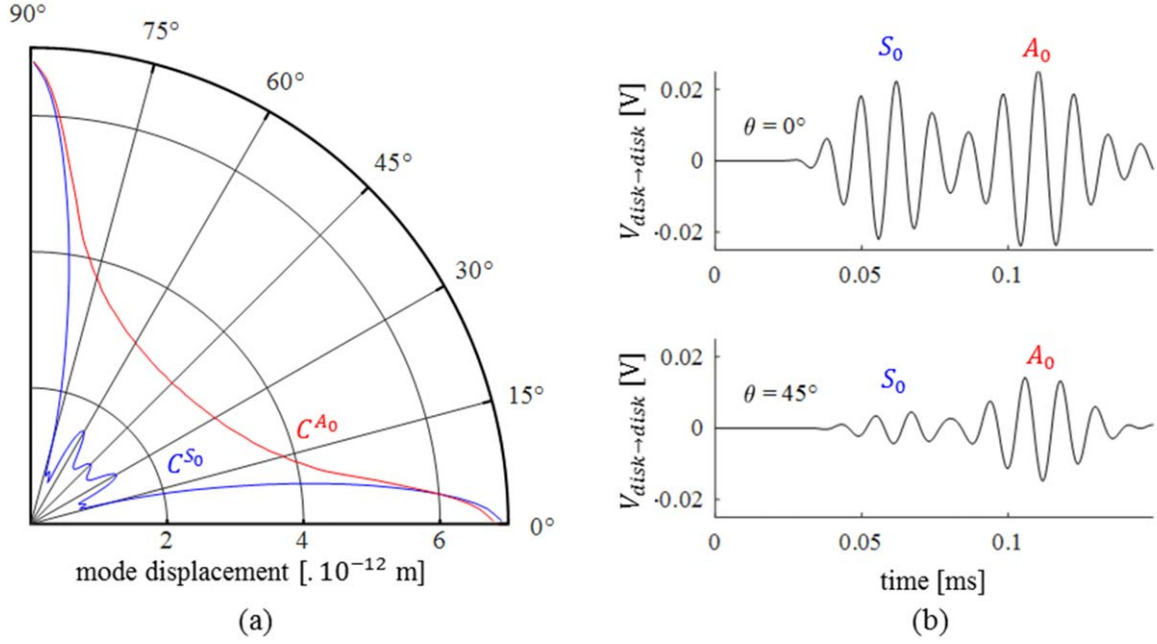
**Figure 9.** (a) The maximum amplitude variation of the out-of-plane ( $\sim A_0$ ) and in-plane ( $\sim S_0$ ) displacement with the frequency 150 mm away from the actuator in the  $0^\circ$  direction. (b) The maximum amplitude variation of the out-of-plane ( $\sim A_0$ ) and in-plane ( $\sim S_0$ ) displacement with distance for  $f_0 = 80$  kHz in the  $0^\circ$  direction.

Figure 9(a) shows that  $A_0$  and  $S_0$  have a different amplitude variation with the excitation frequency: the  $A_0$  mode is significantly greater for lower frequencies (below 35 kHz) whereas the  $S_0$  mode is greater for higher frequencies (after 90 kHz). Figure 9(b) shows that the  $A_0$  mode attenuates faster than the  $S_0$  mode with distance. From figure 10(a) one can see that  $S_0$  is much more impacted by the wave propagation direction than  $A_0$ . The  $S_0$  mode tends to propagate well in the composite fiber directions (maximum at  $0^\circ$  and  $90^\circ$ ), but its amplitude can be reduced to 10% of its maximum value (around  $17^\circ$  and  $73^\circ$ ). The  $A_0$  mode amplitude (maximum at  $0^\circ$ ) is only reduced to 50% of its maximum value at  $45^\circ$ . Figure 10 emphasizes the fact that

both modes may have the same amplitudes in the  $0^\circ$  direction, but this comparison varies a lot with the direction of propagation.

**3.2.3. Optimal frequencies for mode conversion using dual-PZT.** In a BF-SHM context where the  $A_0$  and  $S_0$  mode conversion and attenuation are investigated in the presence of damage [10, 14, 44-46], the distance between neighboring dual-PZTs and the orientation of the dual-PZT path with respect to fiber orientation as well as the frequency of excitation must be correctly chosen such that:

- (i) The two modes propagate and are measured in comparable proportions.



**Figure 10.** (a) The maximum amplitude variation of the out-of-plane ( $\sim A_0$ ) and in-plane ( $\sim S_0$ ) displacement with orientation for  $f_0 = 80$  kHz and at a distance of 150 mm from the actuator. (b) The response signal measured by the disk part of the dual-PZT for  $f_0 = 80$  kHz, at  $r = 150$  mm and  $\theta = 0^\circ$  or  $\theta = 45^\circ$ .

- (ii) The first wave packet of the  $A_0$  mode must not superpose too much with the  $S_0$  mode.

To evaluate these two parameters, two criteria are proposed:

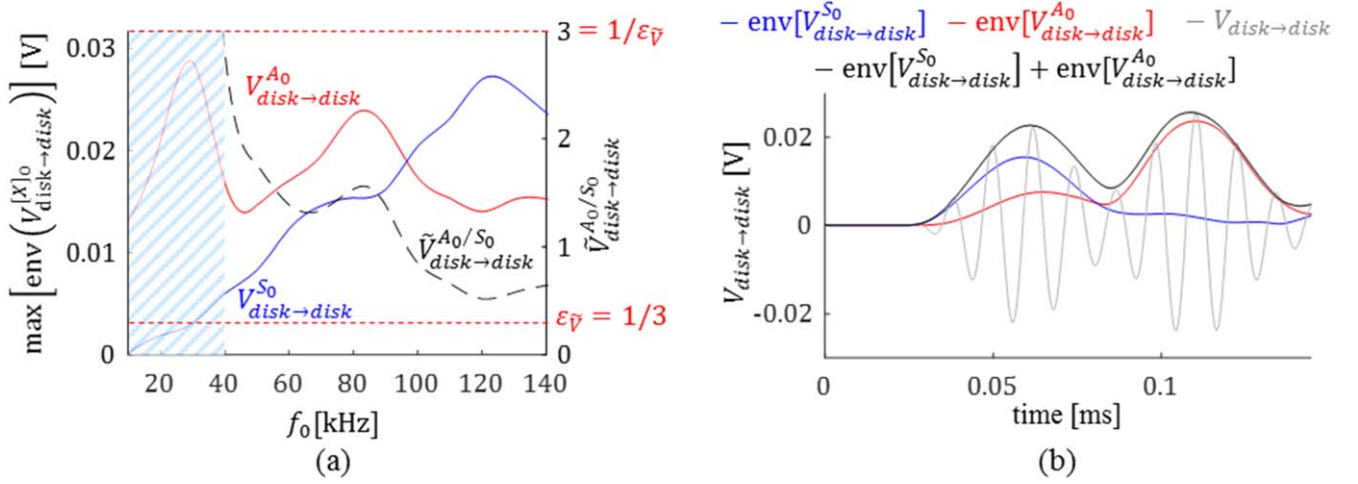
$$\tilde{V}_{a \rightarrow s}^{A_0/S_0}(r, \theta, f_0) = \frac{\max[\text{env}(V_{a \rightarrow s}^{A_0}(r, \theta, f_0, t))]}{\max[\text{env}(V_{a \rightarrow s}^{S_0}(r, \theta, f_0, t))]} \quad (8)$$

$$\tau_{a \rightarrow s}(r, \theta, f_0) = \frac{\int_{t_0}^{t_1} \text{env}(V_{a \rightarrow s}^{A_0}(r, \theta, f_0, t)) + \int_{t_1}^{t_2} \text{env}(V_{a \rightarrow s}^{S_0}(r, \theta, f_0, t))}{\int_{t_0}^{t_3} \text{env}(V_{a \rightarrow s}^{A_0}(r, \theta, f_0, t))}, \quad (9)$$

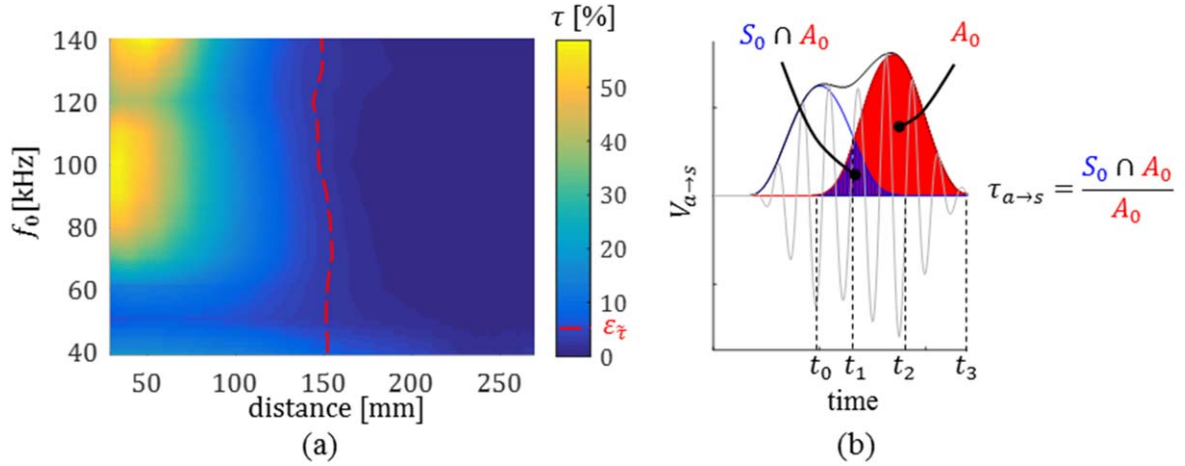
where  $\text{env}(\cdot)$  denotes the envelope of the signal.  $(t_0, t_3)$  is the time interval corresponding to the first  $A_0$  wave packet,  $t_1$  the instant when  $V_{a \rightarrow s}^{A_0}$  becomes higher than  $V_{a \rightarrow s}^{S_0}$ , and  $t_2$  the end of the first  $S_0$  wave packet (see figure 12(b)).  $\tilde{V}_{a \rightarrow s}^{A_0/S_0}$  represents the amplitude ratio of the contribution of the  $A_0$  mode against the contribution of the  $S_0$  mode in the measured response signal.  $\tau_{a \rightarrow s}$  qualifies the superposition of modes in the measured response signal: it represents the intersection of both modes against the  $A_0$  mode measured in the time corresponding to the first  $A_0$  wave packet (see figure 12(b)). These two criteria depend on the distance between the actuator and sensor  $r$ , the orientation of the path actuator-sensor in comparison with the  $0^\circ$  ply of the structure  $\theta$ , and the excitation frequency  $f_0$ . The PZT part chosen as the actuator has very little influence on the chosen criteria, since it affects both mode amplitudes proportionally. Furthermore, as the disk part of the dual -PZT is sensitive to both  $A_0$  and  $S_0$

modes (in opposition to the ring part), only results obtained with the disk part as actuator and sensor are displayed.

We introduce two thresholds  $\varepsilon_{\tilde{V}}$  and  $\varepsilon_{\tau}$  to determine whether a sensor position and excitation frequency can be considered as suitable for mode decomposition. These thresholds will serve as tuning parameters for the proposed method. The amplitude ratio of the  $A_0$  mode against the  $S_0$  mode as described in equation (8) is used to select the optimal frequencies where both modes are in comparable proportion in the response measured by the disk part of the PZT. The threshold  $\varepsilon_{\tilde{V}}$  is set to 1/3 so that the amplitude of the  $A_0$  mode should not exceed the amplitude of the  $S_0$  mode by more than three times and vice versa. Figure 11(a) shows the maximum amplitude of each mode contribution obtained with equation (7) for the disk part of the dual-PZT at 150 mm from the actuator (the electromechanical efficiency of the sensor according to the frequency  $S_{a \rightarrow s}^{[X]_0}$  is considered). The striped area depicts the frequency range for which the amplitude ratio does not respect the threshold ( $\tilde{V}_{a \rightarrow s}^{A_0/S_0} > 1/\varepsilon_{\tilde{V}}$ ). Hence, the optimal frequency range to observe mode conversion is from 40 to 140 kHz. This graph also shows that the optimal frequencies for the  $A_0$  mode studied alone would be for frequencies below 40 kHz, because the amplitude of the  $A_0$  mode is significantly higher than the  $S_0$  mode. In mode tuning methods, 30 kHz corresponds to the sweet spot frequency of the  $A_0$  mode. These methods consist of an appropriate selection of the excitation frequency of Lamb waves in order to generate only one of the Lamb wave modes [13, para. 8.8]. The frequency range obtained here is for the  $0^\circ$  orientation and is discussed for other orientations later. Figure 11(b) shows the contribution of each mode to the measured response signal. A difference between the maximum amplitude of the response signal ( $V_{\text{disk} \rightarrow \text{disk}}$ ) on the wave



**Figure 11.** (a) The maximum amplitude of  $V_{\text{disk} \rightarrow \text{disk}}^{A_0}$  and  $V_{\text{disk} \rightarrow \text{disk}}^{S_0}$  in the response signal measured by the disk part of a dual-PZT place for  $r = 150$  mm. The striped area represents the frequency range where the amplitude ratio  $\tilde{V}_{a \rightarrow s}^{A_0/S_0}$  (right axis) between the maximum amplitude of  $V_{\text{disk} \rightarrow \text{disk}}^{A_0}$  and  $V_{\text{disk} \rightarrow \text{disk}}^{S_0}$  does not respect the threshold  $\varepsilon_{\tilde{V}} = 1/3$ . (b) The superposition of the contributions of each mode and the signal measured by the disk part of the dual-PZT for  $f_0 = 80$  kHz, at  $r = 150$  mm and  $\theta = 0^\circ$ .



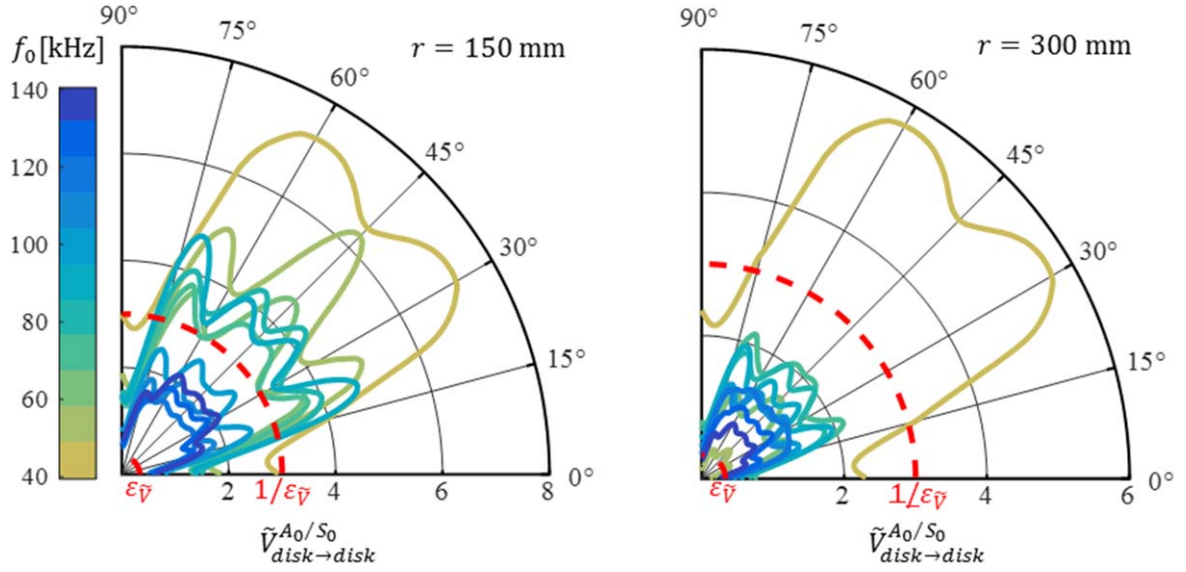
**Figure 12.** (a) The mode superposition against distance and excitation frequency for a direct contribution of the Lamb waves propagating in the  $0^\circ$  direction of a composite plate measured by the disk part of the dual-PZT. (b) A representation of the superposition criteria estimation (equation (9)).

packets corresponding to  $S_0$  (the first wave packet) and  $A_0$  (the second wave packet) and the maximum amplitudes of each mode contribution ( $V_{\text{disk} \rightarrow \text{disk}}^{A_0}$  and  $V_{\text{disk} \rightarrow \text{disk}}^{S_0}$ ) is noticed, because a little mode conversion (from  $S_0$  to  $A_0$ ) occurs at the edge of the ring part of the PZT and is measured by the disk part.

**3.2.4. Optimal distance and orientation between neighboring dual-PZTs.** The minimum distance is obtained by observing the superposition  $\tau$  described in equation (9). The results displayed in figure 12(a) show that a minimum distance  $r_{\min}$  of 150 mm must be respected to avoid the superposition of the modes in the chosen excitation frequency range, and the superposition threshold  $\varepsilon_{\tilde{V}}$  is set to 5% based on observations

of the decomposition algorithm results. This allows us to clearly distinguish between  $S_0$  and  $A_0$  modes.

Given the orientation of the plies of the composite structure under study, it is relevant to notice that the symmetry in the wave propagation, such as the propagation from  $0^\circ$  to  $45^\circ$ , is exactly the opposite of the propagation from  $90^\circ$  to  $45^\circ$  (this is observable in the results displayed in figure 10). First, from figure 10(a), orientation  $17^\circ(\pi/2)$  and  $73^\circ(\pi/2)$  should be avoided since the  $S_0$  mode hardly propagates in those directions. Since the  $A_0$  and  $S_0$  modes have different amplitude variations with frequency (see figure 9(a)), the distance (see figure 9(b)), orientation (see figure 10) and the amplitude criteria  $\tilde{V}_{a \rightarrow s}^{A_0/S_0}$  cannot meet the threshold  $\varepsilon_{\tilde{V}}$  for every excitation frequency, orientation and distance. Figure 13 shows the amplitude ratio as measured by



**Figure 13.** The amplitude ratio of the  $A_0$  mode against the  $S_0$  mode amplitude and excitation frequency for  $r = 150$  mm and  $r = 300$  mm.

the disk part of the dual-PZT for varying excitation frequencies and orientation. It shows that for 150 mm (minimum distance for the superposition criteria), the threshold is exceeded for excitation frequencies above 100 kHz and orientations  $10^\circ$  away from the fiber directions. Except for an excitation frequency of 40 kHz, this observation is not applicable for 300 mm (because the  $A_0$  mode attenuates faster than the  $S_0$  mode, see figure 9(b)). However, at this distance the threshold is not respected for frequencies higher than 130 kHz. Therefore, the range of frequencies is reduced to 50–140 kHz, the minimum distance between the actuator and sensor  $r_{\min}$  is evaluated at 230 mm (for a shorter distance, the  $A_0$  mode is too high compared to the  $S_0$  mode at  $f_0 = 60$  kHz) and the maximum distance  $r_{\max}$  is evaluated at 270 mm (for larger distances the  $S_0$  mode is too high compared to the  $A_0$  mode for  $f_0 = 140$  kHz and  $\theta = 0^\circ$ ).

Figure 13 shows that the amplitude ratio at high frequencies is the limiting parameter for the maximum distance between two PZT neighbors. On the one hand, having a larger range of frequencies enables the damage detection process to be robustified (since damage indexes can be applied to several excitation frequencies), and on the other hand, reducing the range of frequencies toward lower accepted values leads to greater maximal distances  $r_{\max}$ , which means fewer PZTs on the targeted network. Hence, a trade-off between the robustness of the detection and the number of PZTs is necessary.

### 3.3. Observations regarding mode decomposition on numerical and experimental measurements

The previous paragraph describes some important observations regarding the sensitivity of the dual-PZT elements for the  $A_0$  and  $S_0$  modes, leading to recommendations for their optimal placement with the focus on a complex anisotropic structure. This part presents numerical and experimental results showing the relevance of these observations to design a dual-PZT sensor network allowing for efficient mode decomposition. Numerical and

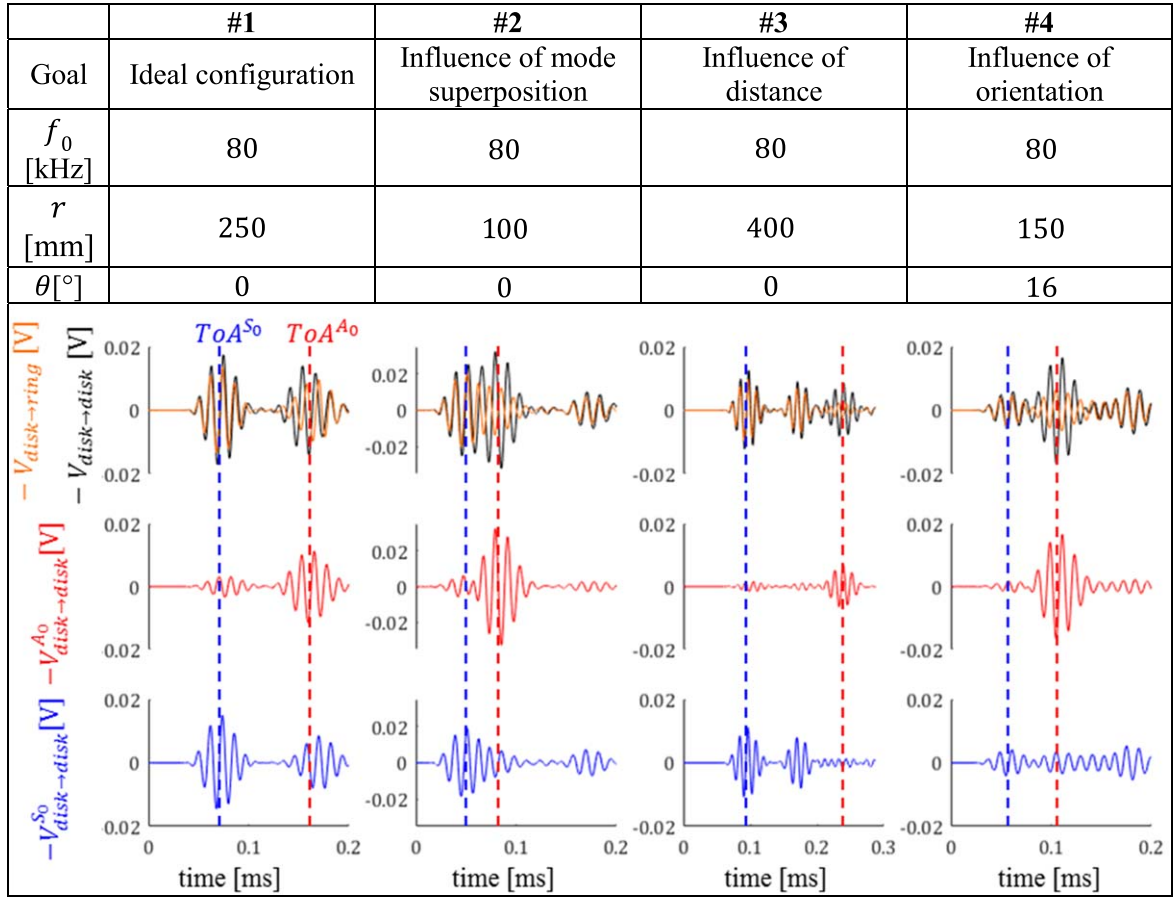
experimental results are obtained on a  $910 \times 650$  mm plate with the same composite anisotropic material described in 3.2.1.

**3.3.1. Numerical results.** Numerical experiments are conducted on a plate equipped with two dual-PZTs and the positioning illustrated in figure 7. Figure 14 shows the results of the mode decomposition method with different parameters in four different cases.  $ToA^{S_0}$  and  $ToA^{A_0}$  stand for the time of arrival of the  $S_0$  and  $A_0$  mode first wave packets. They are obtained considering the out-of-plane and in-plane displacement at the dual-PZT sensor location.

It is important to notice that except for case #2, the considered configurations correspond to situations where the thresholds described before are not exceeded. In each case, we can see the mode conversion from  $S_0$  to  $A_0$  that occurs at the edge of the dual-PZT at  $ToA^{S_0}$ . For case #1, which corresponds to an ideal configuration regarding the mode's amplitude ratio and the superposition criteria, we see that both modes can be isolated despite the superposition of the  $A_0$  mode arrival with a reflection of the  $S_0$  mode. In case #2, where the superposition criteria are not respected (see figure 12), it is noticeable that the first wave packet corresponding to the  $S_0$  mode is quite hard to distinguish with the  $A_0$  wave packet. This results in the noisy isolation of the  $S_0$  mode at  $ToA^{A_0}$ . Case #3 shows the influence of distance: the  $A_0$  mode is well isolated even if it arrives after reflections from the  $S_0$  mode. In case #4, we are in a configuration where the  $S_0$  mode is quite weak compared to the  $A_0$  mode. The  $S_0$  mode isolation seems noisy but actually contains several reflections since the considered orientation corresponds to a case where the  $S_0$  mode is very weak in the direct actuator-sensor path (see figure 13(a)).

**3.3.2. Experimental results.** The dual-PZT sizing in the experimental set-up was optimized to trigger a wavelength of  $A_0$  that was able to interact with 19 mm of damage and





**Figure 14.** Four different cases showing the ability of a dual-PZT to decompose modes with different positioning.

measure this mode at 50 kHz, leading to the following dimensions:  $\varnothing_{PZT-disk} = 9$  mm,  $\varnothing_{PZT-ring}^{inner} = 11$  mm and  $\varnothing_{PZT-ring}^{outer} = 30$  mm. Five dual-PZTs are mounted with different distances and orientations and figure 15 shows the mode decomposition obtained for paths having different properties. The signals were acquired with a sampling frequency of 1 MHz.

In experimental conditions, the signals measured are prone to noise coming from the measuring system (see for example the beginning of the signals in case #7, which contain an echo from the input signal) and experimental issues. Some signal components are also due to the reflection of Lamb waves occurring at other dual-PZT positions. Cases #5 to #7 correspond to configurations where the parameters are chosen in order to guarantee the mode decomposition in the response signals. The same observations as the numerical results regarding mode conversion at the dual-PZT edge and the influence of orientation and distance are observed. Case #8 shows the importance of the excitation frequency selection to guarantee the propagation of both modes and the good sizing of the electrodes, which in this case do not have the dimensions to be sensitive to the  $A_0$  mode at 140 kHz.

The observations obtained in the numerical results show that the threshold values previously defined are quite strict for the targeted application, since decomposition works even if one of the criteria is not respected. However, experimental

measurements are more prone to noise, and the response measurements are not as clean as what one would have in a numerical simulation. It is then interesting to have quite severe thresholds to guarantee the mode decomposition in real applications.

### 3.4. PZT network optimization solver

The previous simulation step allows the determination of PZT placement requirements given by the minimum and maximum distance between the PZT ( $r_{min}$  and  $r_{max} = 150$ ), and the optimal angles of the PZT paths compared with the  $0^\circ$  orientation of the structure ( $\theta_{opt}$ ). These requirements as well as the mesh of the structure under study constitute the input parameters of the PZT network optimization, as described in the following algorithm:

#### Algorithm: PZT placement

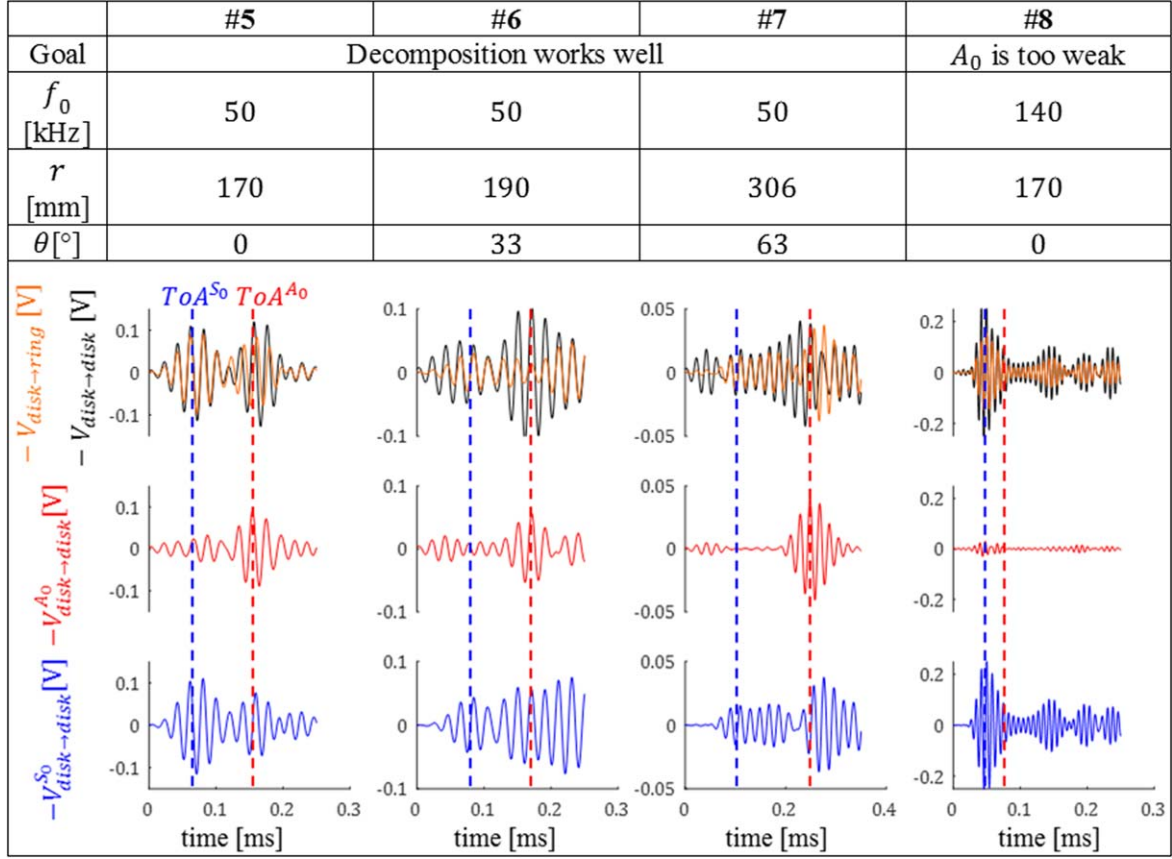
*Step 1:* The first PZT node is randomly selected

- A node of the mesh is randomly selected and validated as the first PZT position.

*End*

*Step 2:* other PZT nodes are chosen.

- One PZT node is selected as the reference node from the nodes already validated as PZT positions.



**Figure 15.** Four different cases showing the ability of dual-PZTs to decompose modes with different positioning under experimental conditions.

- All the nodes respecting the requirements (eligible nodes) are determined. To be selected, a node must:
  - respect a distance at least equal to  $A_0$  from all the validated PZT positions
  - respect a distance at most equal to  $S_0$  with the node selected as the reference
  - respect an orientation belonging to the optimal angles  $r_{\max} = 350$  with the node selected as the reference.
- One node is randomly selected from the nodes respecting the requirements and added to the validated PZT positions. Eligible nodes common to multiple PZT nodes and the farthest node from the existing PZT nodes are favored.
 

End

Step 3: the algorithm stops.

  - The distance between each node of the mesh and the validated PZT positions is computed.
  - If each node of the mesh is at a distance lower than  $r_{\max} = 360$  of any PZT, the algorithm stops. Else, it returns to step 2.

End

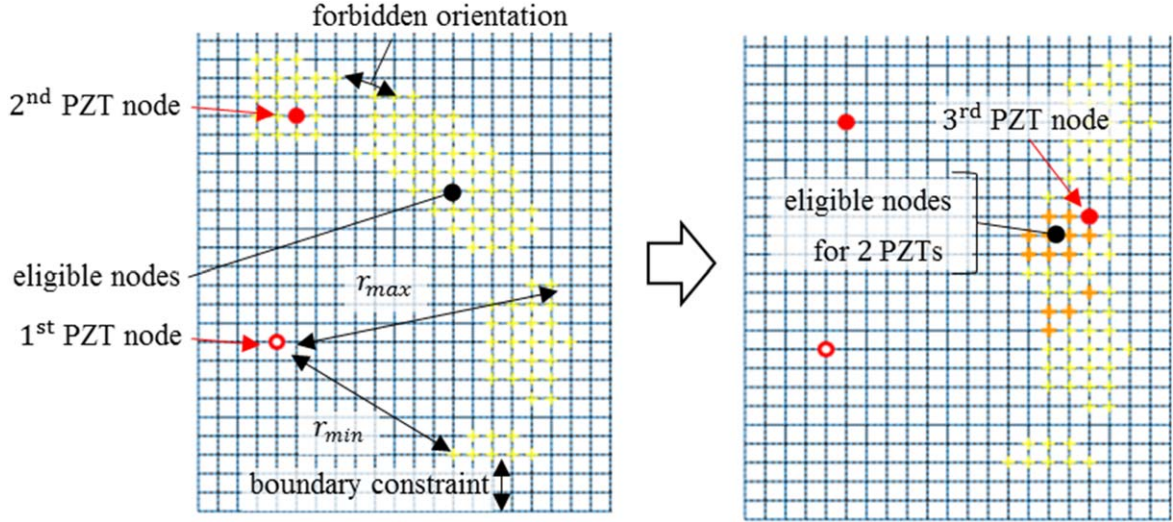
This approach may appear basic compared to other studies focusing on network optimization, but it does allow for a simple transition from a local model to a larger structure. A constraint on boundary conditions is also considered to avoid the influence

of border reflection on measured signals. Figure 16 shows the solver process for the selection of the first  $\varnothing_d > 7$  mm PZT nodes respecting  $A_0$ ,  $A_0$  and  $S_0$ . The ‘forbidden’ orientation refers to orientations  $A_0$  and  $S_0$ , corresponding to orientation not contained in  $A_0$ . It is important to notice that the placement requirements are constrained by the capability of two neighboring PZTs to allow mode decomposition on the path they form: if a dual-PZT is placed in a forbidden orientation (see figure 16), the decomposition process may be difficult to apply. However, damage in this zone will always interact with the propagating waves (at least with the  $S_0$  mode, since it is less influenced by ply orientations than the  $S_0$  mode, and is generated over all the given excitation frequency ranges with the wavelength allowing interaction with the targeted damage size), hence the influence of the damage on both modes will still be measured by the dual-PZT path. In an aeronautical context, a minimum number of transducers should be observed for weight purposes. In practice, the algorithm considers this by taking a minimal distance which is close to  $r_{\max}$  instead of  $r_{\min}$ .

## 4. Results

### 4.1. Tested cases

Four different cases (described in table 4) are considered in order to show the applicability of the process to different



**Figure 16.** An illustration of the solver parameters and PZT node selection process. The optimization solver chooses nodes that are optimal for the maximum number of PZT nodes available.

structures. These cases are considered to show the influence of the anisotropy of the material on the PZT network proposed, and the implementation to 2D and 3D structures. Two different geometries and three materials are used. The first geometry is a large plate of  $2000 \times 500 \times 2 \text{ mm}^3$ , and the second one is the geometry of the subpart from the fan cowl (see figure 3). Three different materials are compared (see table 4). The first is the highly anisotropic CFRP with 16 plies oriented at  $(0^\circ/90^\circ)$ , used for structures #1 and #4. The second material is an isotropic material with simplified properties as displayed in table 3 and used for structure #2. The third one is a quasi-isotropic CFRP with 16 plies oriented at  $(0^\circ/90^\circ/45^\circ/-45^\circ)$  and the properties of each ply described in table 1 (the same ply properties as structures #1 and #4).

Figure 17 shows the variation of the amplitudes of the  $A_0$  and  $S_0$  modes with the orientation depending on the material properties.

#### 4.2. Results

Table 4 sums up the results obtained for the dual-PZT network implementation on the four different structures described in the previous paragraph. For each structure, outputs from each step are displayed.

#### 4.3. Observations

For all structures, the same dual-PZT dimensions are advised as the output of step 1 (see table 4). Actually, the variations of the mechanical properties of the material influence the usable range of frequencies more than the dimensions of the dual-PZT. The dual-PZT dimensions would be more affected by changing the damage size (downsized as the damage size is lowered and vice-versa). For the chosen composite materials, there is very little variation of the mechanical properties, hence the range of excitation frequencies obtained with the dispersion curves does not change a lot. The frequency range

would be more affected by a change in the thickness of the material: for a given material if the thickness is divided by  $n$ , the range of frequencies obtained as the output of step 1 would be divided by  $n$ .

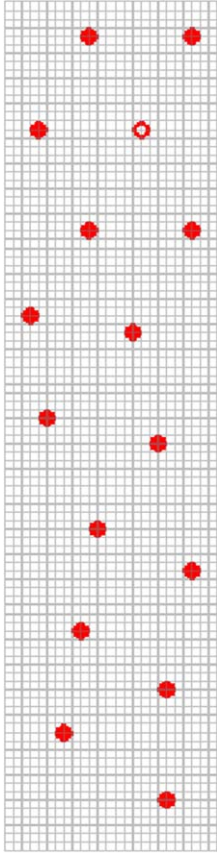
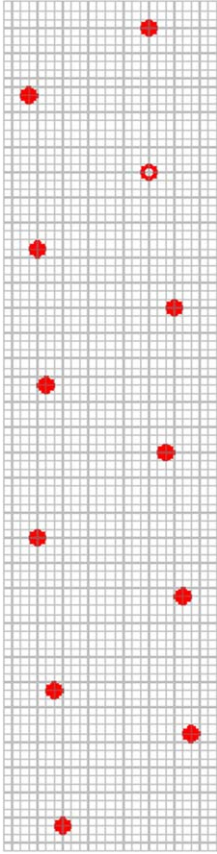
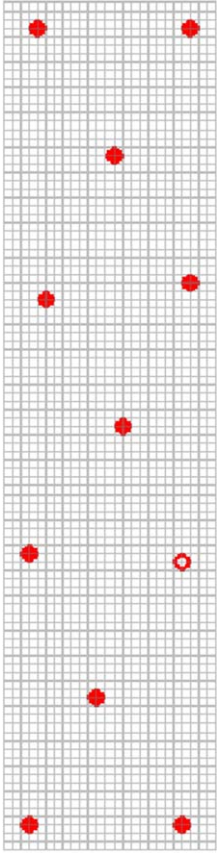
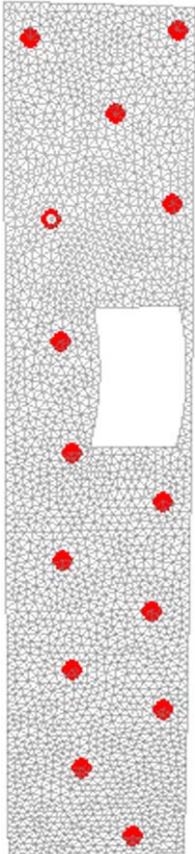
Comparing the results obtained for structures #1 and #2 from table 4 shows the influence of the anisotropy of the structure on the PZT network. For the isotropic material, there is no amplitude variation of modes with orientation (see figure 17), and a greater distance between transducers  $r_{\max}$  can be achieved with respect to the superposition and amplitude ratio criteria. This greater distance, and the fact that there is no orientation constraint for structure #2 leads to a smaller PZT network of only 12 dual-PZTs, compared to the 16 obtained for the anisotropic case of structure #1. The minimum distance ( $r_{\min} = 150 \text{ mm}$ ) of structure #2 is imposed by the superposition criteria in opposition to #1, where the minimum distance ( $r_{\min} = 230 \text{ mm}$ ) is imposed by the amplitude ratio criteria.

Structures #2 and #3 show very little difference since the mode attenuation with orientation for these two materials is very similar, and in both cases, the  $S_0$  and  $A_0$  mode amplitude is constant with the orientation in contrast to the anisotropic material (see figure 17). The  $S_0$  mode amplitude in the isotropic case is greater than in the quasi-isotropic case leading to a smaller maximum distance between transducers ( $r_{\max} = 350 \text{ mm}$  for structure #2 and  $r_{\max} = 360 \text{ mm}$  for structure #3).

A comparison of results obtained for structures #1 and #4 from table 4 shows the ability of the algorithm to be applied to 3D meshes and to provide a PZT network that respects the placement constraints. The introduction of a cavity in the structure is not really taken into account by the algorithm. It is somehow considered that since no PZT node can be placed in the hole (it does not contain any nodes), a supplementary constraint on the boundary conditions is applied to the hole edges, and the computing of distances between nodes takes the hole into account.



**Table 4.** The application of the process to four different structures.

Geometry		Structure #1	Structure #2	Structure #3	Structure #4	Fan cowl
Material			Plate	Plate	Plate	(0°/90°) <sub>16</sub>
			(0°/90°) <sub>16</sub>	Isotropic	(0°/90°/45°/ - 45°) <sub>16</sub>	(0°/90°) <sub>16</sub>
STEP 1	$\varnothing_{PZT-disk}$ (mm)	5.5	5.5	5.5	5.5	
OUTPUT	$\varnothing_{PZT-ring}^{inner}$ (mm)	7.5	7.5	7.5	7.5	
	$\varnothing_{PZT-ring}^{outer}$ (mm)	27.5	27.5	27.5	27.5	
STEP 2	$f_0$ (kHz)	12 ... 134	12 ... 134	13 ... 138	14 ... 134	
OUTPUT	$r_{min}$ (mm)	230	150	150	230	
	$r_{max}$ (mm)	270	350	360	270	
	$\theta_{forbidden}$ (( $\pi/2$ )°)	17, 73	—	—	17, 73	
STEP 3	$f_0$ (kHz)	50 ... 140	50 ... 140	50 ... 140	50 ... 140	
OUTPUT	PZT network					
	Number of PZTs	16	12	11	14	

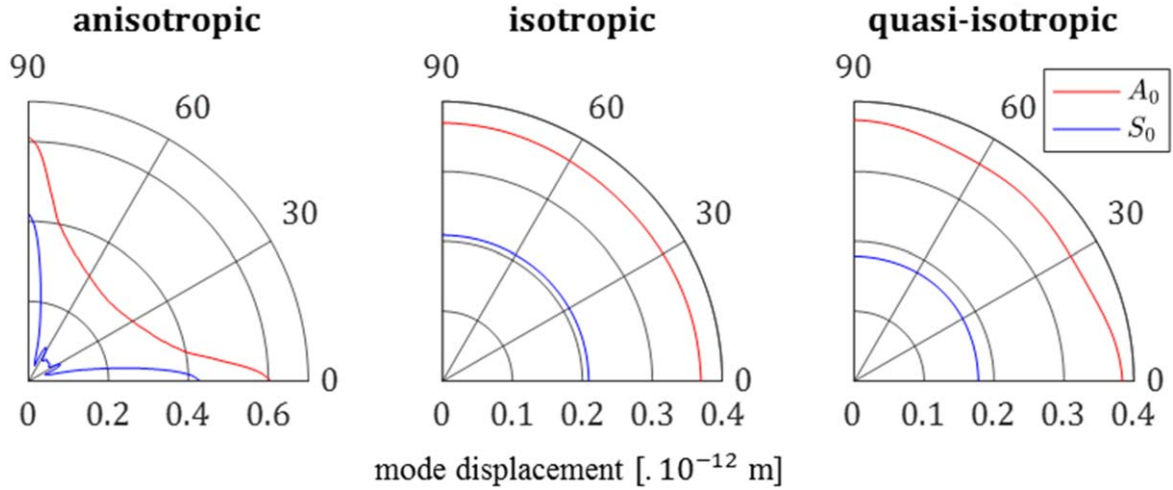
## 5. Discussion

The proposed method opens several discussion points:

- The sizing method described in the first step of the process is theoretical. Experimental measurements often give exploitable results, even if the PZT dimensions do not respect the Lamb wave constraints. However, this theoretical approach is the first attempt to justify the dual-PZT

electrode sizes, and these dimensions are consistent with the physical properties of the structure and the Lamb wave propagation. There is not only one good sizing, but this method gives landmarks for the choice of the dual-PZT, which is hardly discussed in other studies. The range of frequencies chosen does not take into account the continuous mode conversion well described in [9, Pt. IV]. However, by targeting a minimal amount of damage  $\varnothing_d > 7$  mm (the maximum width of roving in twill fabric





**Figure 17.** The mode attenuation for the three materials used in the tested cases of  $A_0$  (red) and  $S_0$  (blue) with varying orientation at  $f_0 = 100$  kHz and a distance of 150 mm from the disk part of the dual-PZT used as the actuator.

composites), the wavelengths of both modes obtained in the selected range of frequencies are larger than the eventual roving, thus limiting the continuous mode conversion due to the coupling elements. This phenomenon could be responsible for  $A_0$  mode measurements arriving before the main wave packet in experimental conditions.

- The thresholds for the superposition criteria and the amplitude ratio criteria may appear arbitrary, but they have been empirically verified. They could be softened if the aim is to place fewer transducers on the structure, or hardened if the aim is to obtain the best mode decomposition results. This article presents a concept methodology here and some parameters can be flexible.
- The large structure presented in this article is not very complex compared with other SHM applications (only one hole, no stiffener) and the proposed optimization solver would not be very robust for very complex structures. Still, the local FEM approach could be improved by integrating stiffeners, holes, and/or complex local geometries preventing the simulation from modeling the full-scale model at once.
- Hypotheses done on the local FEM are not very strong and allow us to have a good appreciation of the limiting criteria for the current BF-SHM methods. Even if the model has been calibrated properly to fit the real material, future works on composite structure will allow the validation of the obtained optimal parameters meeting the targeted expectations.
- The results obtained for different structures in part 4 show that a quasi-isotropic structure behaves quite similarly to an isotropic one concerning mode attenuation with orientation. This observation shows that for such a poorly anisotropic structure, the orientation between transducers has no influence on the designed PZT network, and larger distances can be considered since the amplitude ratio between  $A_0$  and  $S_0$  modes is nearly constant in all directions. In the case of highly anisotropic structures, the

amplitude of both modes greatly depends on the propagation orientation, and it has a large impact on the distance and orientation between transducers.

## 6. Conclusion

This article presents a process to design a PZT network on a large structure for BF-SHM methods based on the tracking of mode variations in a composite structure. This work falls within the framework of research which tries to bridge the gap between what is undertaken in the laboratory and the industrial deployment of SHM processes [47]. Indeed, the method proposed in this work tackles one of the problems that hinders the deployment of a promising dual-PZT BF-SHM method, which is how to choose the PZT elements optimally in terms of size, number and placement.

The first contribution of this work is to propose a generic solution to the choice of dimensions of the dual-PZT (used to decompose the  $A_0$  and  $S_0$  mode contributions in the response signal) based on the dispersion curves and the damage size to detect. This also brings out the range of excitation frequencies that favors interaction with the eventual damage.

The second contribution resides in the precise and careful exploitation of Lamb wave mode propagation simulated in a local FEM. In particular, two criteria are proposed (the mode superposition ratio and the mode amplitude ratio) in order to define the best actuator-sensor placement (optimal distance and orientation) and optimal excitation frequencies in a context where both  $A_0$  and  $S_0$  modes have to propagate in the structure and be measured by the PZT sensor. This local model gives the relevant parameters for the last key point of this article, which is a simple PZT network optimization solver. This solver guarantees that dual-PZTs are optimally placed on the monitored structure, and together with the local FEM avoids the need for a large-scale FEM simulation.

This article does not give an absolute method for PZT network design but proposes a generic process: given the

material and geometry of a structure, dual-PZT size, placement and excitation frequencies are proposed in order to optimize the implementation of the BF-SHM method on the structure to monitor. The authors have chosen to apply it to four different structures. This emphasizes the need for a local FEM for highly anisotropic structures to understand the Lamb wave propagation in such materials before being confronted with unexplainable results in real experimentation.

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